

22/12/2015

The ATLAS Muon System

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Summary

- Overall design
- Track reconstruction
- Performance measurements
- Trigger
- Outlook

“specifications”

“Physics Requirements” from the Technical Design Report (TDR) :

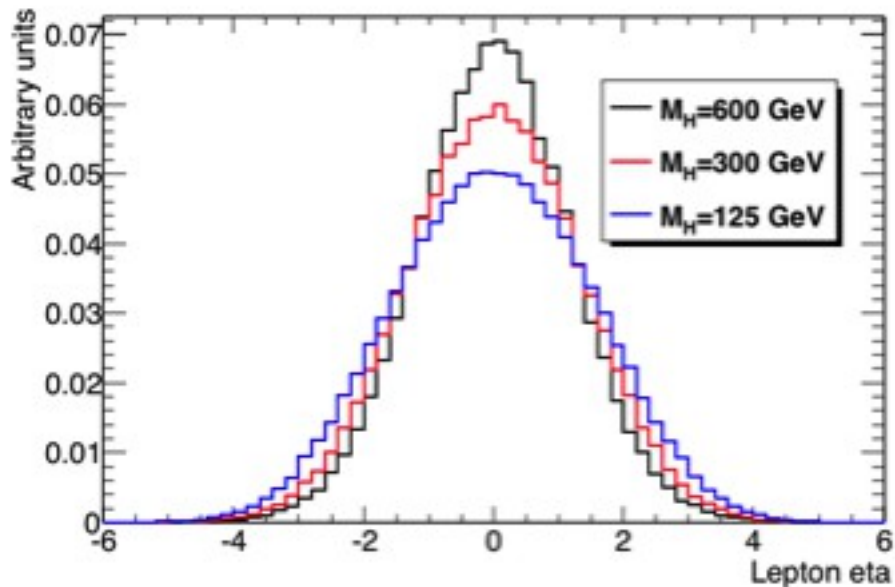
- Identify and reconstruct muon tracks, measure their momenta, and provide matching information for association with inner-detector data [...].
- Trigger on single- or multi-muon event topologies [...].
- Unambiguously associate the muon with its parent bunch crossing.

The scale of the performance requirements is set by a number of benchmark reactions:

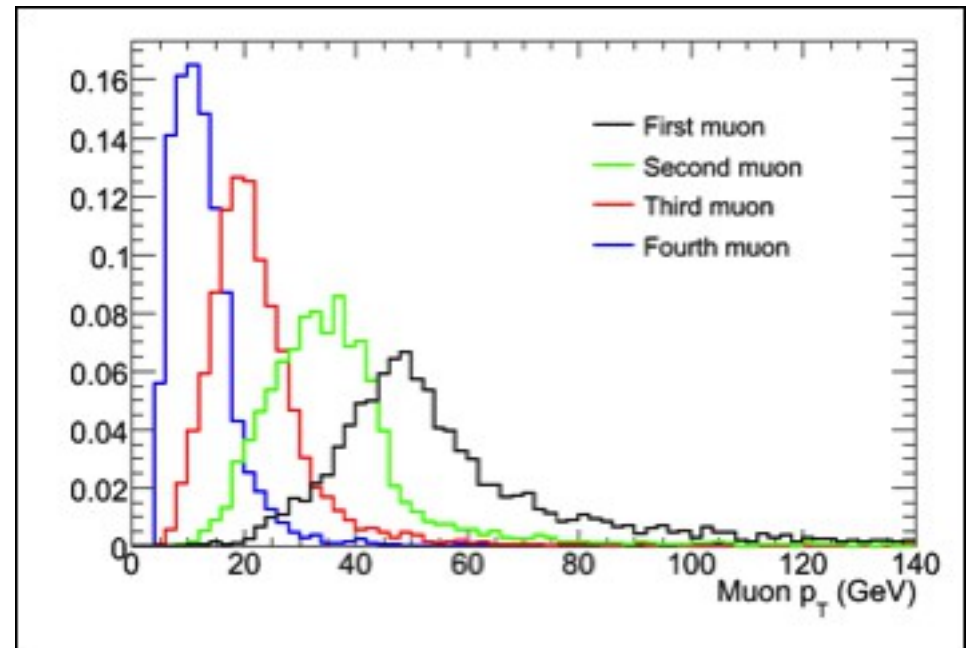
- 1) $H \rightarrow ZZ^* \rightarrow \mu\mu ll$ SM ($120 < m_H < 170$ GeV) [...]
- 2) $H \rightarrow ZZ^* \rightarrow \mu\mu ll, A \rightarrow \mu\mu$ MSSM ($180 < m_H < 2m_{top}$) [...]
- 3) New vector bosons $Z' \rightarrow \mu\mu, W' \rightarrow \mu\nu$ ($1 < m < 5$ TeV);
- 4) B-physics

Example: muons from $H \rightarrow ZZ^* \rightarrow 4\mu$

Pseudo-rapidity η :



Transverse momentum p_T



In practice for Higgs analysis we would like a detector with

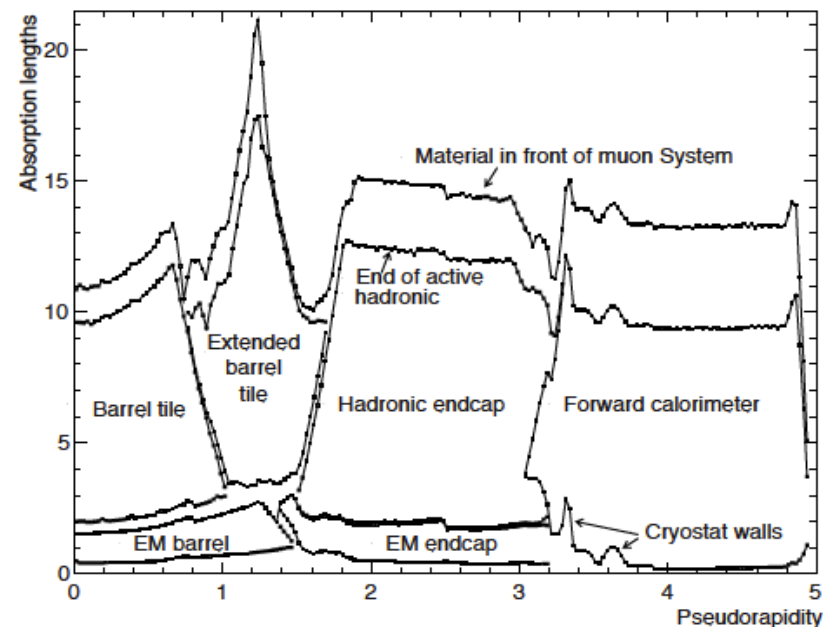
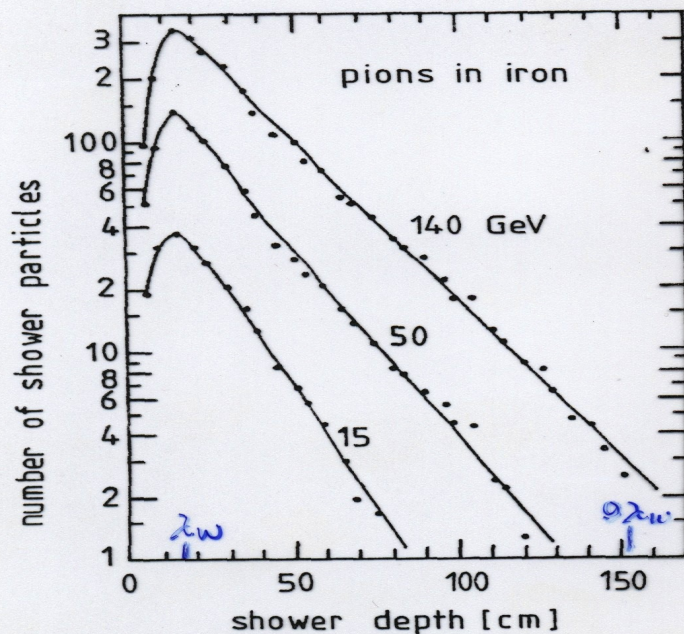
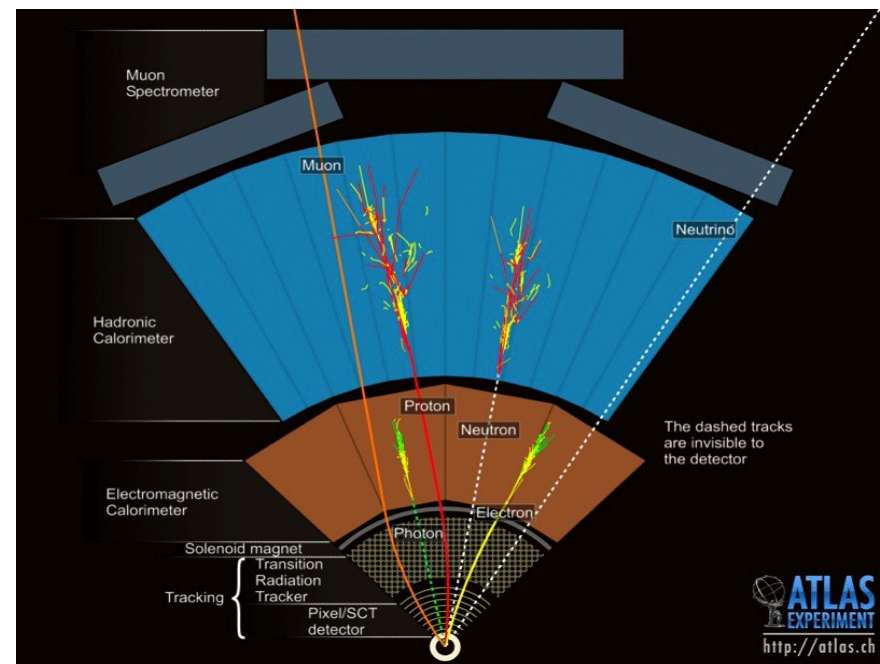
- coverage $|\eta| < \sim 3$
- Reco from $p_T > \sim 5$ GeV
- Trigger: 1 mu with $p_T > \sim 25$ GeV or 2 mu with $p_T > \sim 10$ GeV
- Best possible momentum resolution

The actual choices are driven by cost-performance optimization

Parameter	Main physics criteria	Performance		Comments
		desired	actual	
Momentum measurement				
$\Delta p_T/p_T$ at 20 GeV	$H \rightarrow ZZ^* \rightarrow 4l$	1-2%	~2.5%	Muon spectrometer only limited by energy loss and multiple scattering
			~1.6%	Combined with inner tracker
$\Delta p_T/p_T$ at 75 GeV	$H \rightarrow ZZ \rightarrow 4l$ (MSSM)	1-2%	~2.4%	Muon spectrometer only limited by energy loss and multiple scattering
			~2.0%	Combined with inner tracker
$\Delta p_T/p_T$ at 1000 GeV	$Z' \rightarrow \mu\mu$	few %	~11%	Resolution limited by cost-performance optimization; charge determination is the driving criterion
Rapidity coverage	Above processes	~3	2.7	Limited by system integration and shielding
Trigger				
Low- p_T threshold	b physics and CP violation (event rate)	~ 5 GeV	6 GeV	Limited by muon energy loss for triggering behind the calorimeter and by hadron decays in flight
High- p_T threshold	$H \rightarrow ZZ^* \rightarrow 4\mu$ (event rate)	20 GeV	20 GeV	Background-dependent, tunable
Rapidity coverage	$b \rightarrow \mu x$ (event rate)	~ 2.7	2.4	Limited by low- p_T rate and accidentals
	$H \rightarrow ZZ^* \rightarrow 4\mu$	~2.5	2.4	Single-muon trigger at high p_T provides good trigger efficiency
	$Z' \rightarrow \mu\mu$ (asymmetry)	~2.0	2.4	Single-muon trigger at high p_T provides good trigger efficiency
Bunch crossing identification	Event matching	$\sigma < 5$ ns	$\sigma < 4$ ns	

Muon identification

- Muon identification is based on the absorption of other particles producing EM and Hadronic showers in the calorimeters
- In ATLAS $> \sim 10$ interaction lengths provide shower containment [$< 95\%$ of energy] for pions up to approx 100 GeV
- Simulations provide the number and momentum of charge particles “leaking” from showers

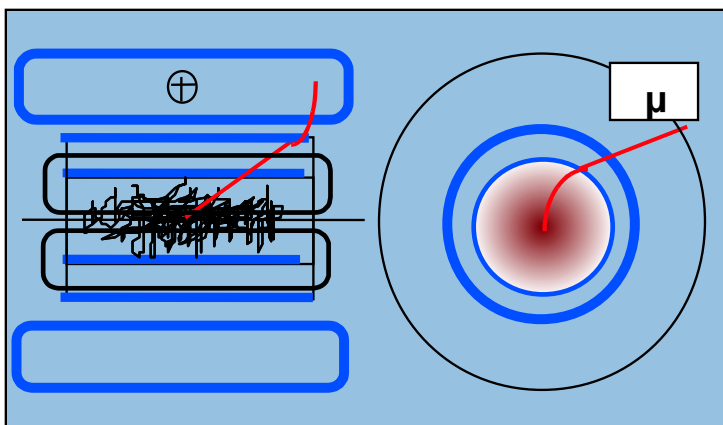
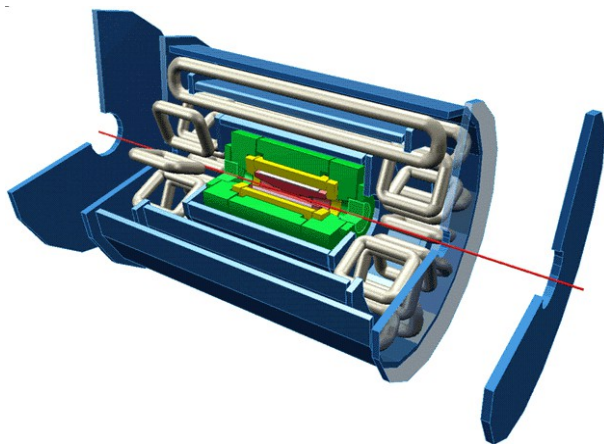


Magnetic field configuration: two different choices

ATLAS:

- thin solenoid inside EM CALI, $B \sim 2\text{T}$
- muon system in large air-core toroidal field
- “smaller” inner tracker
- precise stand-alone muon momentum measurement in the MS

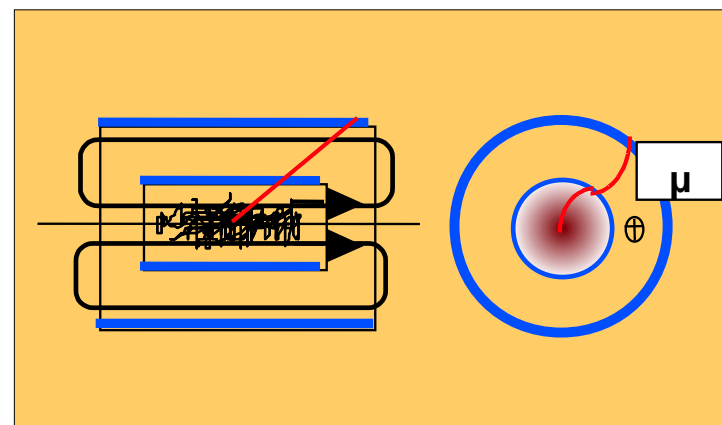
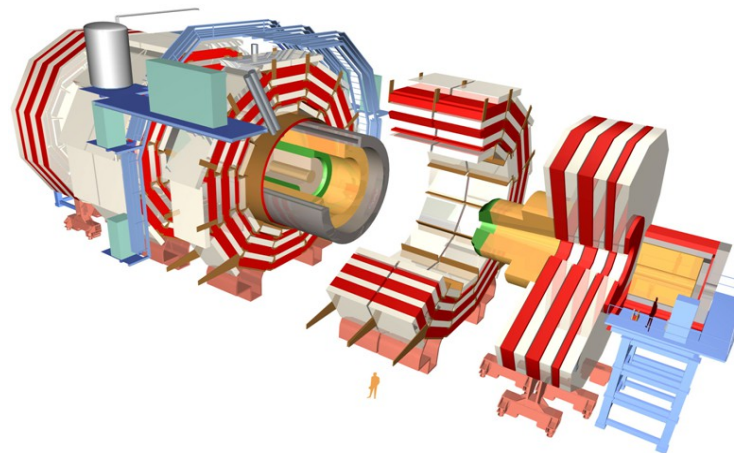
ATLAS A Toroidal LHC Apparatus



CMS:

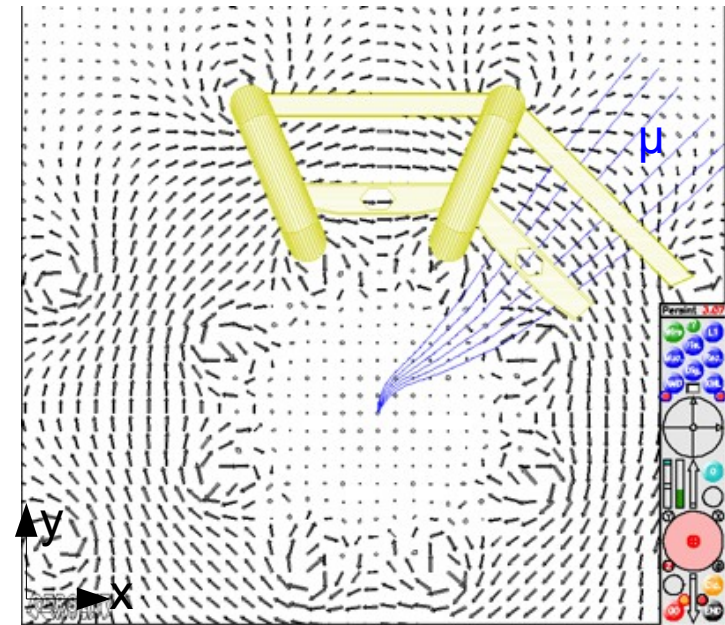
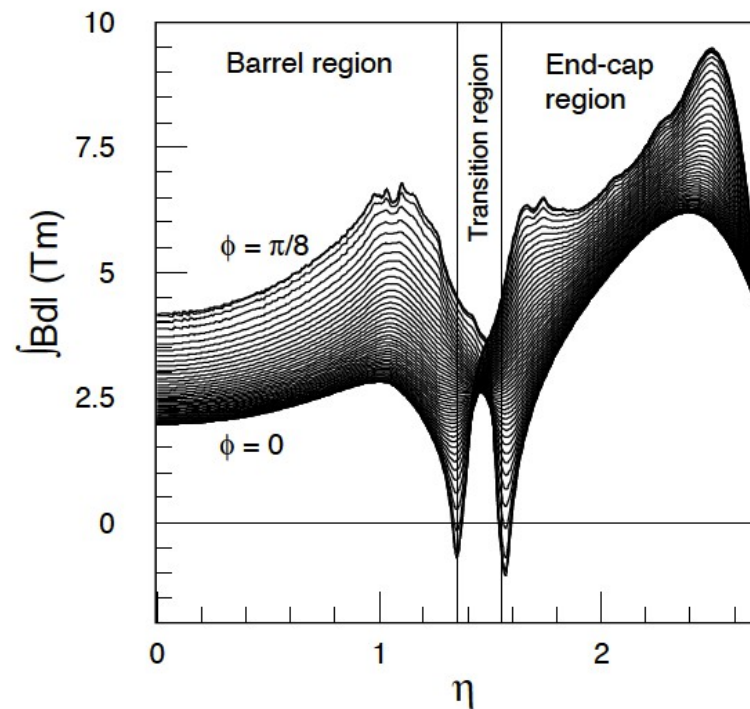
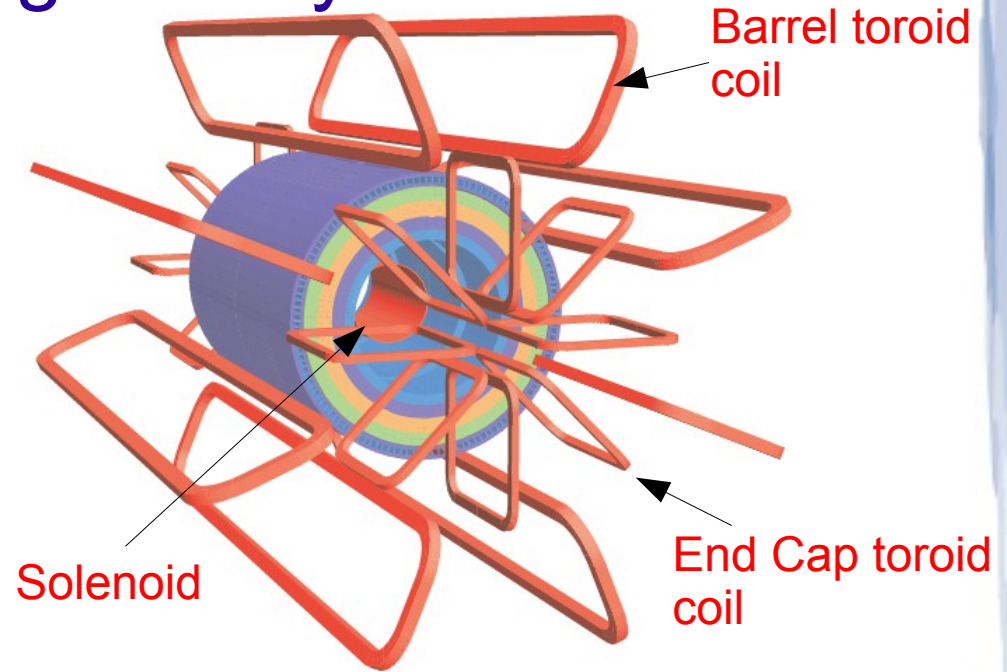
- large solenoid outside calorimeters, $B \sim 4\text{T}$
- muon system in the iron yoke for magnetic field return
- large inner tracker

