

# A Comparative Study of Nucleon Structure in Light-Front Quark Models in AdS/QCD

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## Abstract

We present the nucleon electromagnetic form factors, using the light-front wave functions of a quark-diquark model for nucleon predicted by the soft-wall model of AdS/QCD. The results are compared with the soft-wall AdS/QCD model. Then we show a comparative study of the nucleon charge and anomalous magnetization densities in the transverse plane. Flavor decompositions of the form factors and transverse densities are also presented.

## LF quark-diquark model

- In the quark-diquark model, nucleon is considered to be a bound state of a single quark and a scalar diquark state.
- For a spin  $\frac{1}{2}$  composite system the Dirac and Pauli form factors  $F_1(Q^2)$  and  $F_2(Q^2)$  are identified to the helicity-conserving and helicity-flip matrix elements of the vector current  $J^+$

$$\langle P+q, \uparrow | \frac{J^+(0)}{2P^+} | P, \uparrow \rangle = F_1(Q^2),$$

$$\langle P+q, \uparrow | \frac{J^+(0)}{2P^+} | P, \downarrow \rangle = -(q^1 - iq^2) \frac{F_2(Q^2)}{2M_n},$$

- Using the two particle Fock states for  $J^z = +\frac{1}{2}$  and  $J^z = -\frac{1}{2}$ , the Dirac and Pauli form factors for the quarks can be written in the light-front representation as

$$F_1^q(Q^2) = \int_0^1 dx \int \frac{d^2\mathbf{k}_\perp}{16\pi^3} \left[ \psi_{+q}^{+\lambda}(x, \mathbf{k}'_\perp) \psi_{+q}^{+\lambda}(x, \mathbf{k}_\perp) + \psi_{-q}^{+\lambda}(x, \mathbf{k}'_\perp) \psi_{-q}^{+\lambda}(x, \mathbf{k}_\perp) \right],$$

$$F_2^q(Q^2) = -\frac{2M_n}{q^1 - iq^2} \int_0^1 dx \int \frac{d^2\mathbf{k}_\perp}{16\pi^3} \left[ \psi_{+q}^{+\lambda}(x, \mathbf{k}'_\perp) \psi_{+q}^{-\lambda}(x, \mathbf{k}_\perp) + \psi_{-q}^{+\lambda}(x, \mathbf{k}'_\perp) \psi_{-q}^{-\lambda}(x, \mathbf{k}_\perp) \right],$$

where  $\mathbf{k}'_\perp = \mathbf{k}_\perp + (1-x)\mathbf{q}_\perp$ . For the frame  $q = (0, 0, \mathbf{q}_\perp)$ ,  $Q^2 = \mathbf{q}_\perp^2$ .

- $\psi_{\lambda q}^A$  are the LFWF with nucleon helicity  $\Lambda$  and quark  $q$  with helicity  $\lambda$ . Wavefunctions are constructed from LFWF predicted by AdS/QCD [1].

## AdS/QCD Model I

- This model is originally proposed by Brodsky and Téramond [2].
- Action for Dirac fields in AdS<sub>5</sub> in this model

$$S = \int d^4x dz \sqrt{g} \left( \frac{i}{2} \bar{\Psi} e_A^M \Gamma^A D_M \Psi - \frac{i}{2} (D_M \bar{\Psi}) e_A^M \Gamma^A \Psi - (\mu + U(z)) \bar{\Psi} \Psi \right)$$

where  $e_A^M = (z/R)\delta_A^M$ ,  $\sqrt{g} = (R/z)^5$ , and  $\Gamma_A = \{\gamma_\mu, -i\gamma_5\}$ .

- The action leads to the AdS solutions  $\psi_+(z)$  and  $\psi_-(z)$  (correspond to different orbital angular momentum  $L^z = 0$  and  $L^z = +1$ ),

$$\psi_+(z) = \frac{\sqrt{2}\kappa^2}{R^2} z^{7/2} e^{-\kappa^2 z^2/2}, \quad \psi_-(z) = \frac{\kappa^3}{R^2} z^{9/2} e^{-\kappa^2 z^2/2}.$$

- The Dirac (spin non-flip) form factors

$$F_1^p(Q^2) = R^4 \int \frac{dz}{z^4} V(Q^2, z) \psi_+^2(z),$$

$$F_1^n(Q^2) = -\frac{1}{3} R^4 \int \frac{dz}{z^4} V(Q^2, z) (\psi_+^2(z) - \psi_-^2(z)).$$

- The Pauli (spin flip) form factor is modeled as

$$F_2^{p/n}(Q^2) = \kappa_{p/n} R^4 \int \frac{dz}{z^3} \psi_+(z) V(Q^2, z) \psi_-(z).$$

- We use the value of  $\kappa = 0.4 \text{ GeV}$  and  $V(Q^2, z)$  is the bulk to boundary propagator, related to electromagnetic field.