CMS Compact Muon Solenoid

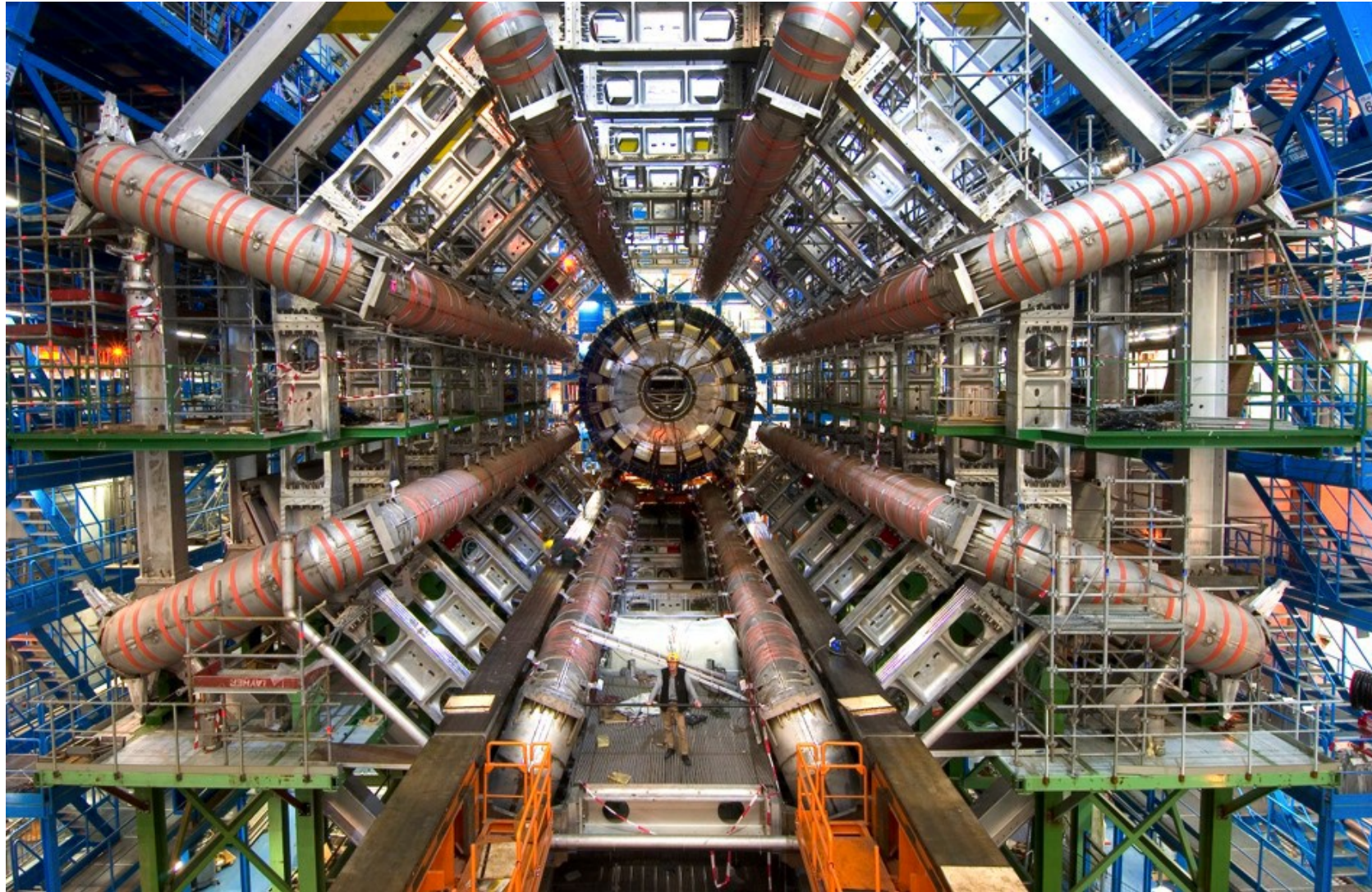


The ATLAS magnetic system

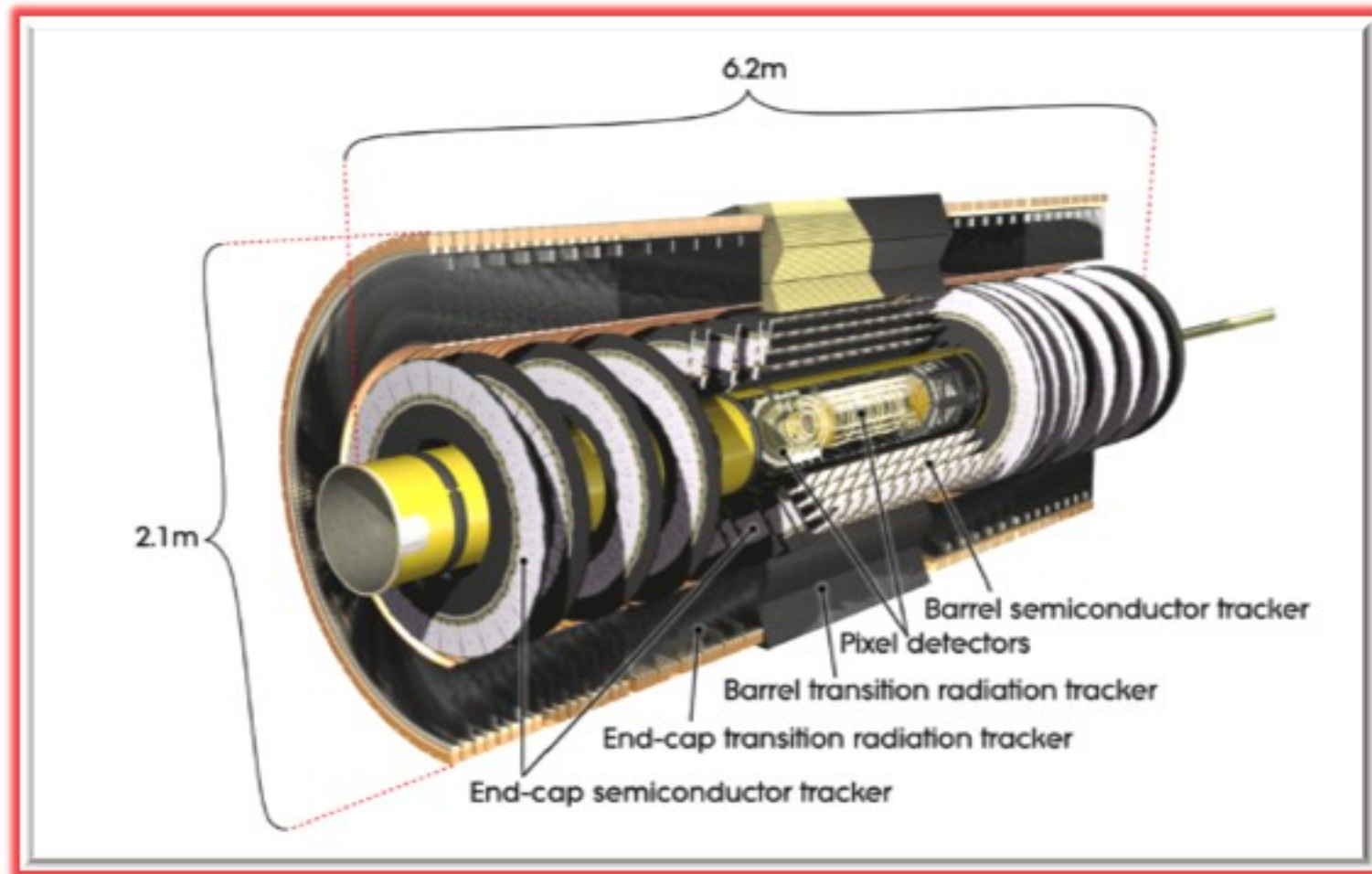
- Central barrel solenoid $B=2T$, $R=1$ m
- Barrel toroid (8 coils)
- Two End-Cap toroids (8 coils)
- MS: bending in η , straight tracks in ϕ
- Complex field configuration due to “few” coils and Barrel/End Cap transition
- Field integral seen by a muon in MS:
 $\int B dl = 2.5 : 10$ Tm



The Barrel toroid coils during construction



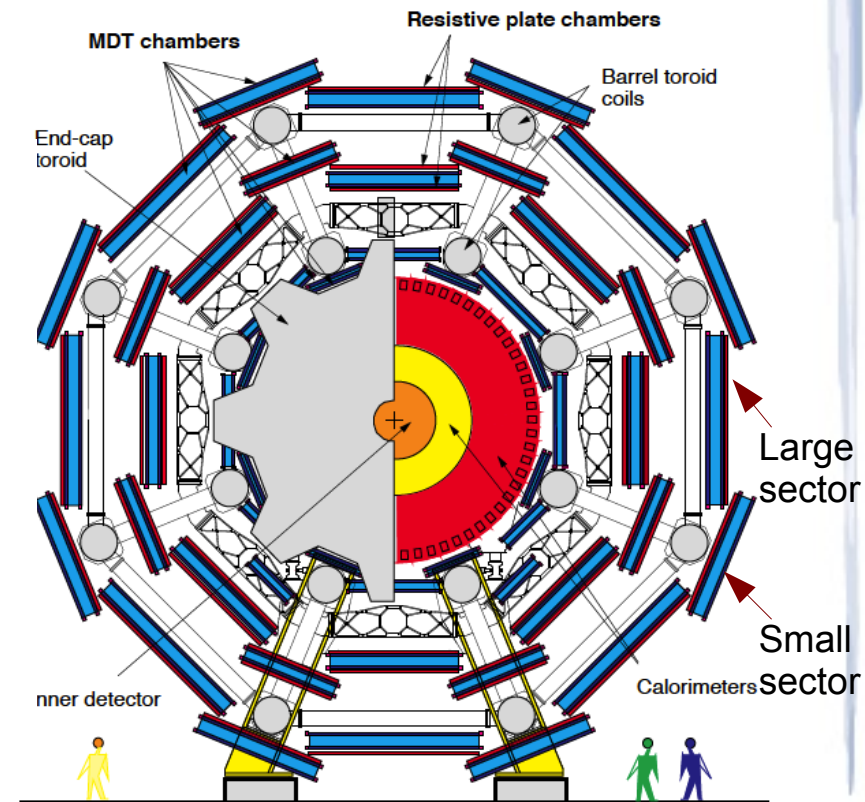
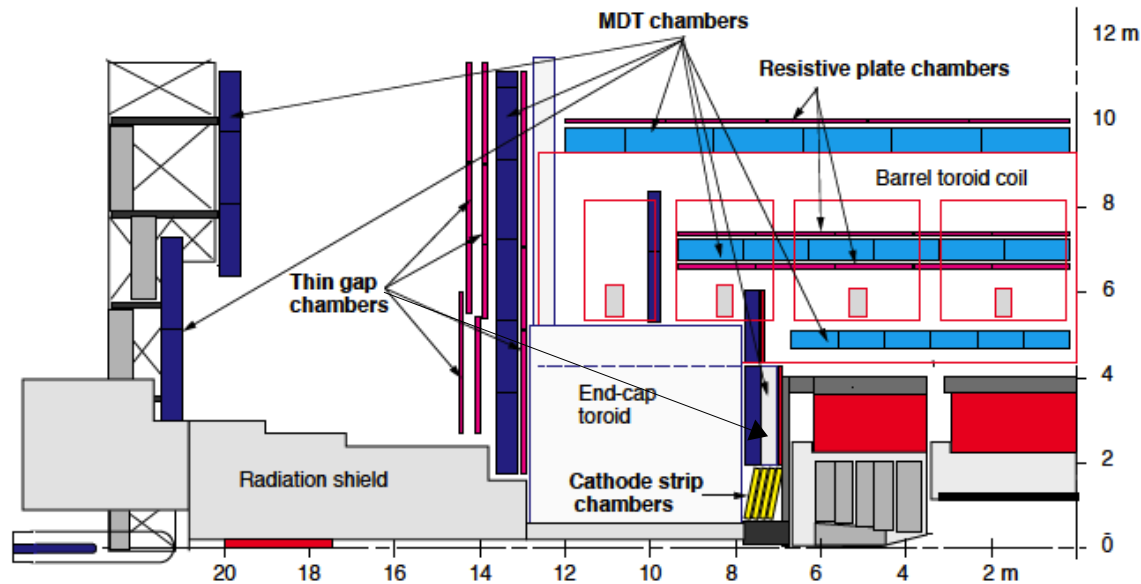
The Inner Detector (ID)



- **Pixel**: 3 layers 2D - Hit resolution bending plane: $10 \mu\text{m}$ (NOW 4 Layers with the IBL)
- **Silicon Strips**: (SCT) 4-9 layers 2D (stereo strips) - Hit resol bending plane: $17 \mu\text{m}$
- **Straw Tubes** (TRT): up to 160 planes -1D Hit resol bending plane: $130 \mu\text{m}$
- Magnetic Field : Solenoid 2 T, Angular Coverage $\eta < 2.5$

The muon system (MS)

- Three layers of “precision” chambers for precise measurement in the bending plane
 - MDT (monitored drift tubes)
 - CSC (cathode strip chambers) inner layer $|\eta| > 2$
- 3(4) layers of “trigger” chambers for triggering and ϕ coordinate
 - RPC (resistive plate chambers) in barrel
 - TGC (Thin gap chambers) in endcaps



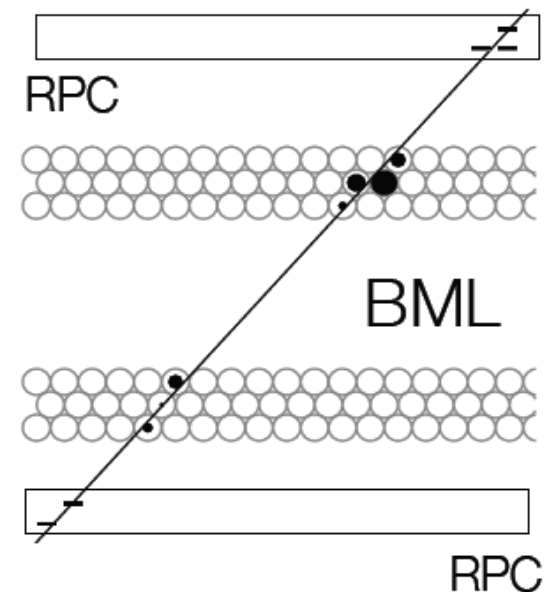
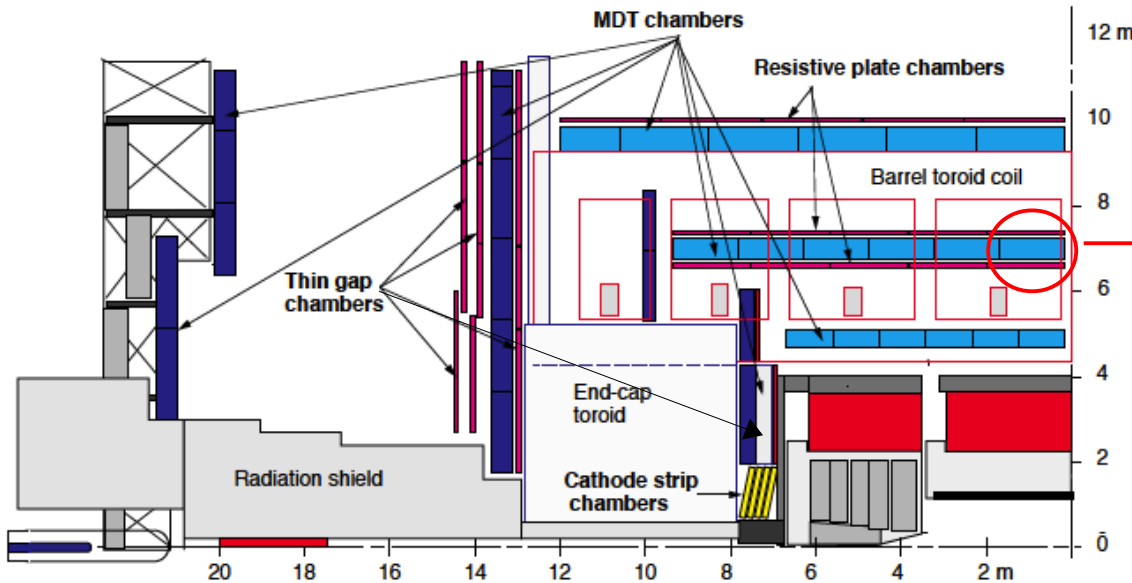
The muon system (MS)

- Three layers of “precision” chambers for precise measurement in the bending plane
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 - TGC (Thin gap chambers) in endcaps

Total hits along track:
~ 20 precision hits
~ 6 (barrel) to 12 (endcap) trigger hits

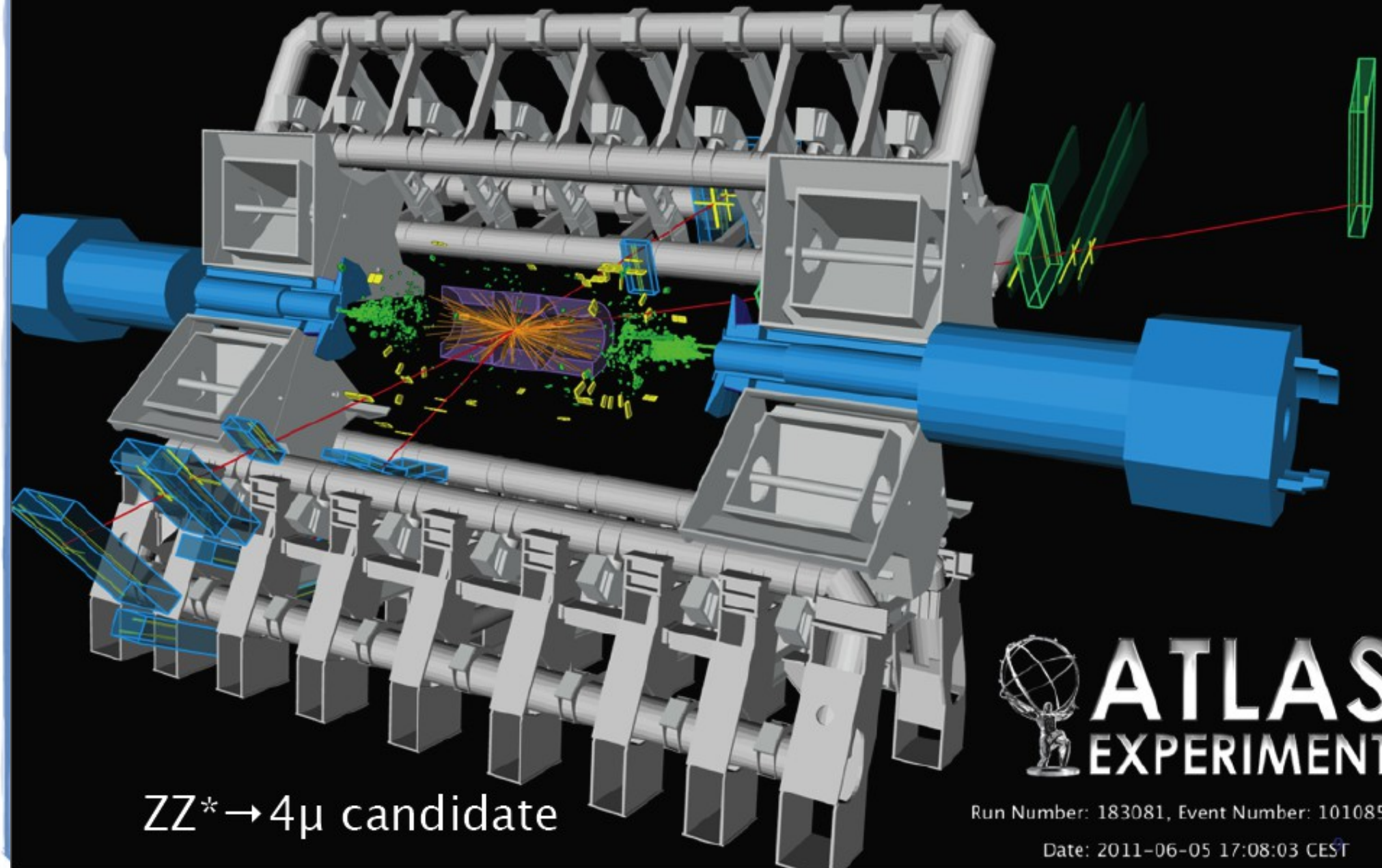
Example:

“Barrel Middle Large” station:
3+3 precision and 2+2 trigger points



In Reality ?