## AdS/QCD Model II

- This model has been proposed by Abidin and Carlson [3].
- They add the following extra gauge invariant term to the action

$$\eta \int d^4x dz \sqrt{g} \frac{i}{2} \bar{\Psi} e_A^M e_B^N [\Gamma^A, \Gamma^B] F_{MN} \Psi.$$

- This term produces the Pauli (spin flip) form factors  $F_2$ .
- Also provides a contribution to the Dirac (spin non-flip) form factor  $F_1$ .

$$F_1^p(Q^2) = C_1(Q^2) + \eta_p C_2(Q^2), \quad F_1^n(Q^2) = \eta_n C_2(Q^2),$$

$$F_2^p(Q^2) = \eta_p C_3(Q^2), \quad F_2^n(Q^2) = \eta_n C_3(Q^2).$$

- $\eta_p C_2(Q^2)$  is an additional contribution to the Dirac form factor.
- In this model, the value of  $\kappa = 0.350 \text{ GeV}$  and the other parameters  $\eta_p = 0.224$  and  $\eta_n = -0.239$ .

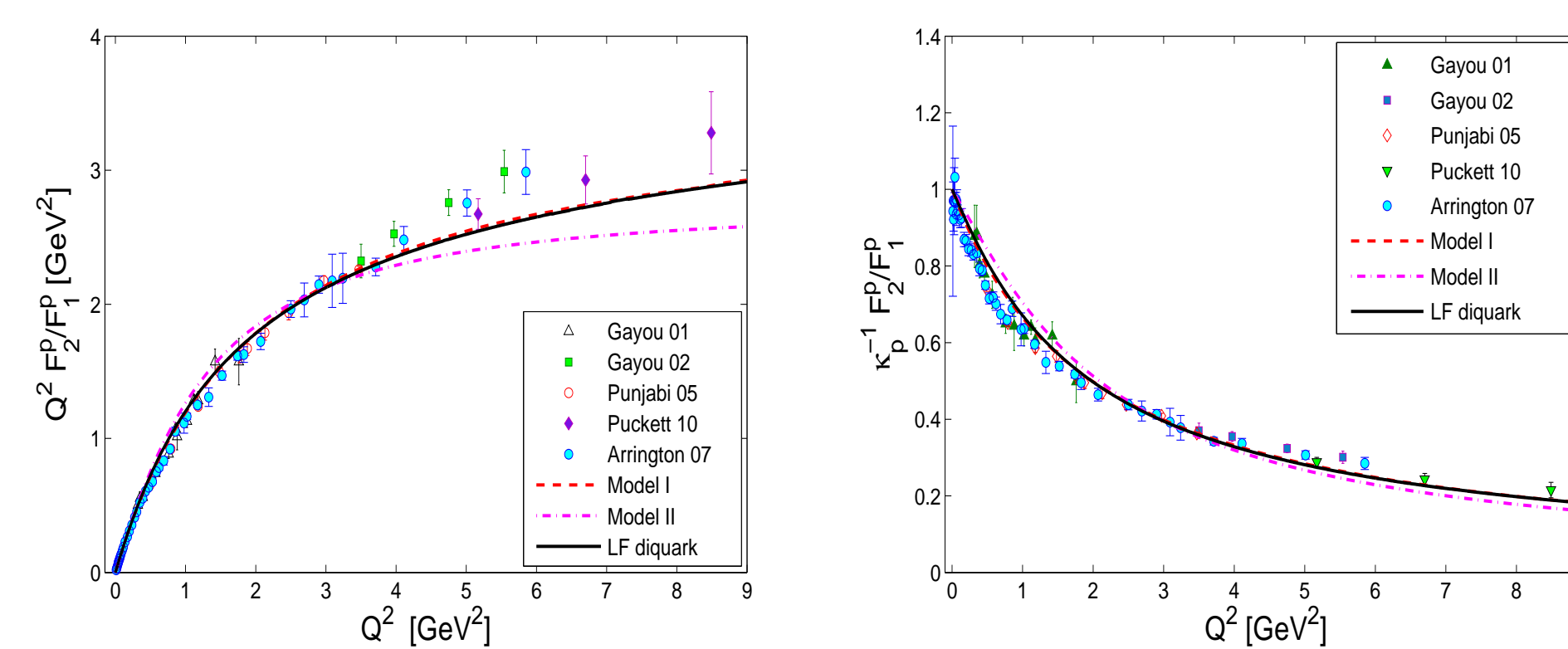


FIGURE 1: Ratio of Pauli and Dirac form factors for proton [4].

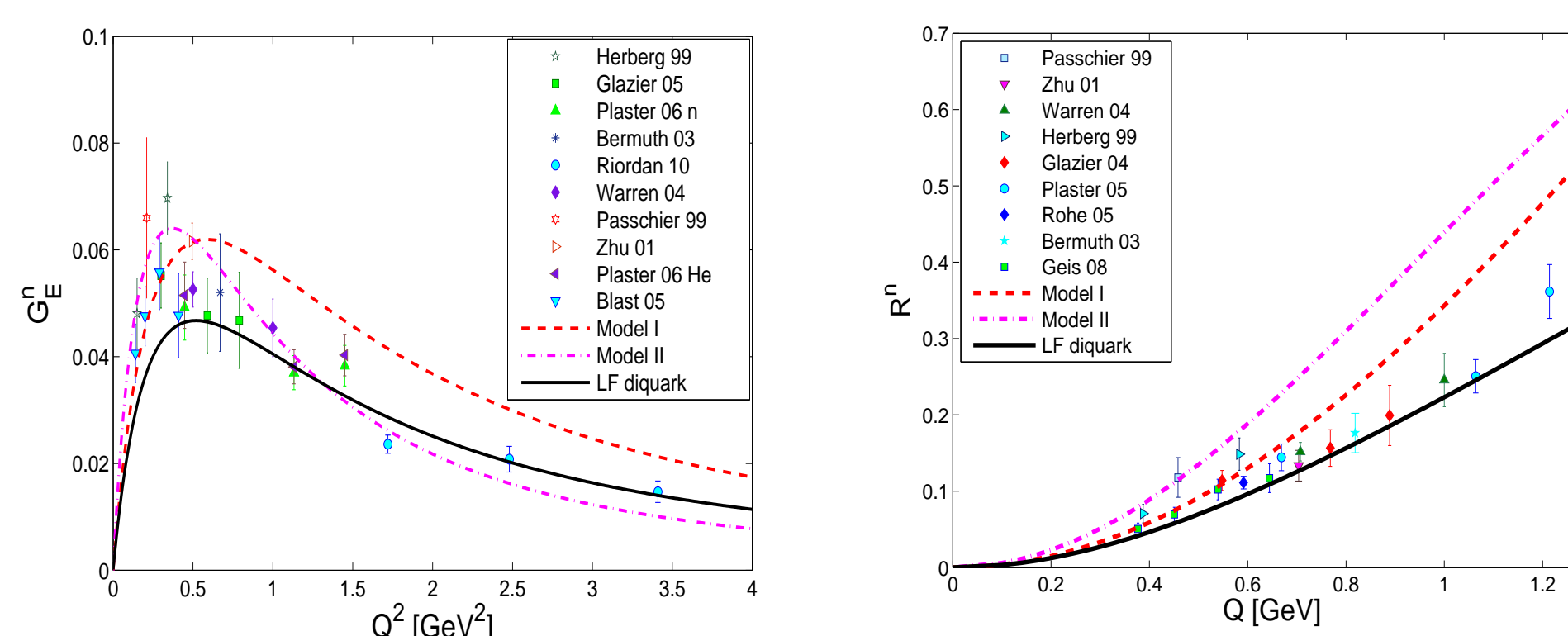


FIGURE 2: Neutron Sachs form factor  $G_E^n(Q^2) = F_1^n(Q^2) - \frac{Q^2}{4M^2} F_2^n(Q^2)$  and the ratio  $R^n = \frac{\mu_n G_E^n}{G_M^n}$  [4].

## Flavor decompositions

- Under the charge and isospin symmetry, the flavor decompositions of the nucleon FFs [Cates *et al.* PRL 106, 252003 (2011)]

$$F_i^p = e_u F_i^u + e_d F_i^d, \quad F_i^n = e_u F_i^d + e_d F_i^u$$

- Normalizations :  $F_1^u(0) = 2$ ,  $F_1^d(0) = 1$  and  $F_2^u(0) = \kappa_u$ ,  $F_2^d(0) = \kappa_d$ .
- Anomalous magnetic moments :  $\kappa_u = 2\kappa_p + \kappa_n = 1.673$  and  $\kappa_d = \kappa_p + 2\kappa_n = -2.033$ .

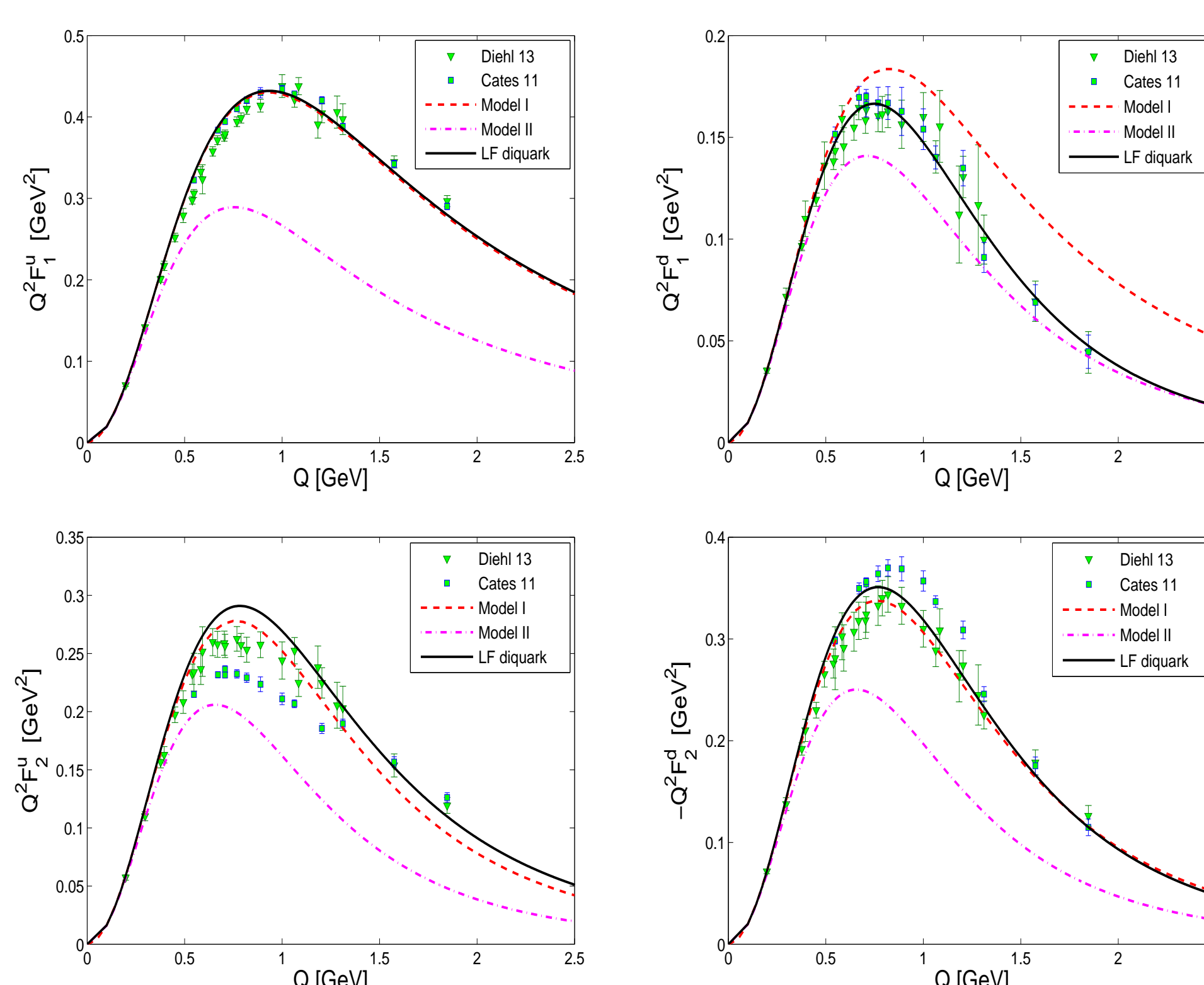


FIGURE 3: Plots of flavor dependent form factors for  $u$  and  $d$  quarks [4].