... a bit more complicated



$ZZ^* \rightarrow 4\mu$ candidate



ATLAS EXPERIMENT

Run Number: 183081, Event Number: 10108572

Date: 2011-06-05 17:08:03 CEST

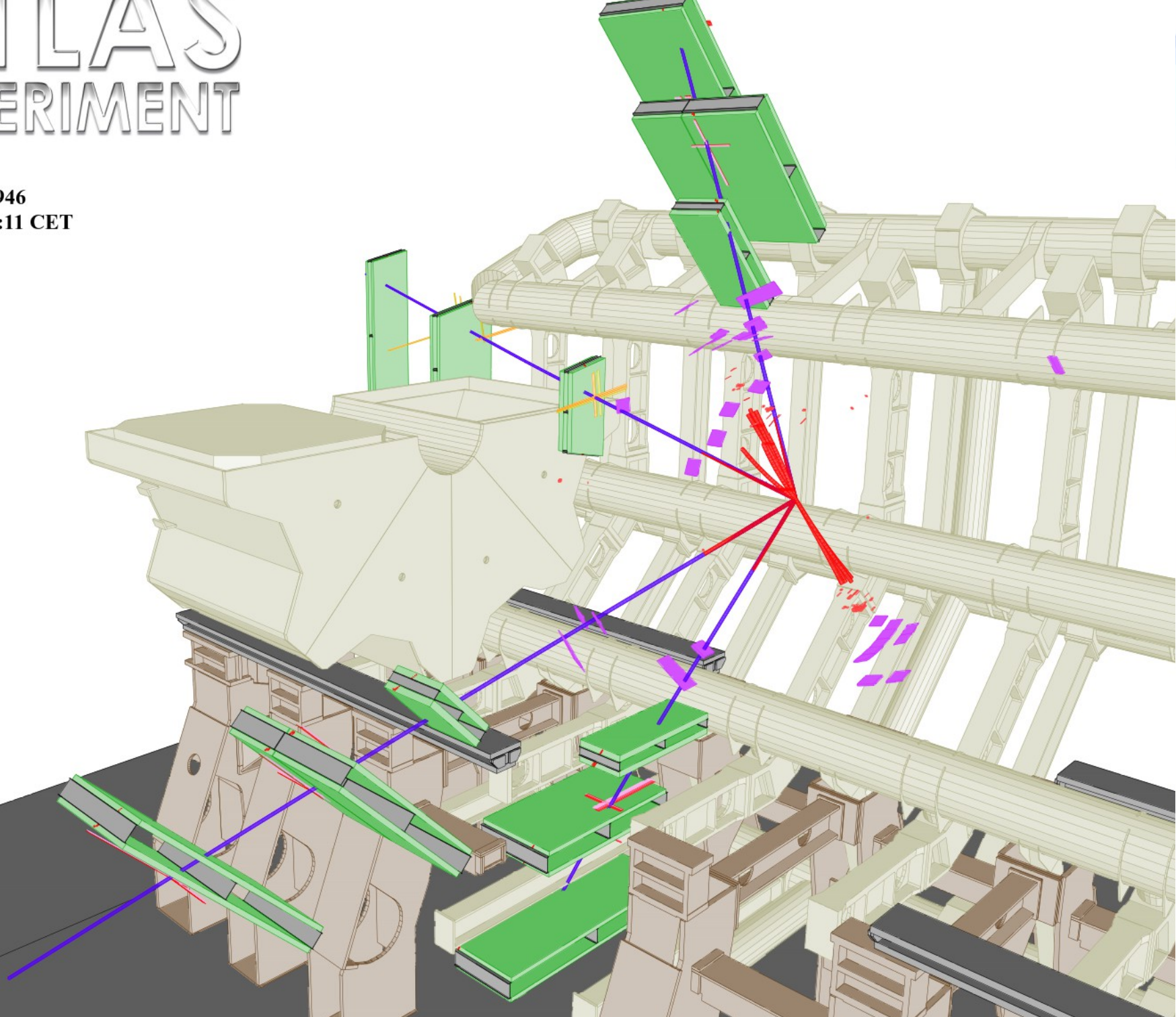


ATLAS EXPERIMENT

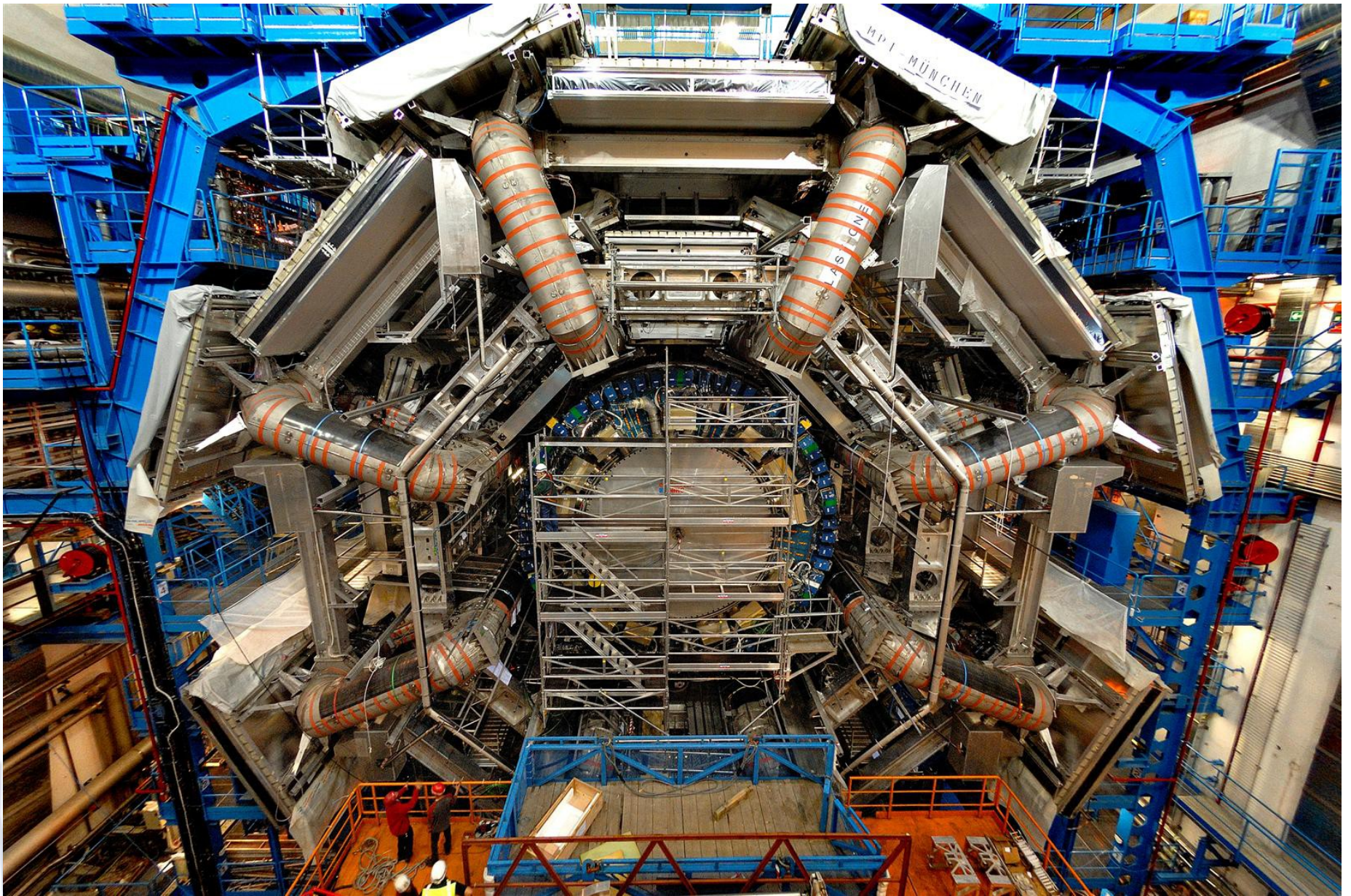
Run Number: 189280,
Event Number: 143576946
Date: 2011-09-14, 11:37:11 CET

EtCut > 0.3 GeV
PtCut > 3.0 GeV
Vertex Cuts:
Z direction < 1cm
Rphi < 1cm

Muon: blue
Cells: Tiles, EMC



Persint



MDTs

Drift tubes:

- $d=30$ mm, wire $d=50$ μm
- $P=3$ bar (abs)
- Ar-N₂-CH₄ (91%/4%/5%)
- HV=3270 V

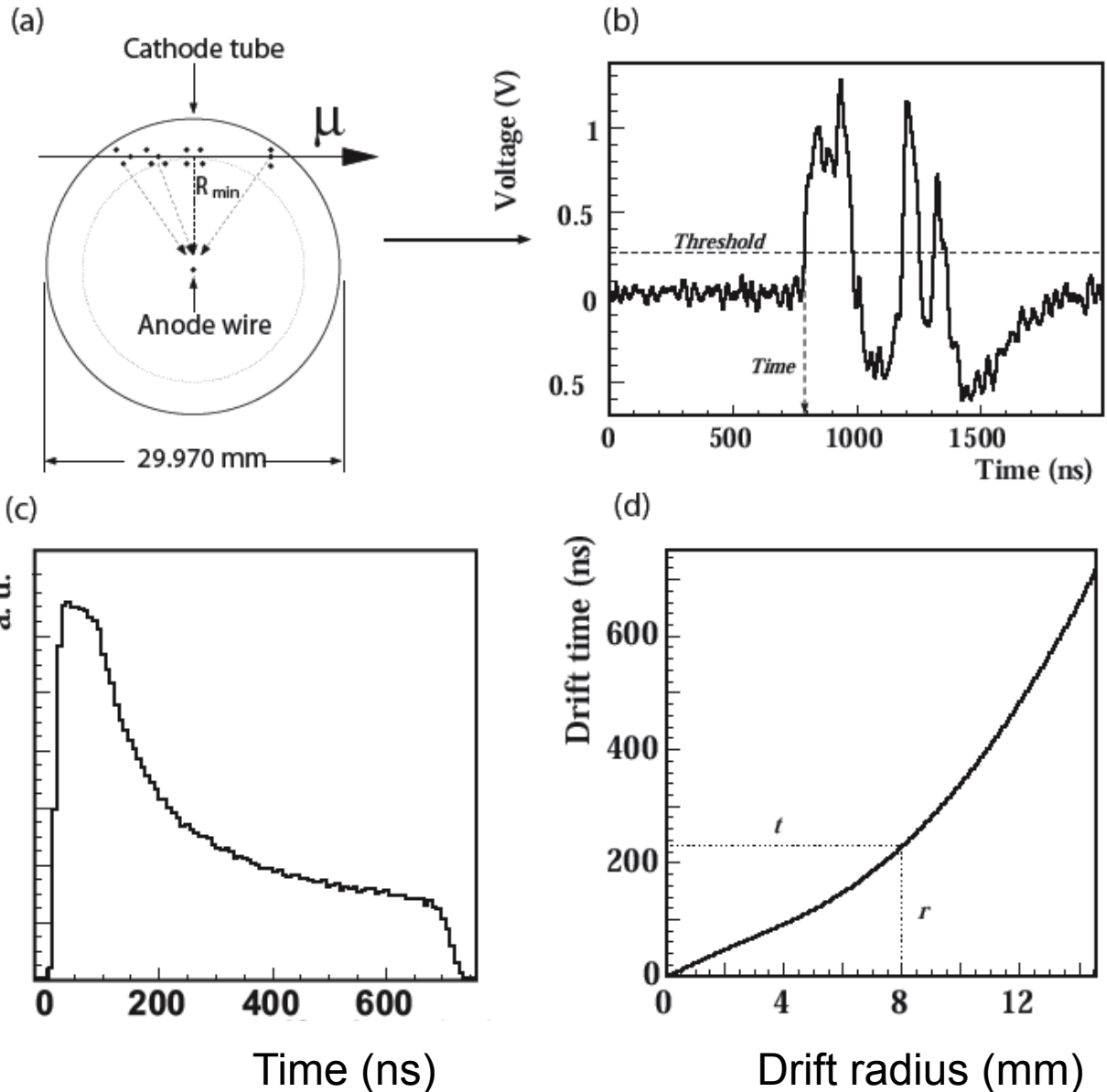
The time of first electron gives is converted into a “drift radius” using the known $r(t)$ relation.

Max drift time ~ 700 ns

Space resolution ~ 80 μm (*)

NB:

* knowing the “start” time, and the position of the muon along the tube



Trigger chambers example: RPCs

Operating Conditions

($E_{\text{gas}} \sim 5 \text{ KV/mm}$)

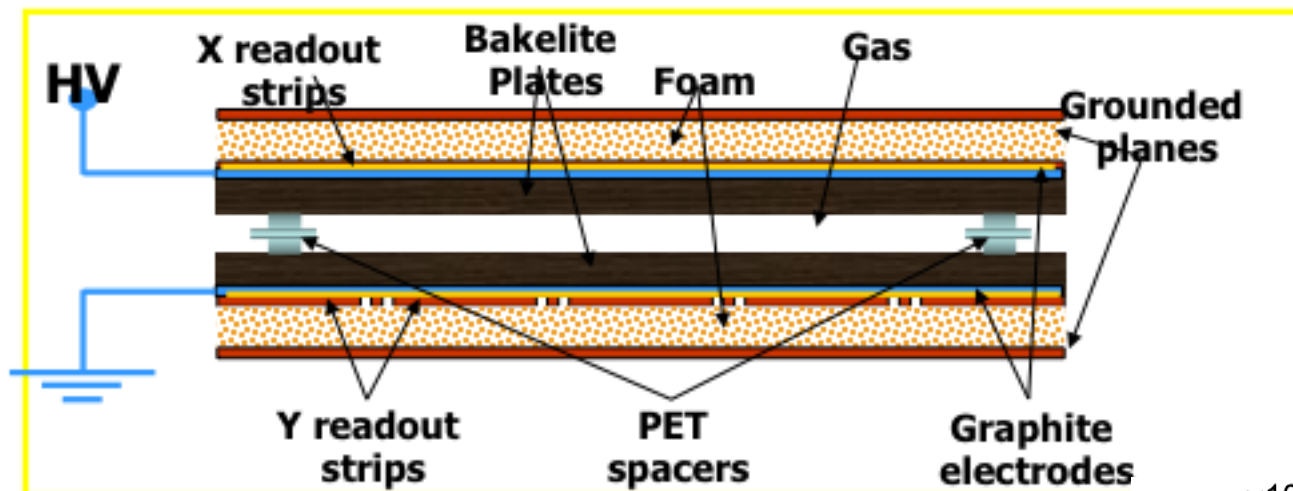
Gas: $\text{C}_2\text{H}_2\text{F}_4$ 95% - C_4H_{10} 4.5% - SF_6 0.5% ;

$\rho_{\text{bakelite}} \sim 2 \times 10^{10} \Omega\text{cm}$;

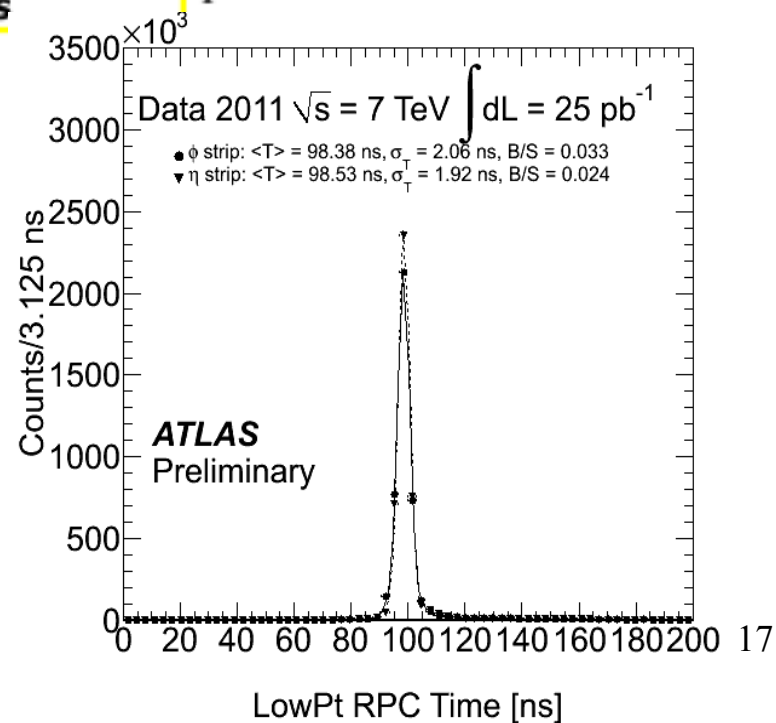
Gas Gap $d = 2 \text{ mm}$;