## Transverse densities

- The transverse charge density inside the nucleons :

$$\rho_{ch}(b) = \int \frac{d^2q_\perp}{(2\pi)^2} F_1(q^2) e^{iq_\perp \cdot b_\perp} = \int_0^\infty \frac{dQ}{2\pi} Q J_0(Qb) F_1(Q^2).$$

- The anomalous magnetization density :

$$\rho_m(b) = -b \frac{\partial \bar{\rho}_M(b)}{\partial b} = b \int_0^\infty \frac{dQ}{2\pi} Q^2 J_1(Qb) F_2(Q^2).$$

- For transversely polarized nucleon with transverse spin :  $\vec{S}_\perp = \cos \phi_s \hat{x} + \sin \phi_s \hat{y}$ , the charge density is given by

$$\rho_T(b) = \rho_{ch}(b) - \sin(\phi_b - \phi_s) \frac{1}{2Mb} \rho_m(b).$$

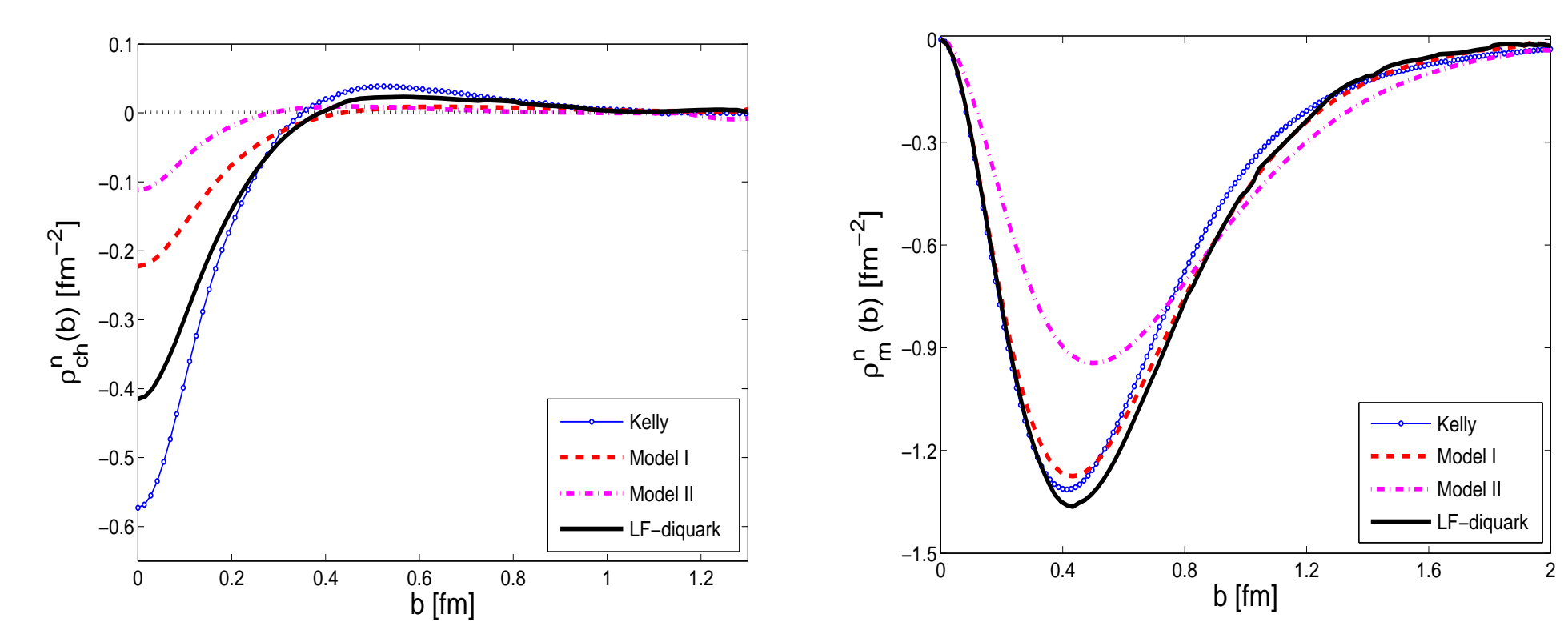
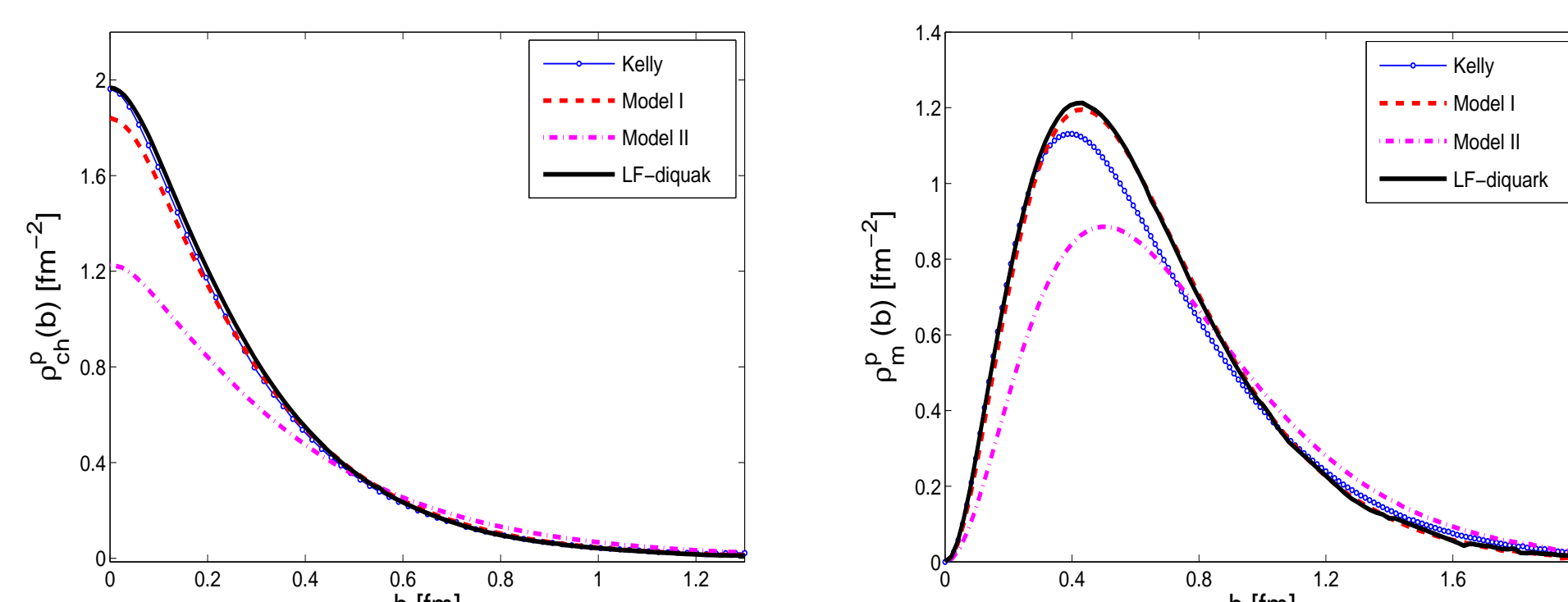


FIGURE 4: Plots of charge and anomalous magnetization densities for proton and neutron in transverse plane [4].