Graphite coated HV electrodes

Cu read out strips 30 mm pitch

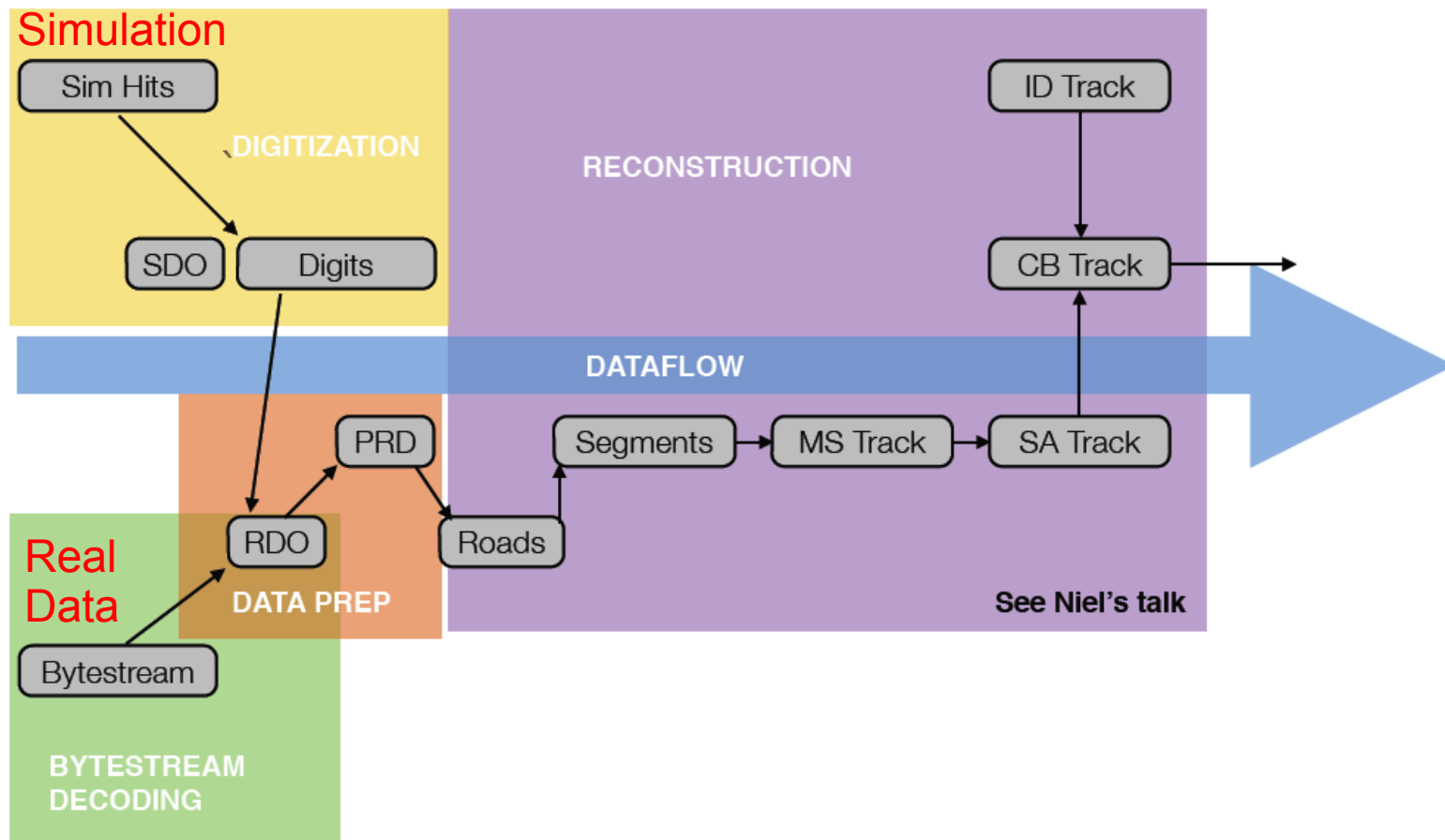


- Readout strips in η and ϕ
- pitch $\sim 3 \text{ cm}$
- time resolution $\sim 2 \text{ ns}$



Track Reconstruction

- Once hits are produced in the detectors
The ATLAS reconstruction program should
 - identify the muons with high efficiency and purity
 - reconstruct muon parameters (charge, momentum, direction)
- Online version to be run in the trigger should also be fast



Pattern recognition

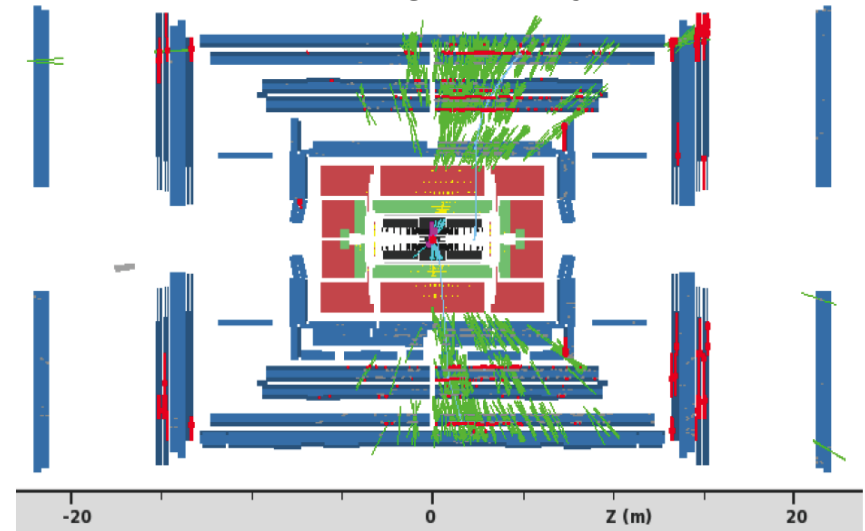
The MS is filled by hits,
not only muons but also :

- Tails of hadronic showers
- Neutron and photons from hadronic int. including a long-lifetime component from slow neutrons (cavern background)
- electronic noise

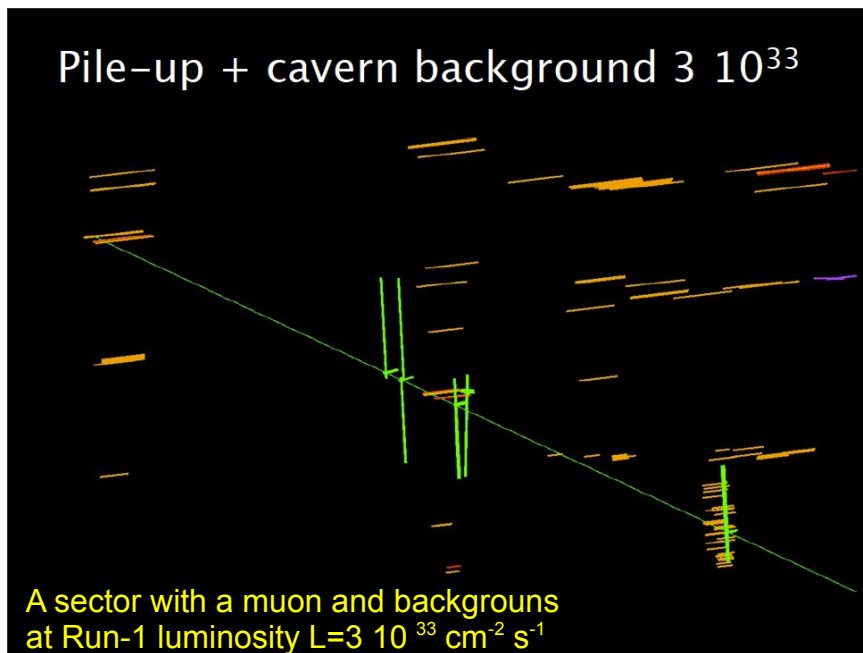
Not possible to try all hit combinations
(CPU time would diverge)

Need to find patterns from charged tracks

High-Et Dijet event



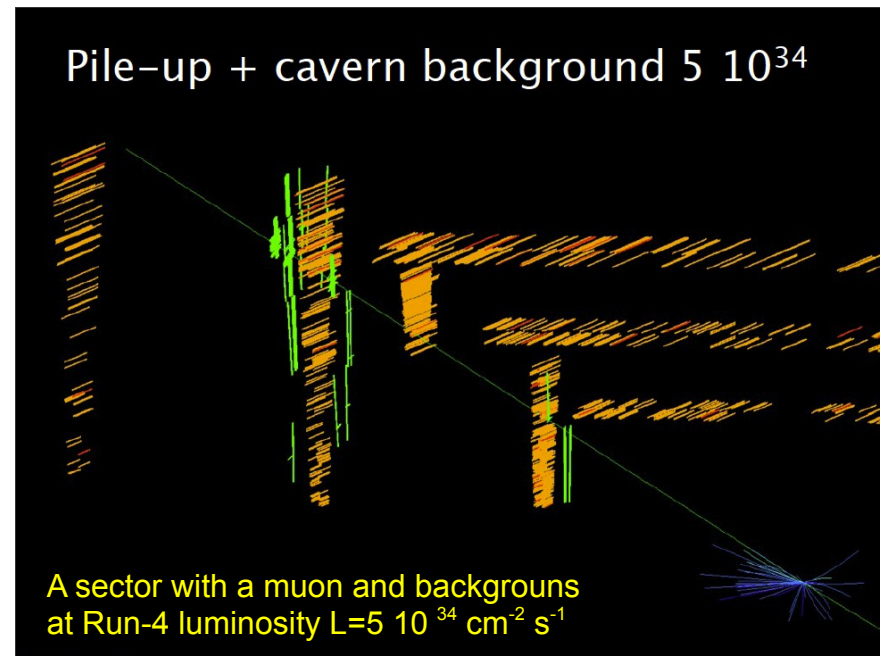
Pile-up + cavern background $3 \cdot 10^{33}$



A sector with a muon and backgrounds
at Run-1 luminosity $L=3 \cdot 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$

Sunday, November 10, 13

Pile-up + cavern background $5 \cdot 10^{34}$



A sector with a muon and backgrounds
at Run-4 luminosity $L=5 \cdot 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$

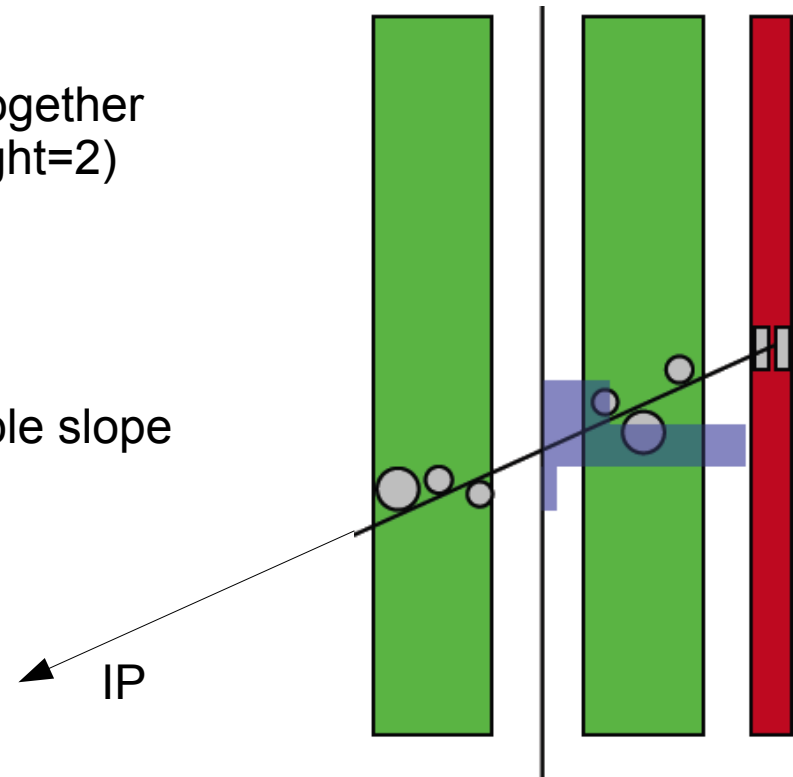
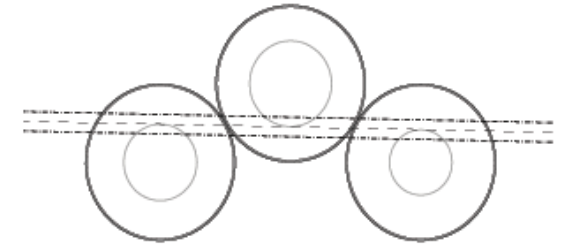
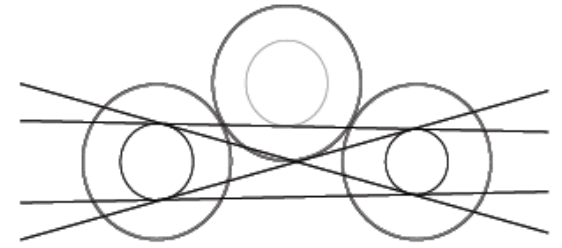
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Pattern recognition

3 tubes MDT multilayer :

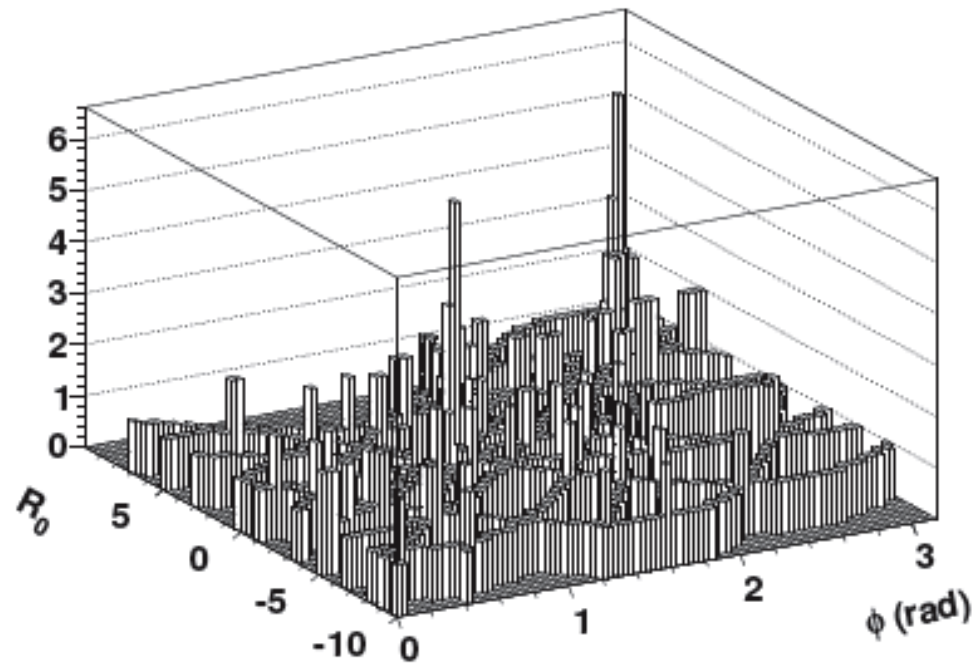
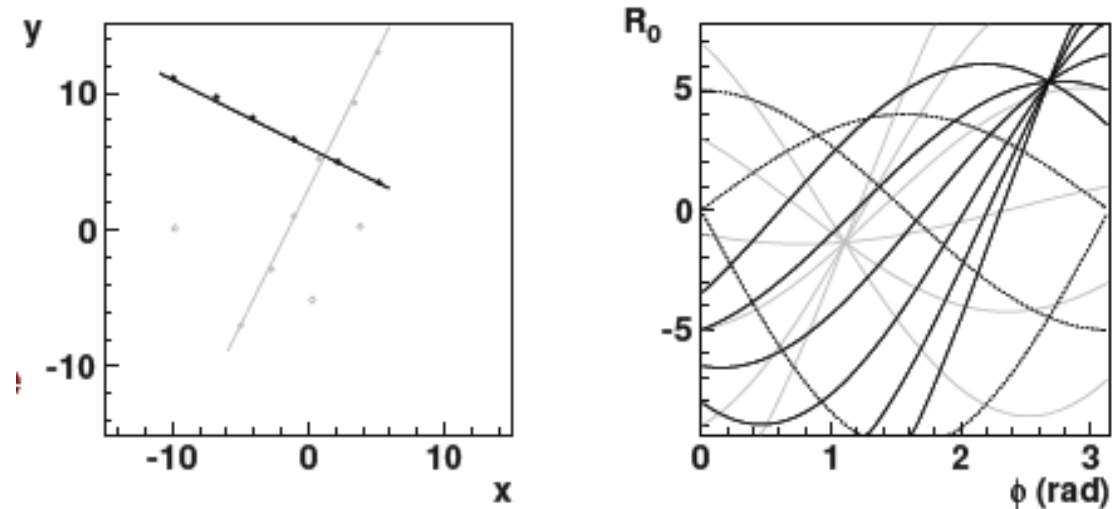
- combining tangents to drift circles gives many possibilities
- Need to consider that tubes are very efficient and accept only the possibilities with more hits
- Simple histogramming technique:
 - project hits along the direction pointing to the interaction point (IP)
 - Select cases with $N_{\text{hits}} \geq 5$
 - Can use trigger and precision chambers together with different weights (e.g. trigger hits weight=2)

=> very fast, linear with num of hits
=> works only for straight tracks from IP
- Extension :
do different histograms, one for each possible slope
=> Hough Transform



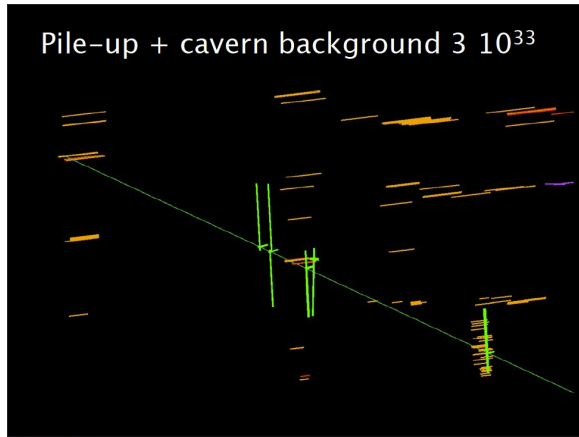
Hough transform

- Each point in x, y belongs to a family of straight lines identified by slope Φ and intercept R_0 , i.e. it is represented by a curve in the $R_0 \Phi$ plane
- The curves in $R_0 \Phi$ from points on the same segment cross at the same $R_0 \Phi$ point
- fill histograms in $R_0 \Phi$
- select maxima
- Very simple and general approach

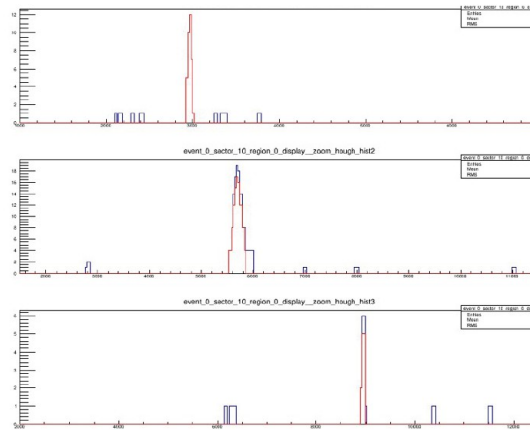


Hough transform, example

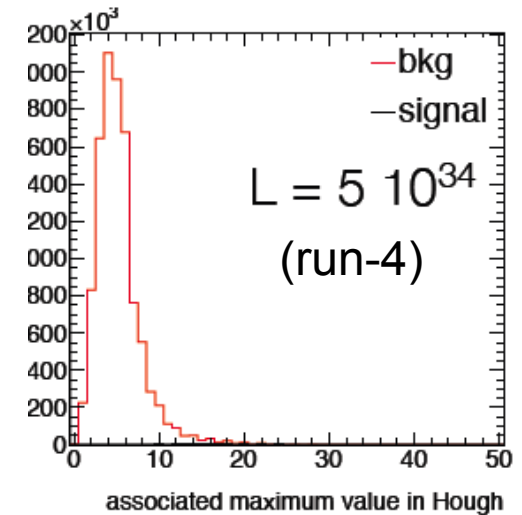
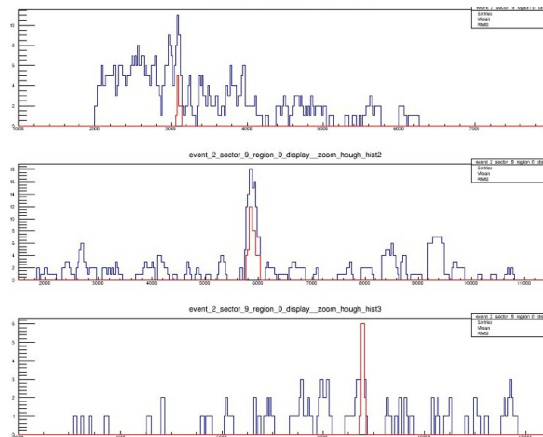
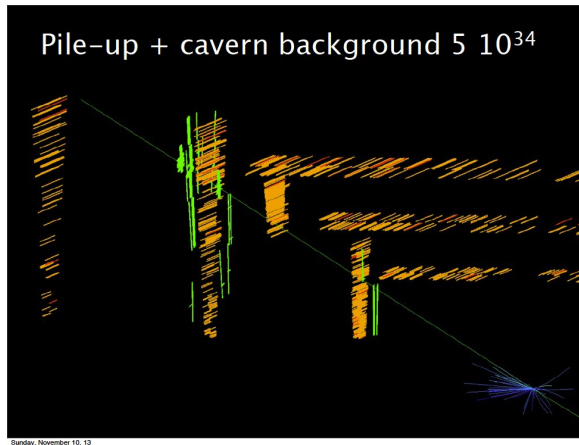
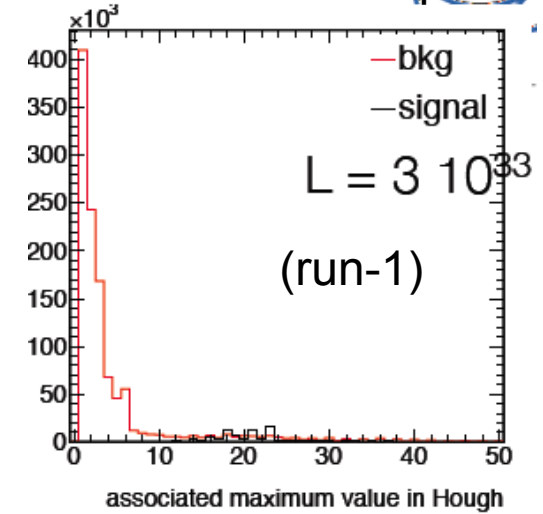
Hits in a sector