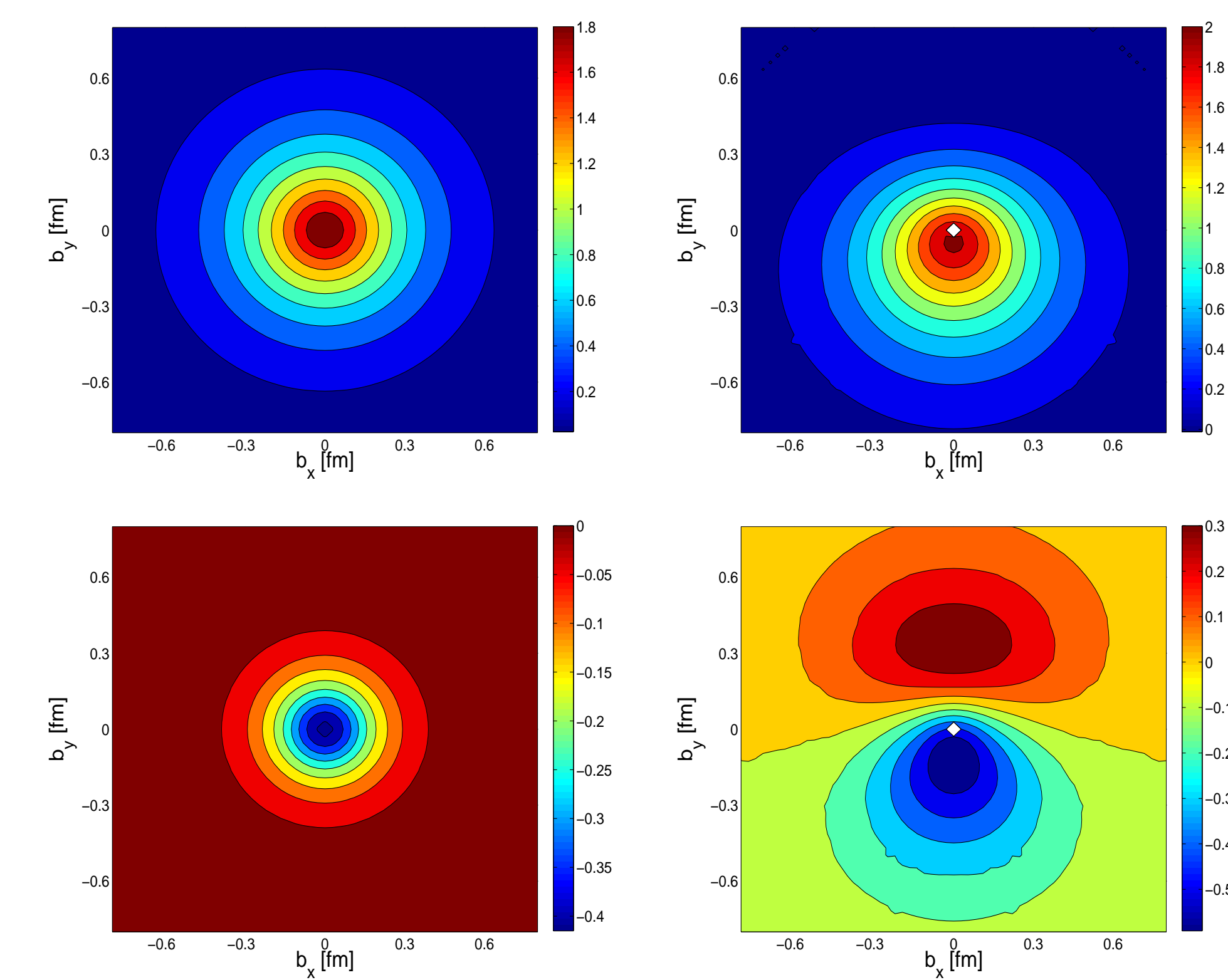


FIGURE 5: Unpolarized and transversely polarized charge densities for proton (upper) and neutron (lower) in LF diquark model [4].