Histograms along R_0 , fixed Φ



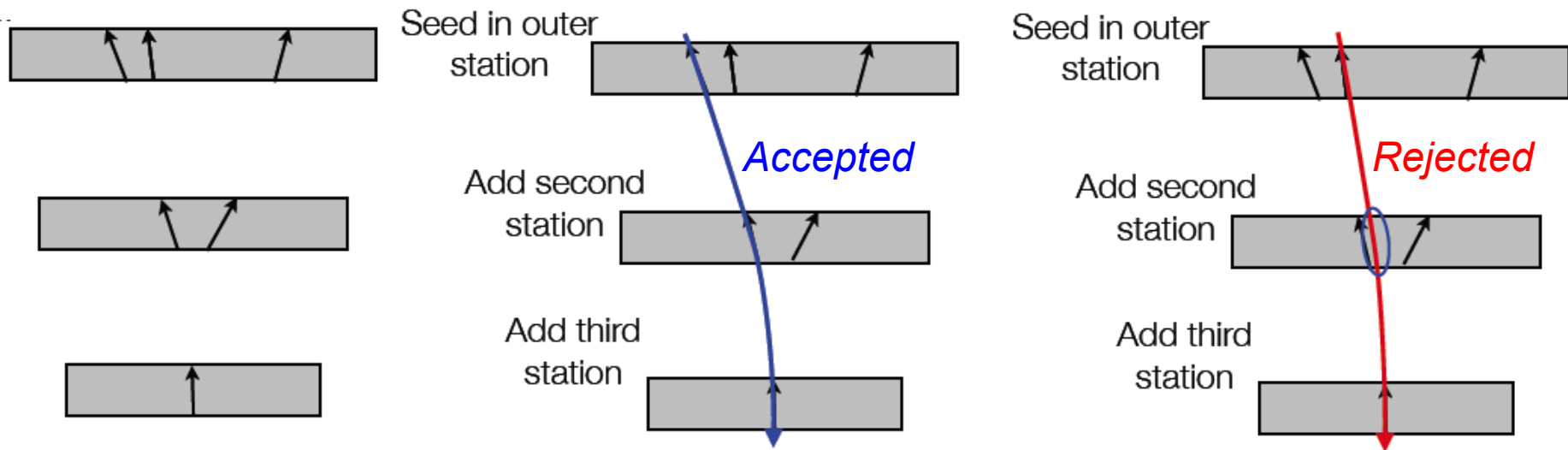
Distribution of hits per bin



Good signal/bkg separation for Run-1/2 backgrounds
Will need to add more constraint for Run-2/3

Pattern recognition in the full MS

- Segments found in different chambers are combined starting from outer layers and following the the track trajectory inward
- Combinations with common segments are removed based on number of holes
- Finally we are left with “MS-only” tracks



Momentum from sagitta measurement

Momentum component perpendicular to B is related to local curvature R by

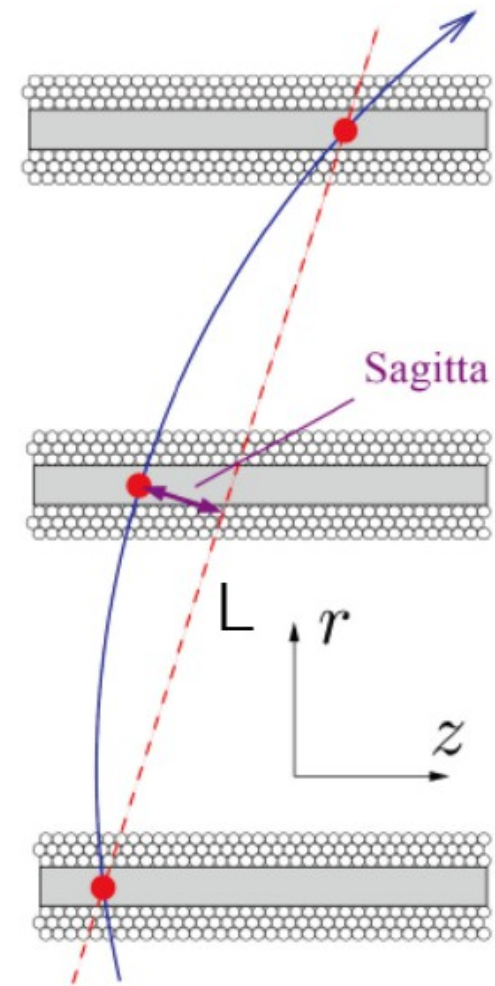
$$p_{\perp} \sim k q B R$$

$$k \sim 0.3 \text{ GeV/T/m}$$

- With three points in a magnetic field we can measure the muon momentum from the sagitta S:

$$q/p \sim 8 S / (k B L^2)$$

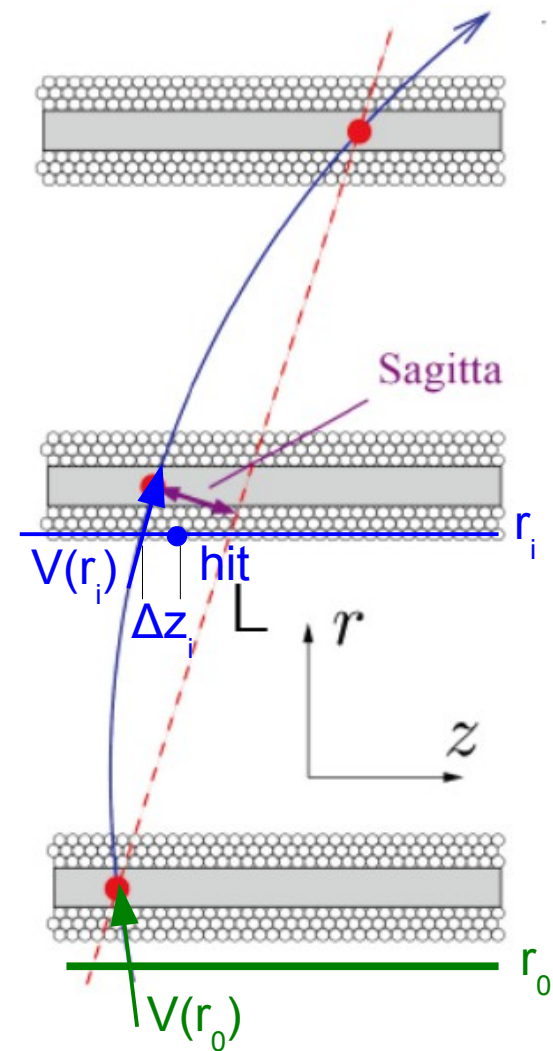
- Typically 1 TeV correspond to $S \sim 1 \text{ mm}$
- In practice the B field is not homogeneous, there are more than 3 measurements, we need to extract the track parameters from a fit



Track fit in the MS

- A track is characterised by 5 parameters,
- e.g. choosing as a reference surface the cylinder with radius r_0 corresponding to the MS entrance :
$$V(r_0) = (q/p, z(r_0), \varphi(r_0), dz/dr(r_0), d\varphi/dr(r_0))$$
- Given $V(r_0)$ it is possible to extrapolate the track to any i^{th} detector layer using a precise numerical transport code, together with a precise map of the B field and of detector positions:
$$V(r_0) \Rightarrow V(r_i)$$

the track covariance matrix is propagated as well.
- A global χ^2 can be calculated from the residuals Δz_i between extrapolated track at r_i and the actual i^{th} measurements.
- A global χ^2 minimization gives the best parameters at MS entrance



Multiple scattering in the MS

- The previous picture is complicated by multiple scattering: the MS contains many radiation lengths of support structures and detectors

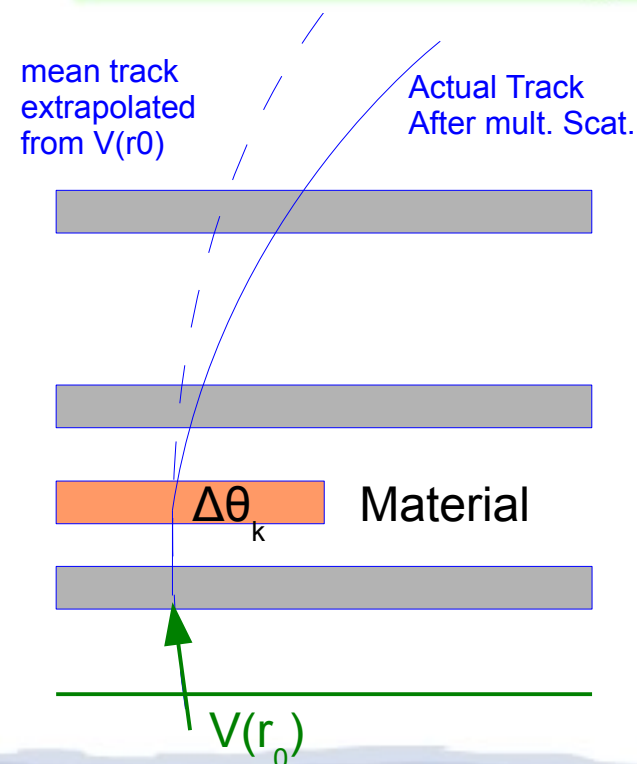
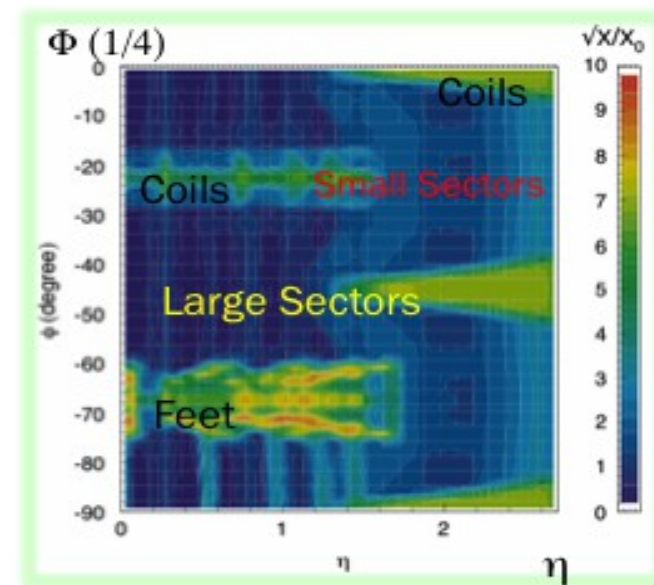
- Multiple scattering: RMS angular deflection from material:

$$\theta_0 = \frac{13.6 \text{ MeV}}{\beta c p} z \sqrt{x/X_0} \left[1 + 0.038 \ln(x/X_0) \right]$$

- Multiple scattering is a stochastic phenomenon that introduces irreversibility into track transport
- Additional “kink” parameters in global fit that allow a deflection (θ_k, φ_k) at reference scattering planes.

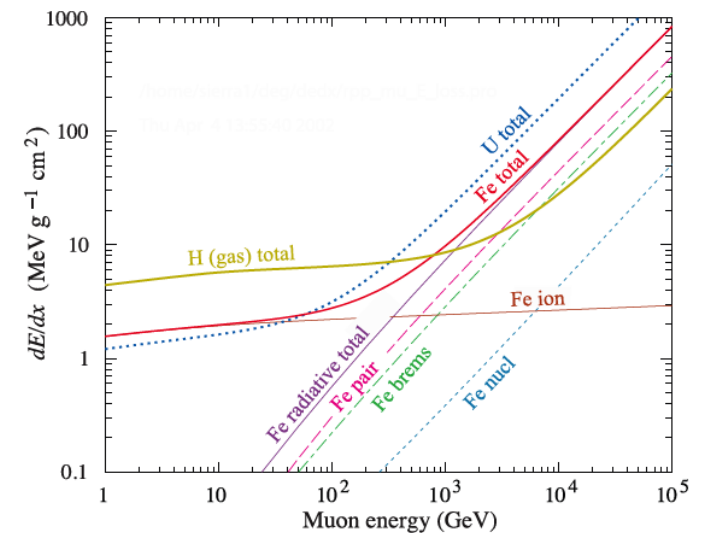
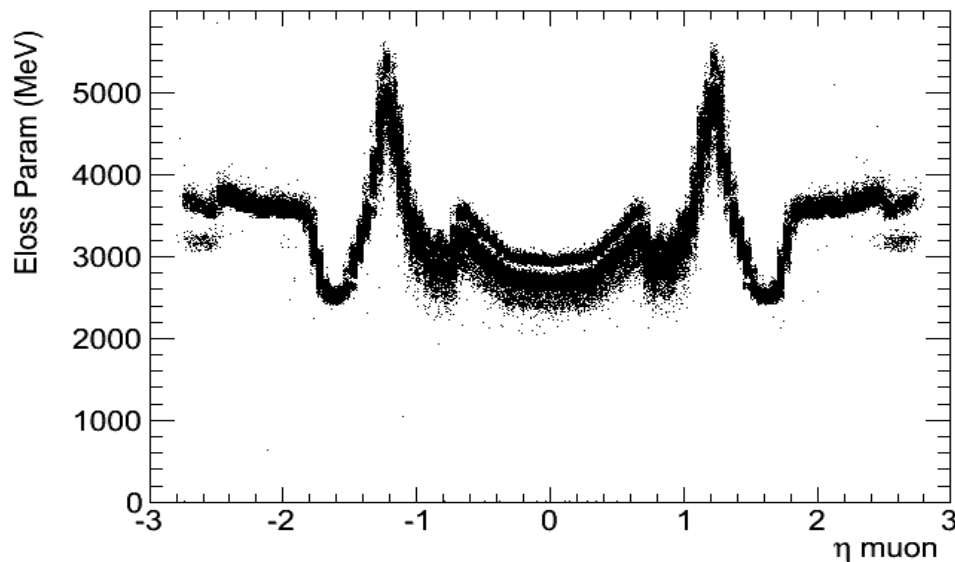
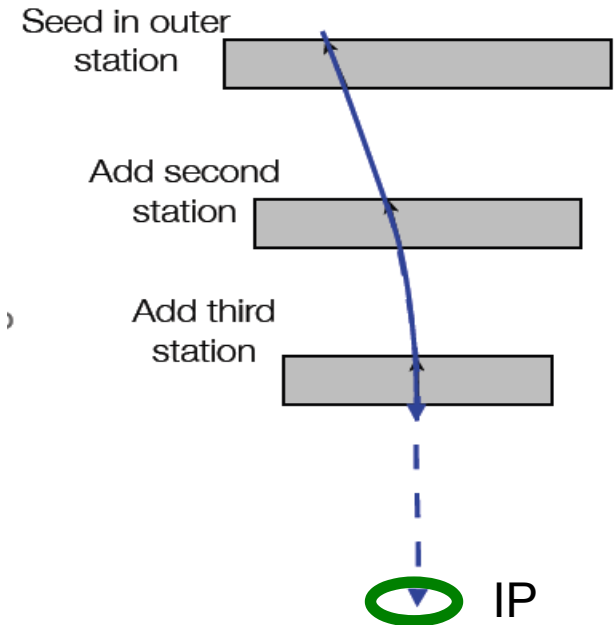
They are treated as nuisance parameter in the fit giving a penalty if the deflection is too large. $\Delta(\chi^2)_k = (\theta_k / \theta_{\text{rms},k})^2 + (\varphi_k / \theta_{\text{rms},k})^2$

- At this point we have the best track at the MS entrance



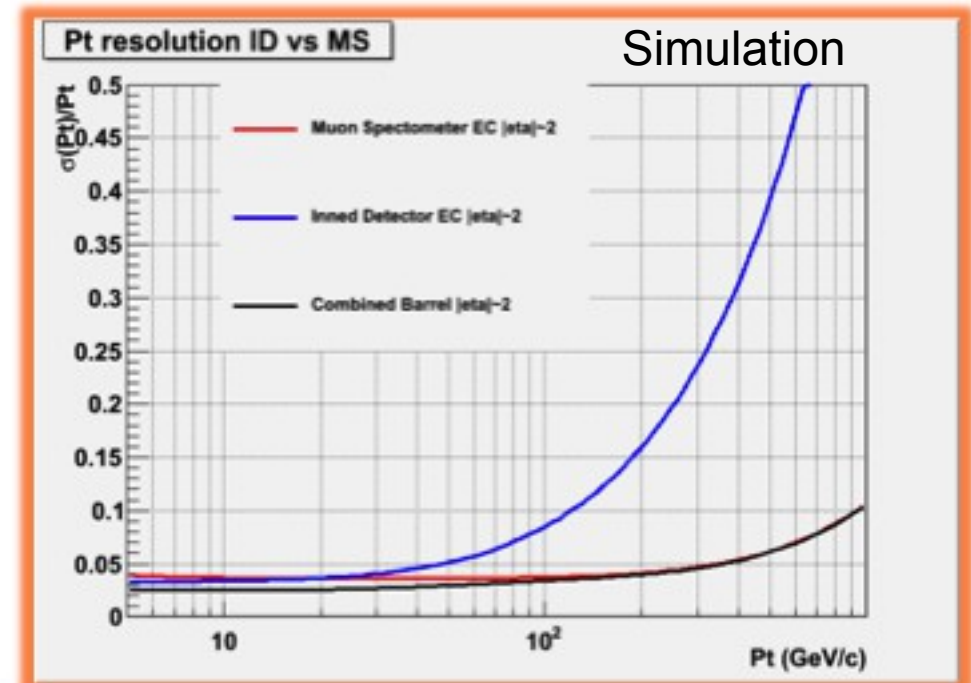
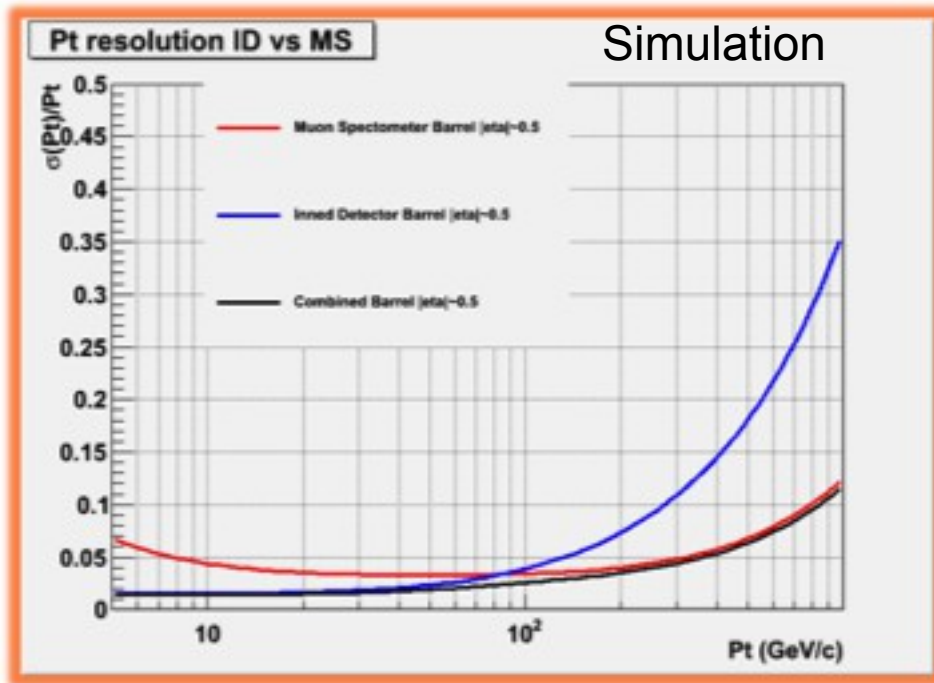
Extrapolation to the IP: energy loss

- What we are really interested in are the muon parameters at the IP: need to back-extrapolate through the calorimeters
- Energy loss : approx. 3 GeV (eta dependent)
- Landau tail: large fluctuations
- Use a combination of parametrization and CAL measurement to estimate energy loss
- CAL measurement used only for isolated muons and for large losses
- Final “Muon Extrapolated” fit including the IP constraint gives the track parameters at the IP



ID-MS combined muons

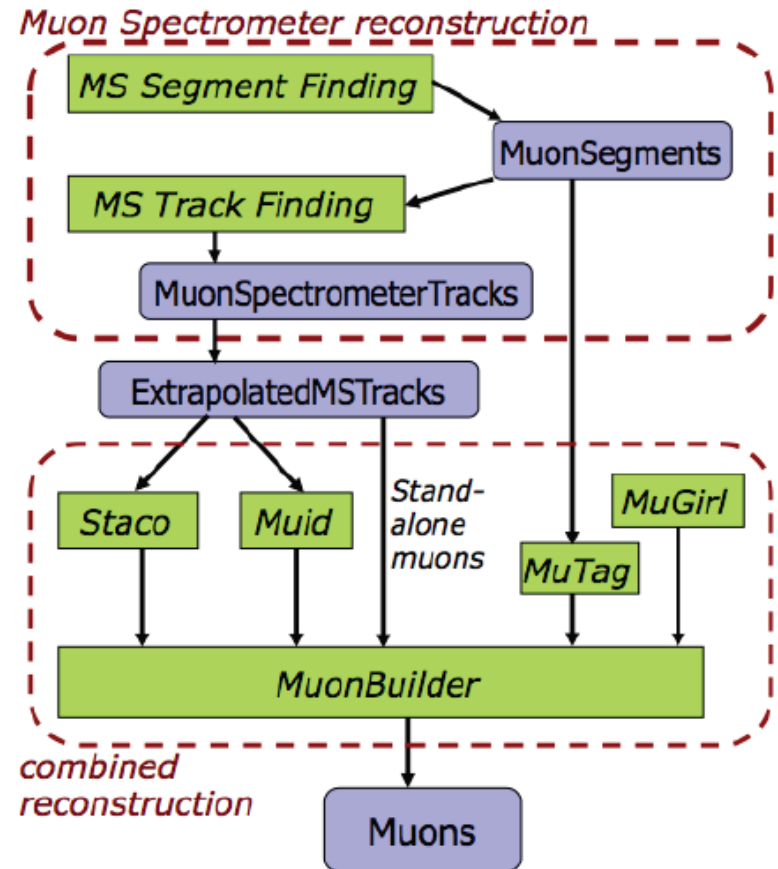
- Outside-In reconstruction: “Muon-Extrapolated” tracks are matched with Inner Detector (ID) tracks to form “Combined Muons”, a full combined fit of ID and MS hits is performed to obtain the final parameters.
- MS dominates the CB measurement for $p_T > 80$ (20) GeV, depending on η
- Inside-Out reconstruction: start from ID tracks and add hits in the MS allows to recover acceptance for low-quality muons with few hits in the MS



Reconstruction output

Final Muon “collection” for analysis

- Outside-in CB muons (best quality): two possible combination algorithms: Muid (track refit) or Staco (statistical combination)
- Inside-out muons: (two algorithms MuGirl, MuTag)
- Stand-Alone Extrapolated muons (recover ID failures, $|\eta| > 2.5$)



Momentum resolution of the MS

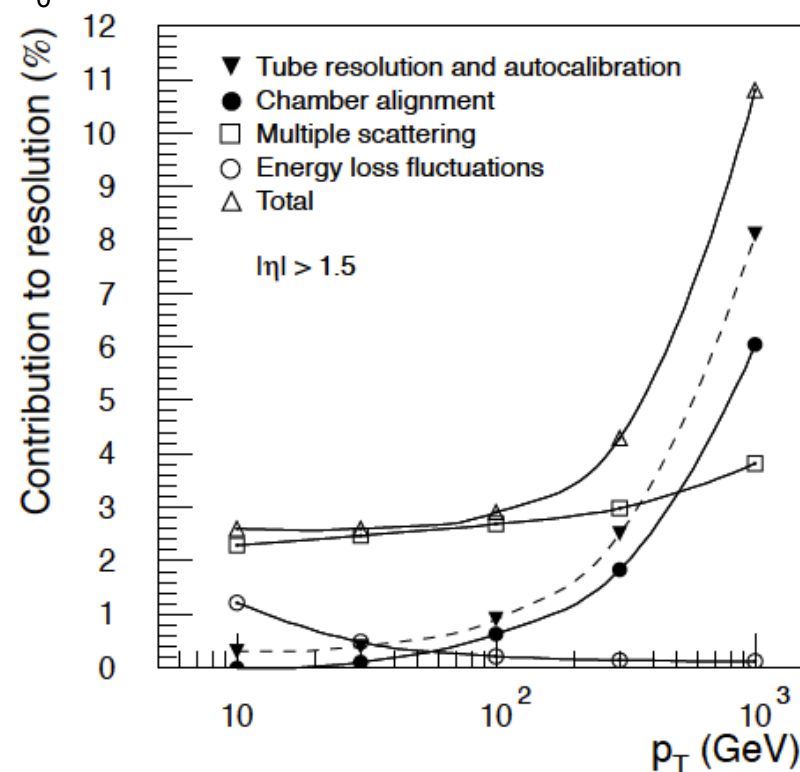
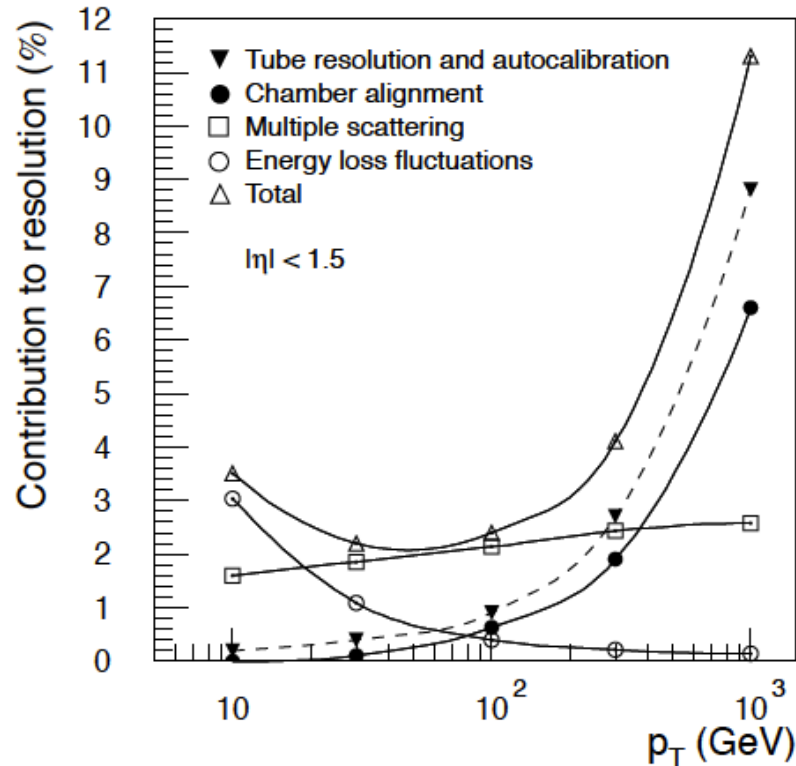
Main contribution to momentum error:

- Error on hit measurements (e.g. uncertainty on sagitta): $\Delta p/p \sim k_2 p$
- Multiple scattering : $\Delta p/p \sim k_1$
- Energy loss fluctuations : $\Delta p/p \sim k_0/p$

For $p_T < 100$ GeV multiple scattering dominates ($k_1 = 2-2.5\%$)

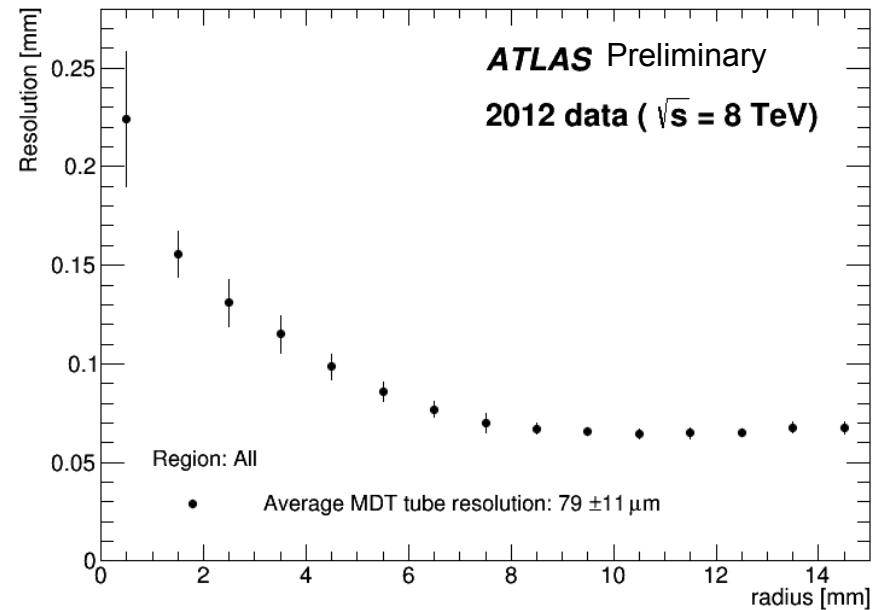
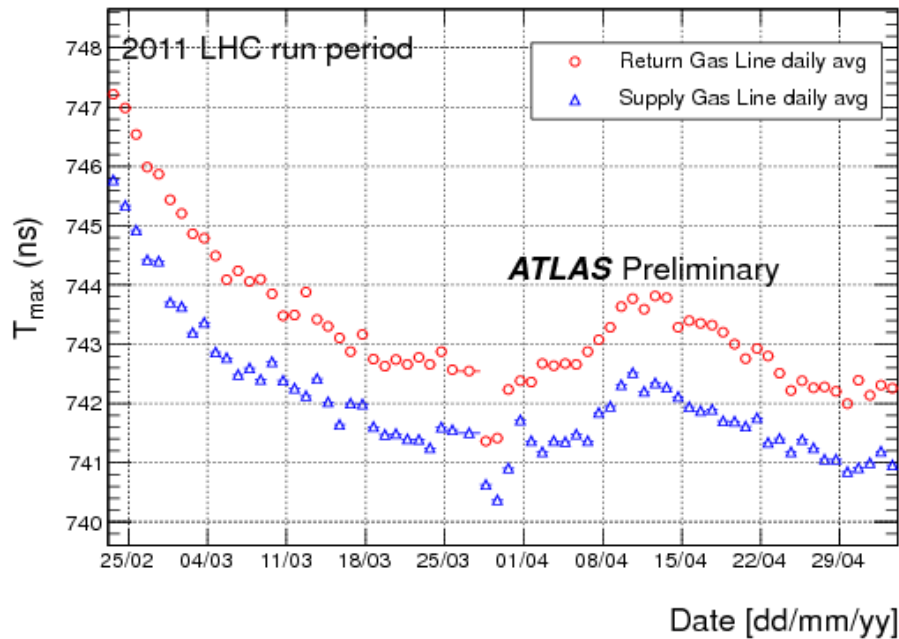
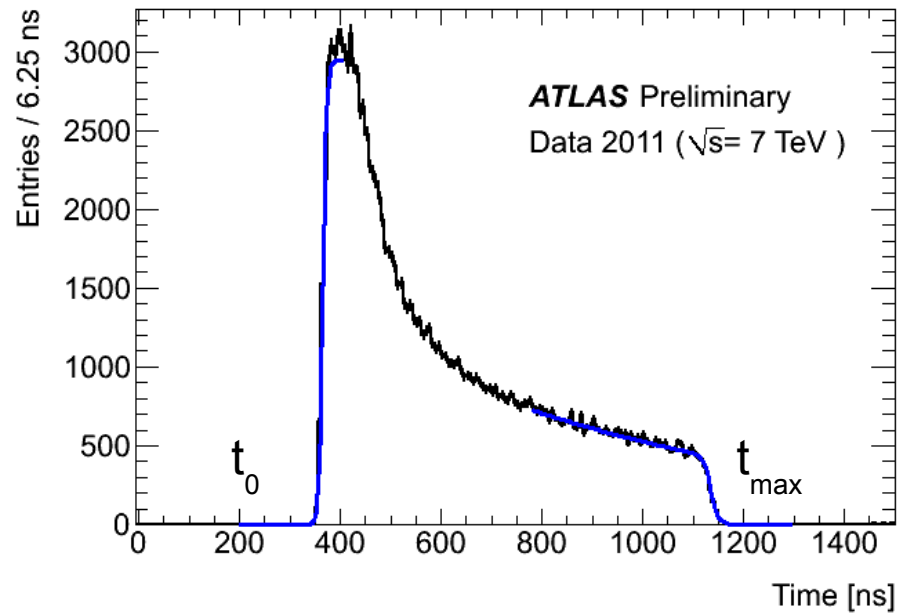
For $p_T > 100$ GeV the “intrinsic term” ($k_2 \sim 10\%/TeV$)

Energy loss fluctuations relevant at low p ($k_0 \sim 250$ MeV)



MDT $r(t)$ calibration

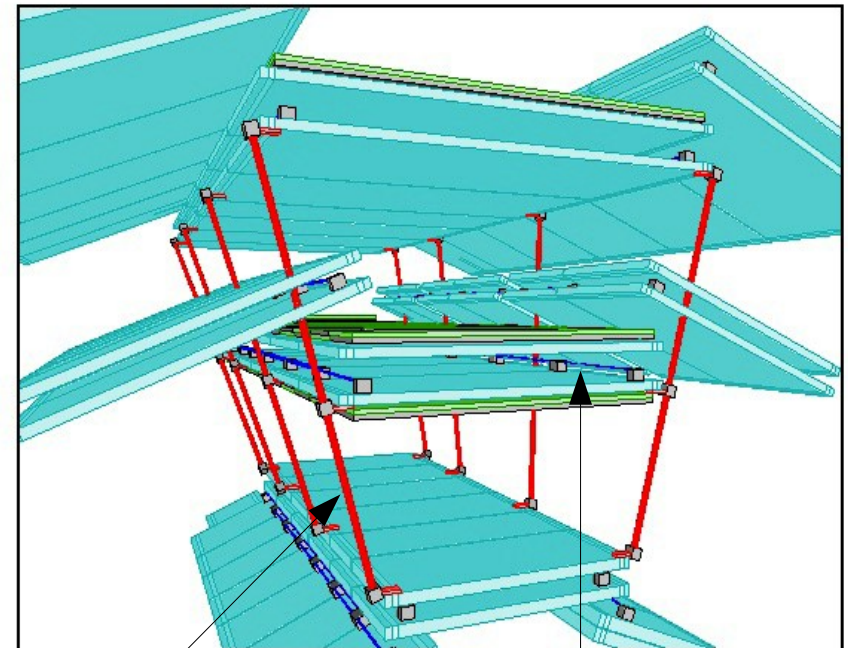
- To arrive to the nominal hit resolution $r(t)$ relation needs to be calibrated, In particular:
 - t_0 tube-by-tube variations
 - t_{\max} : drift velocity variations due to gas temperature, pressure, composition
- Final hit resolution $\sim 80 \mu\text{m}$



Alignment System

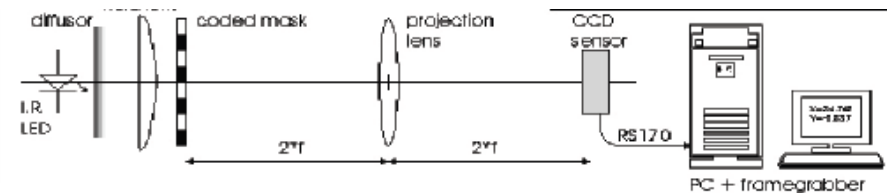
- The “intrinsic” term of p resolution has two components:
 - hit resolution
 - knowledge of detector position: alignment
- The MDT chambers are constructed as precision objects: wires can be located inside a chamber within few tens of μm
- Location and orientation of MDT chambers in ATLAS not trivial: we aim at precision $<50 \mu\text{m}$ over distances of $O(\sim 10 \text{ m})$
- “Absolute” alignment based on tracks (next page)
- Optical alignment system used to
 - follow the relative displacements between different alignment runs
 - Constraint “weak modes”