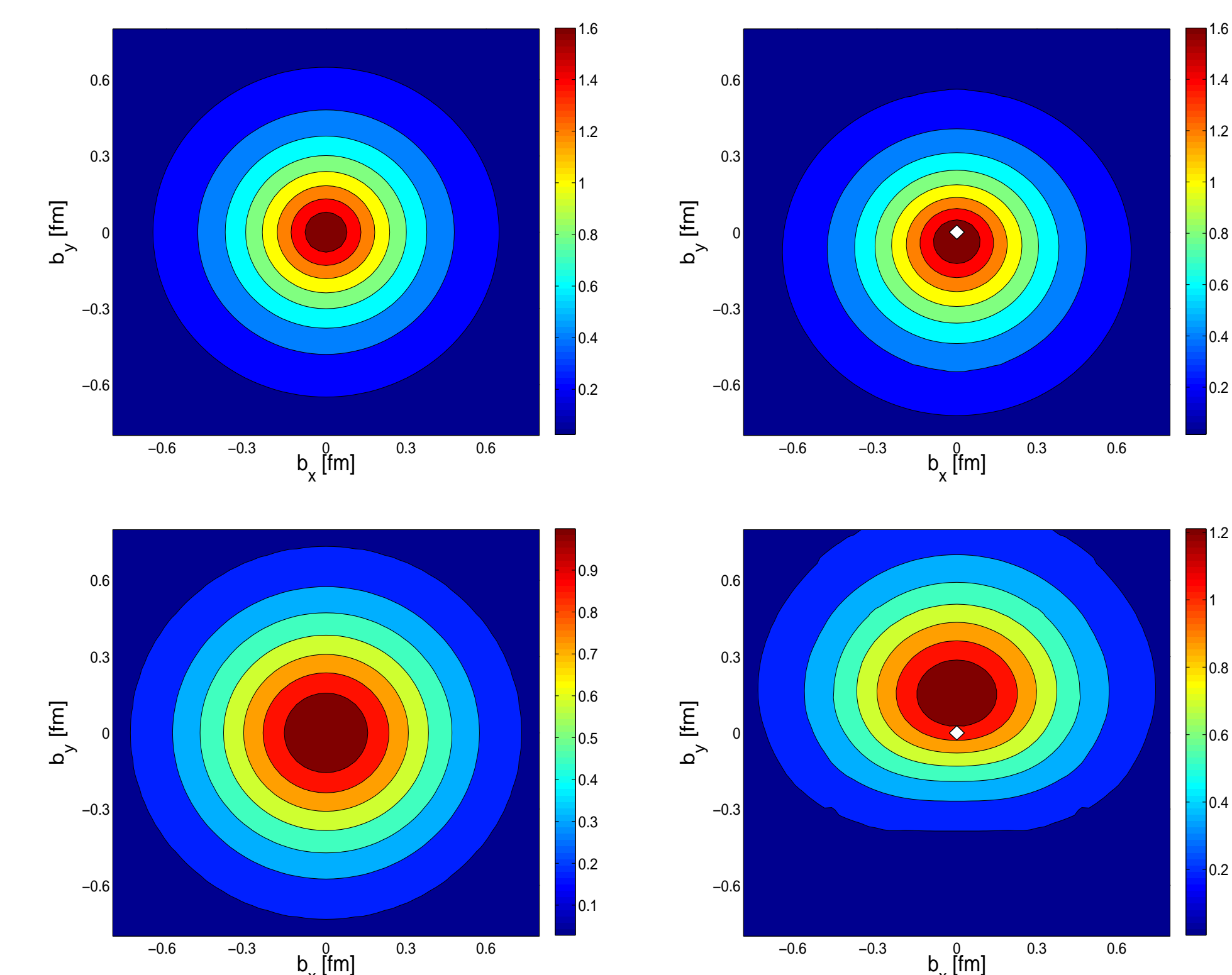


FIGURE 6:  $u$  (upper) and  $d$  (lower) quarks charge densities in a unpolarized and transversely polarized nucleon in LF diquark model [4].

## Electromagnetic radii

$$\langle r_E^2 \rangle^N = -6 \frac{dG_E^N(Q^2)}{dQ^2} \Big|_{Q^2=0},$$

$$\langle r_M^2 \rangle^N = -\frac{6}{G_M^N(0)} \frac{dG_M^N(Q^2)}{dQ^2} \Big|_{Q^2=0}$$

The Sachs form factors are defined as

$$G_E^N(Q^2) = F_1^N(Q^2) - \frac{Q^2}{4M_N^2} F_2^N(Q^2),$$

$$G_M^N(Q^2) = F_1^N(Q^2) + F_2^N(Q^2).$$

Quantity	Model I	Model II	LF diquark	Measured data
$r_E^p$ (fm)	0.810	0.980	0.786	$0.877 \pm 0.005$
$r_M^p$ (fm)	0.782	0.921	0.772	$0.777 \pm 0.016$
$\langle r_E^2 \rangle^n$ (fm <sup>2</sup> )	-0.088	-0.123	-0.085	$-0.1161 \pm 0.0022$
$r_M^n$ (fm)	0.796	0.937	0.7596	$0.862^{+0.009}_{-0.008}$

## References

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