Barrel optical alignment



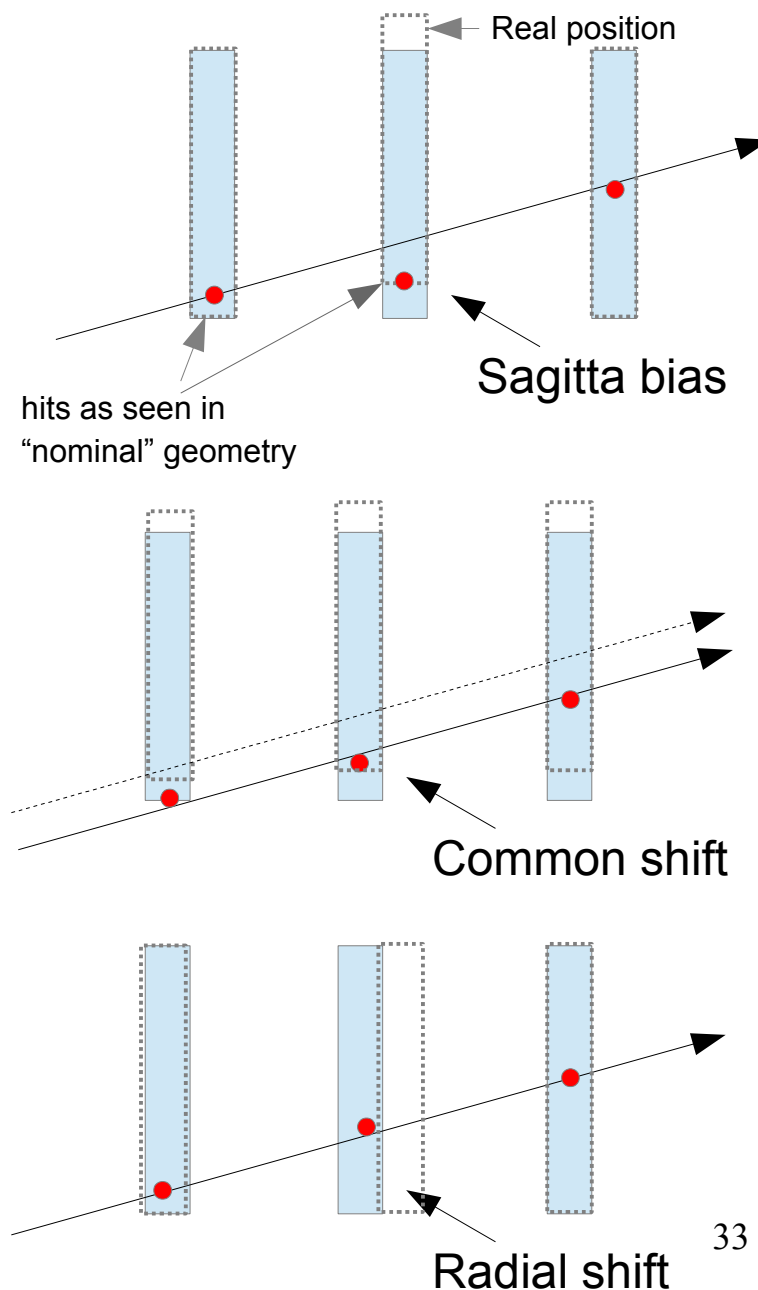
Projective lines

Praxial lines



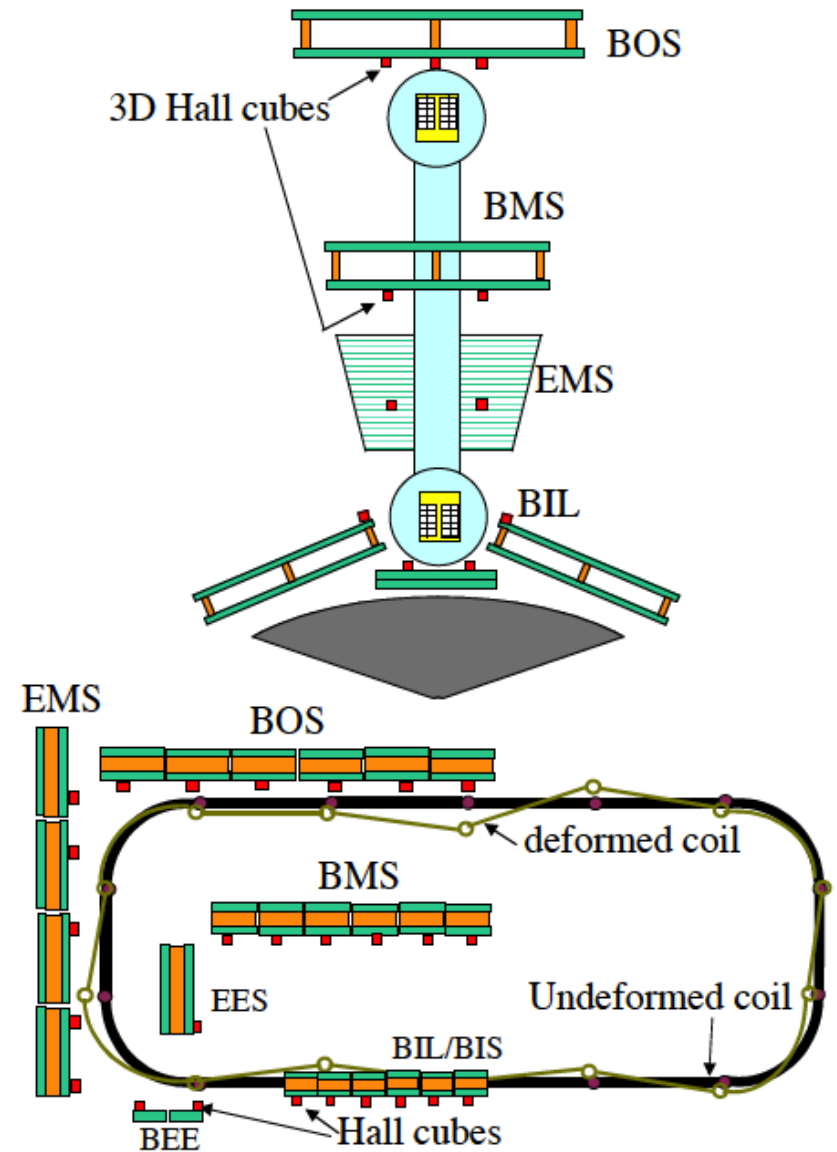
Track-based alignment

- “Absolute” alignment is performed in special runs with the toroidal magnetic field off
 - all tracks are approx. straight lines
 - no sensitivity to knowledge of B field or material
 - ID (solenoid) allows to select high-pT tracks to reduce multiple scattering
- Cosmic rays are used for the Barrel
- Special collision runs for End-Caps (expensive !)
- In practice only sensitive to sagitta bias
- weak modes: common shifts + radial distortions partially recovered from overlaps between sectors and cosmic rays crossing different sectors
- Current precision on sagitta bias $\sim 40 \mu\text{m}$ RMS using $\sim 50\text{M}$ events from collisions with toroid off



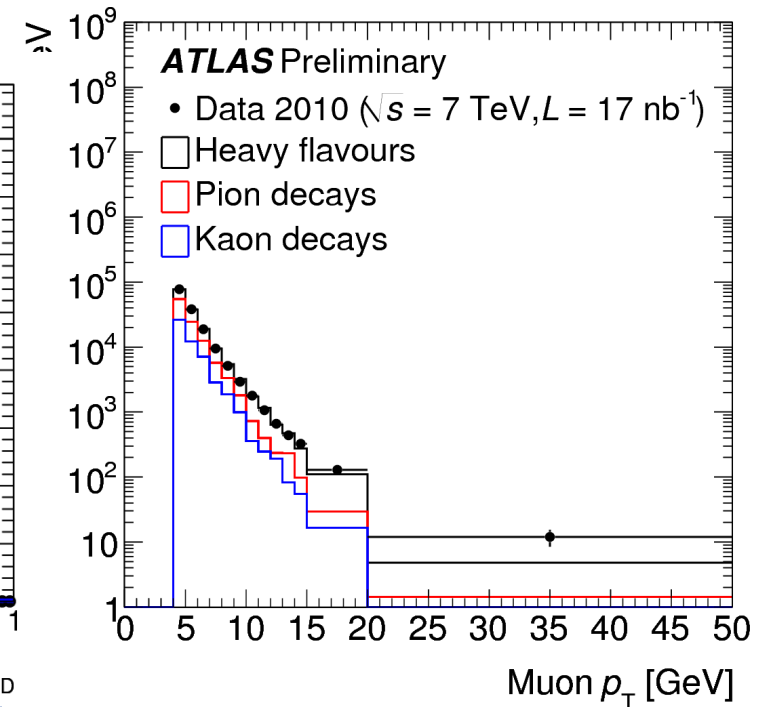
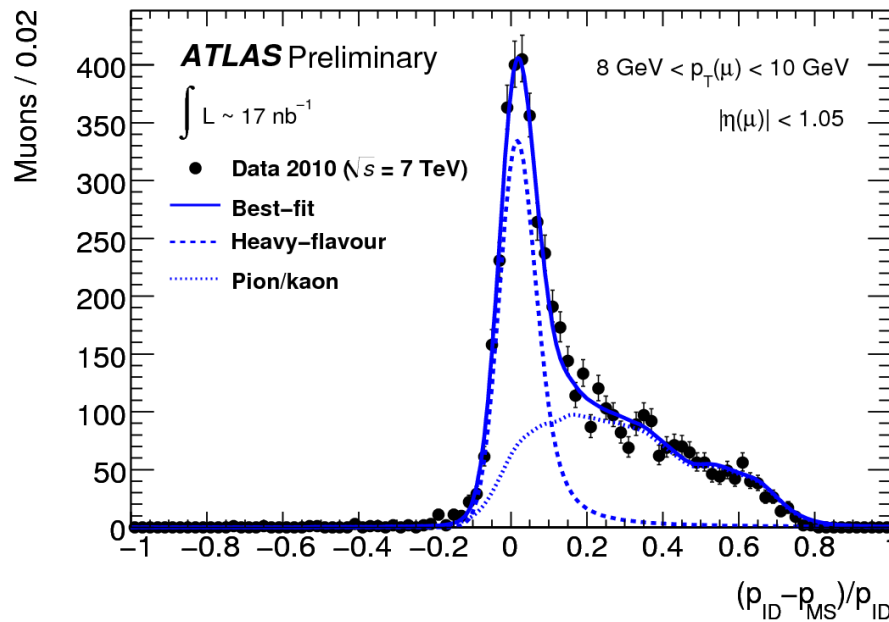
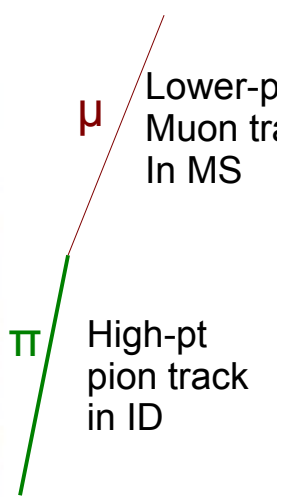
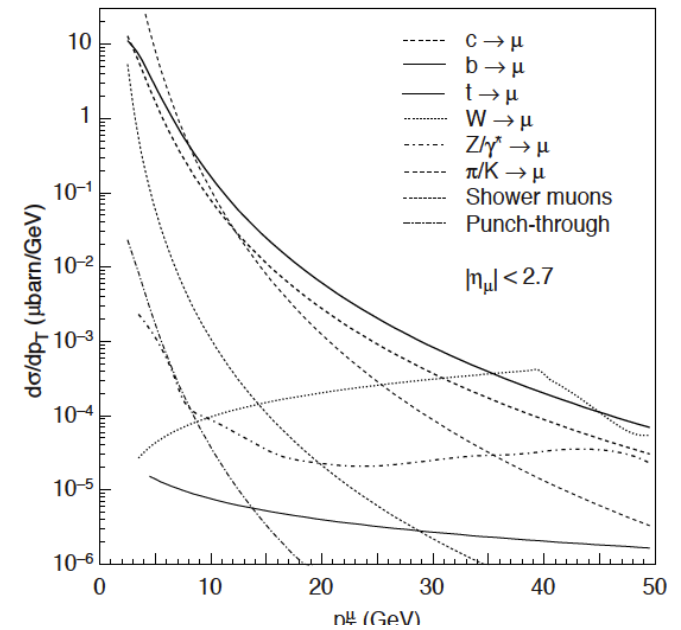
Magnetic field measurement

- The B field integral (actually $\int B L dL$!) should be known precisely to avoid momentum scale biases
- B field maps are made with numerical codes based on the Biot-Savart law, plus non-linear perturbations from ferromagnetic materials (e.g. calorimeters, iron supports of the calorimeters, iron inside concrete walls, cranes etc.)
- The currents in the coils are known precisely, the actual shape of the superconductor coil inside the cryogenic vessel is not so well known
- 3D Hall probes measure B on each chamber => coil shape is fitted to get the best model/measurements agreement
- B known to $\sim 3 \cdot 10^{-3}$ T



Backgrounds from π , K decays

- Muons in the MS originate from
 - π /K decays
 - heavy quark decays
 - (Z,W decays)
- π /K decays can be removed with tighter cuts on ID-MS momentum difference
- Probability that a π /K is identified as a muon ($p_T=20$ GeV) : $\sim 0.2\%$, $\sim 0.1\%$ with tight cut on ID-MS momentum matching

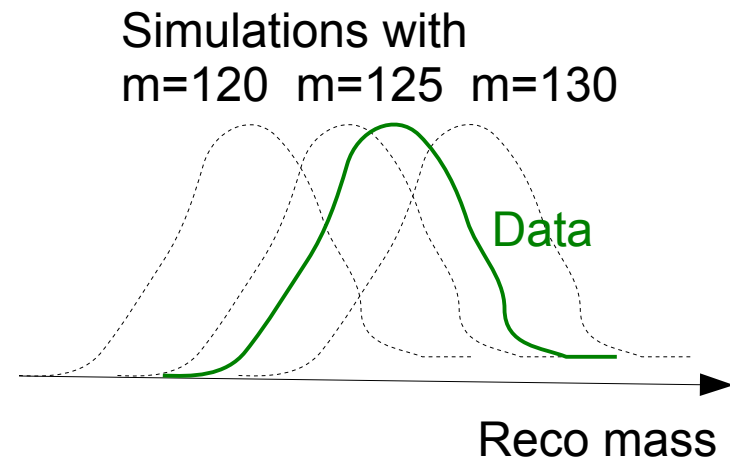


Measurement of performance

- Main Performance parameters to be measured in data
 - Efficiency
 - Momentum resolution and scale
 - In physics analyses MC simulations are used to unfold the detector response
 - The differences in efficiency and momentum resolution/scale are compared between data and simulation
- => corrections are applied to MC simulation to give the best description of the data

Example Higgs mass:

*Not so important that reco mass is correct
But that Simulations reproduces the data*



Efficiency: Tag and Probe method

- Need a sample of “unbiased” muons
- “Tag and probe” method : select $Z \rightarrow \mu\mu$ by
 - Tag: isolated high- p_T CB muon
 - Probe: ID track making the correct invariant mass once combined with the Tag

- Use the probe to check combined muon efficiency (given an ID track):

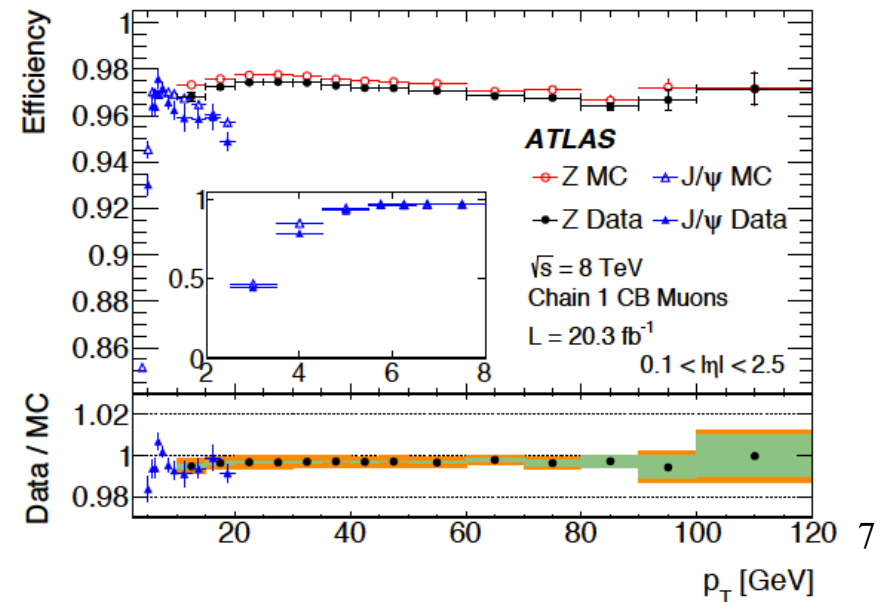
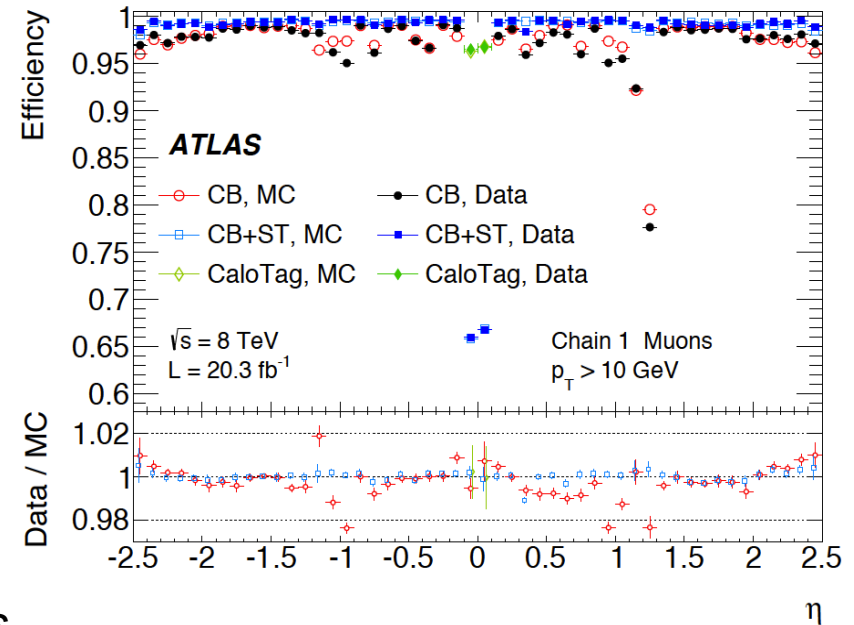
$$P(\text{CB} | \text{ID}) = \frac{N(\text{probes matched to CB})}{N(\text{probes})}$$

- The approach can be inverted using MS muons as probes to check the ID efficiency (TP approximation):

$$\begin{aligned} \text{eff}(\text{CB}) &= P(\text{CB} | \text{ID}) P(\text{ID} | \text{true-}\mu) \\ &\sim P(\text{CB} | \text{ID}) P(\text{ID} | \text{ME}) \end{aligned}$$

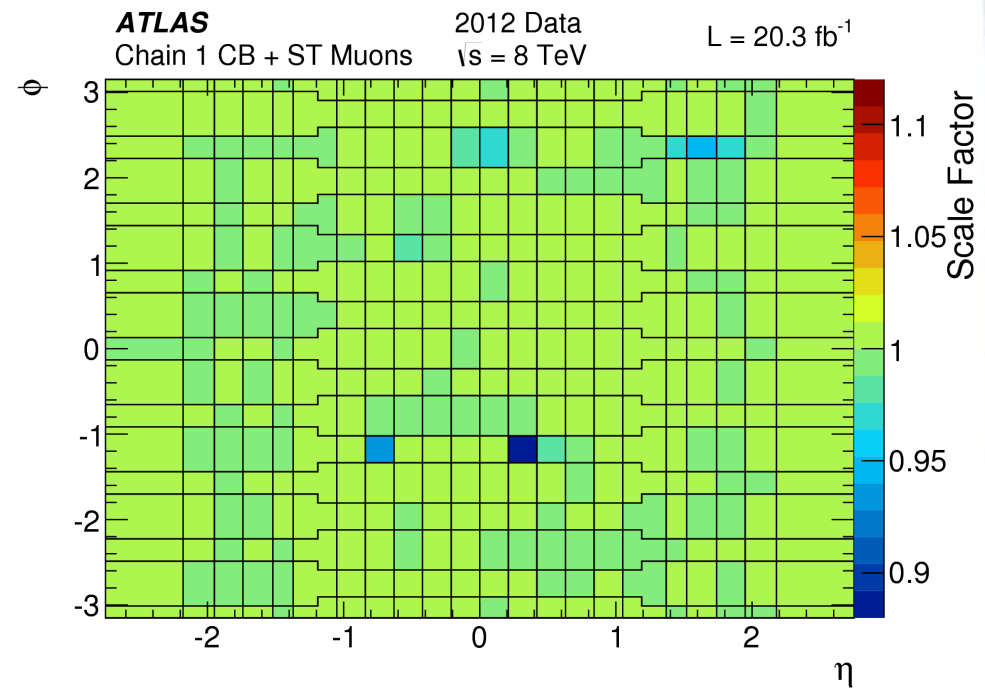
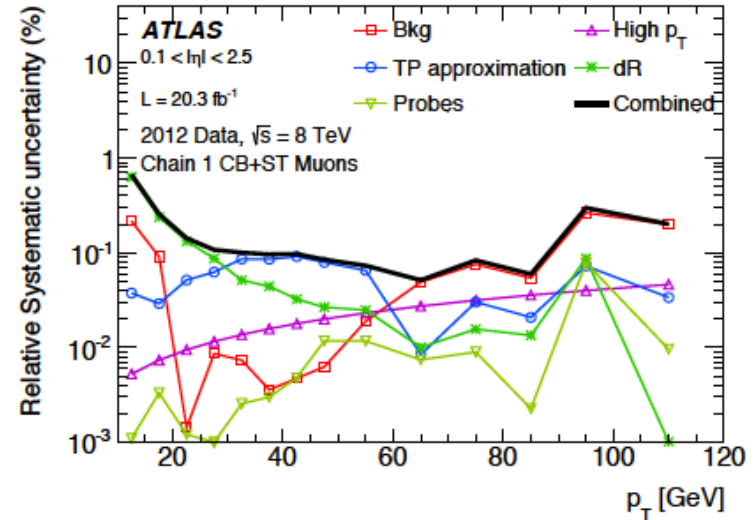
- Requirement that Calorimeter deposit associated to ID probes is compatible with a muon (CaloTag) to reduce the remaining backgrounds

- J/ψ decays used for low- p_T



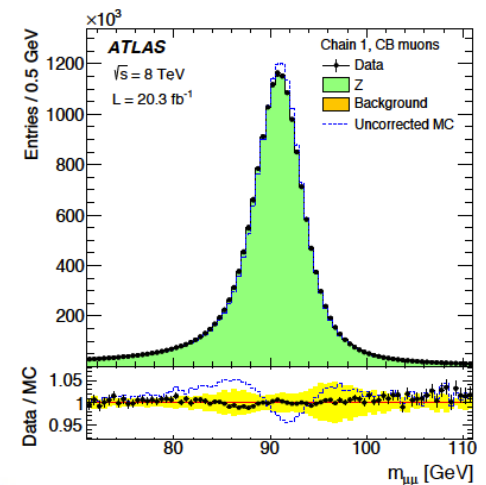
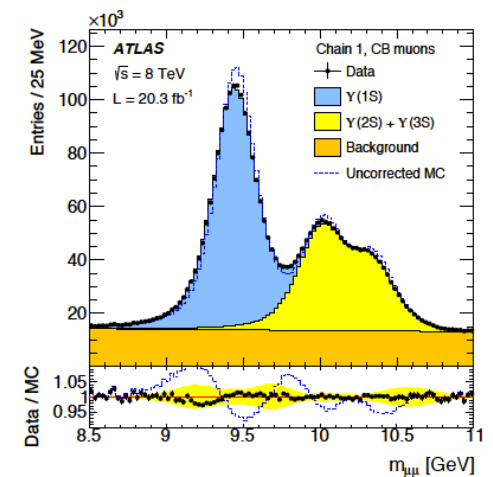
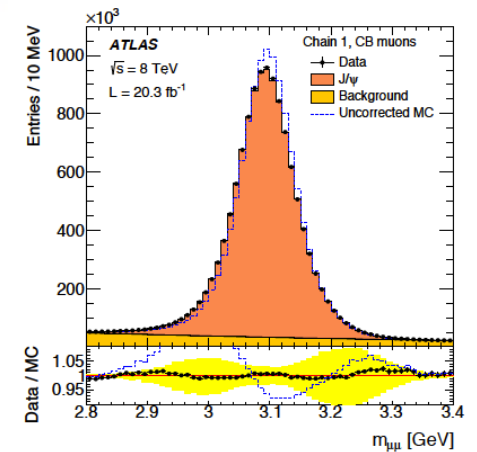
Efficiency: uncertainties and Scale Factors

- Main systematics uncertainties from
 - “TP approximation” estimated comparing measured and true efficiency in MC
 - backgrounds at large p_T estimated from same-sign dimuons and MC
- Scale factors for physics analysis: η - ϕ maps of $\text{eff}(\text{Data})/\text{eff}(\text{MC})$ to be used to correct MC in physics analyses
- Data/MC differences in general within 1%
Few differences due to problematic chambers, or to low efficiency of trigger chambers.



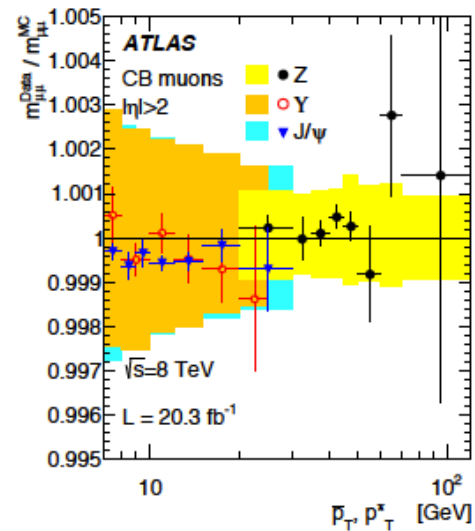
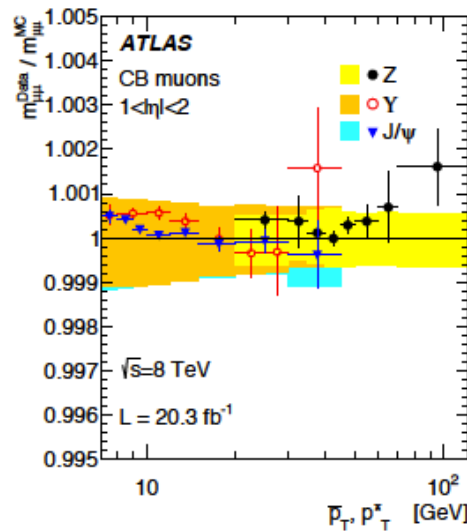
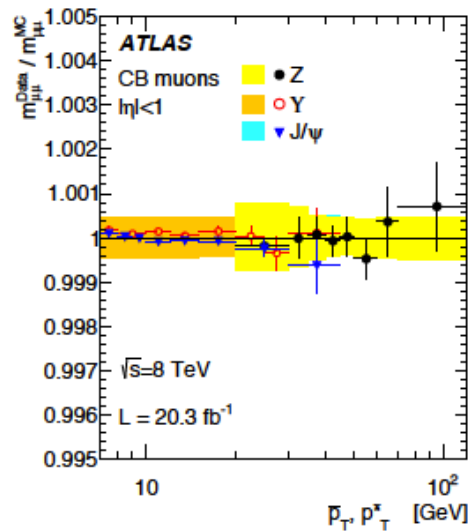
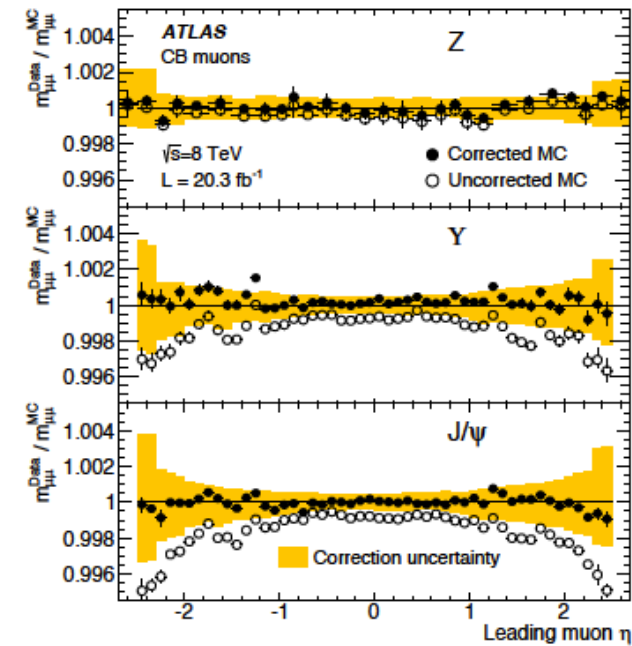
Momentum corrections

- Resonances of well known mass are used to “calibrate” the momentum response.
- MC corrections (in η bins):
 Scale: $p_T \rightarrow s_0 + p_T (1 + s_1)$
 Resolution: add random smearing terms to p_T
 of sigma $\Delta r_0, \Delta r_1 p_T, \Delta r_2 p_T^2$
- The best parameters are obtained by comparing data and smeared MC distributions for a set of invariant mass distributions from J/ψ and Z samples.
- Same approach repeated for ID and MS measurements separately. Then comined to obtain the correction for CB muons



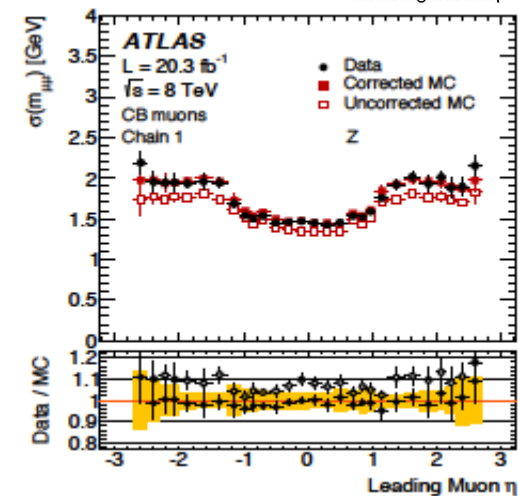
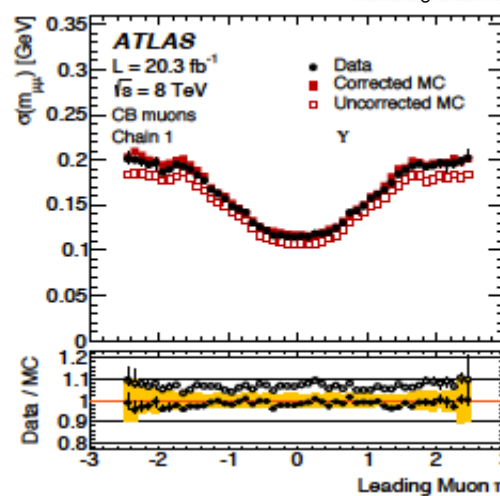
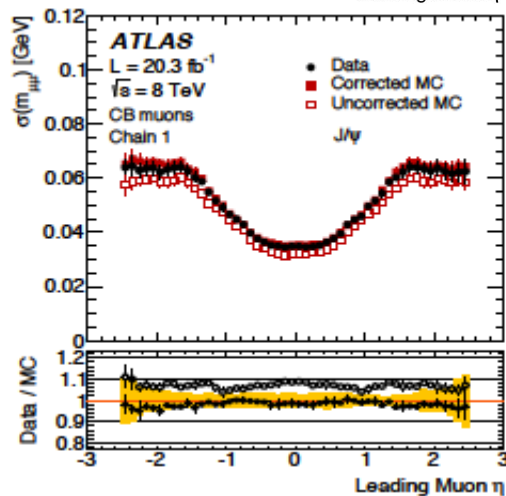
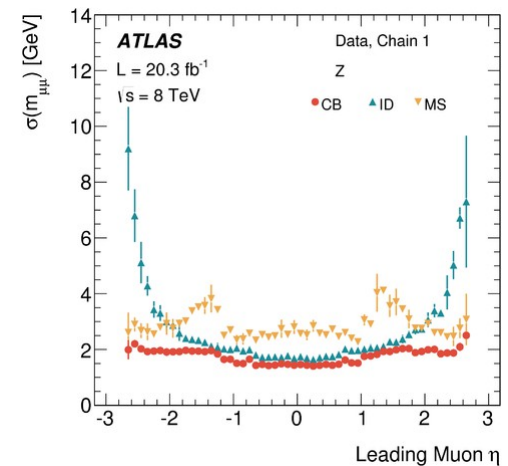
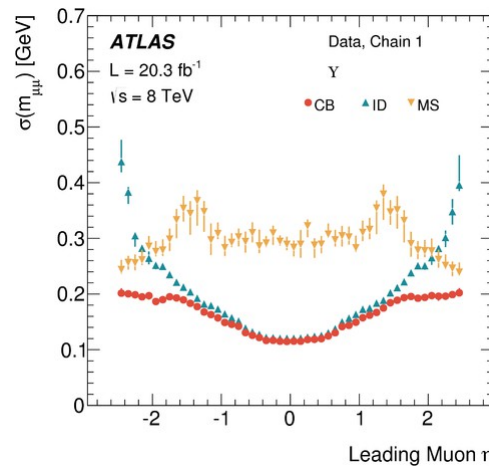
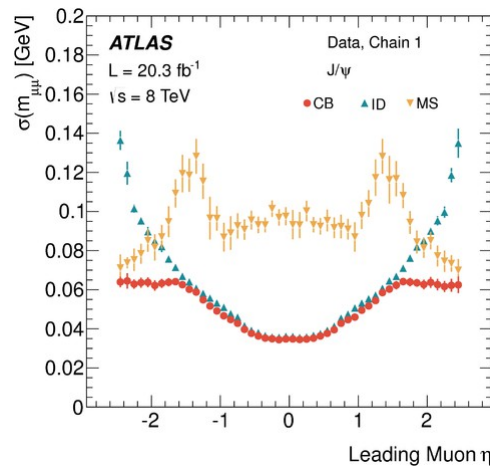
Momentum scale: results

- Momentum scale in data wrt ideal MC :
 - 0.1% offset in ID scale
 - bias at low-pt (E-loss)
- Corrected MC agrees with data within uncertainties (<0.1% from fit syst + stat)



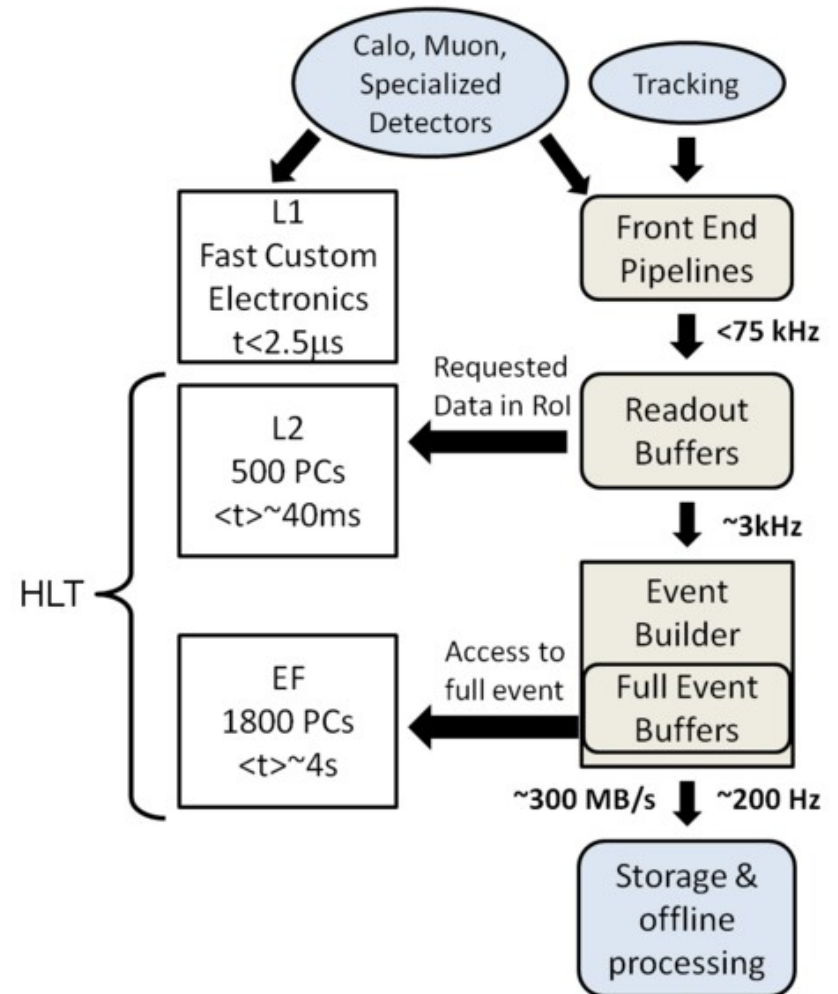
Momentum resolution: results

- Good data/MC agreement after correction
- Mass resolution $\sigma(m)/m \sim 1/\sqrt{2} \sigma(p)/p$
- At the Z: $\sigma(m)/m \sim 1.5$ to 2%.



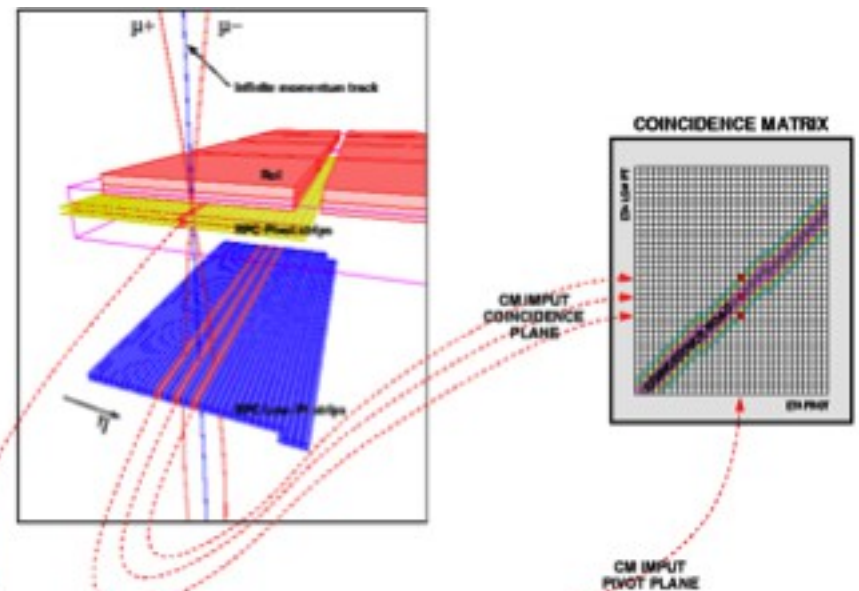
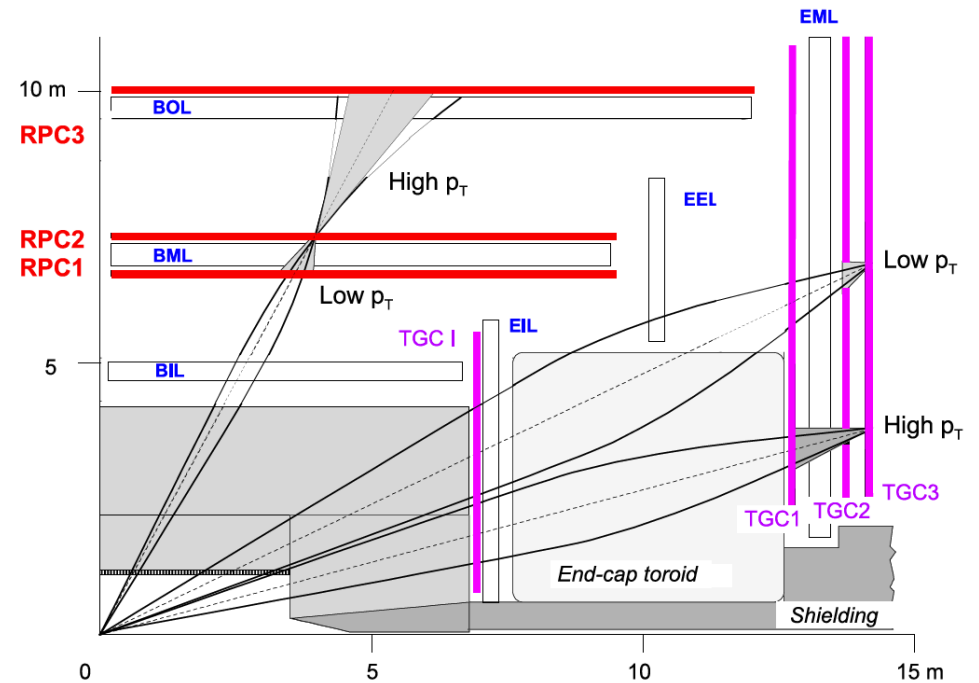
ATLAS Trigger

- In Run-2 ATLAS has a two-level trigger:
- Level 1, hardware,
Input 40 MHz => Output 100 kHz
Time to take a decision (latency) < 2.6 μ s
- High-Level Trigger (HLT) software on a computer farm
Input 100 kHz => output ~1 kHz
- In Run-1 the HLT was further divided in level-2 (reading only partial data) and Event Filter (EF)



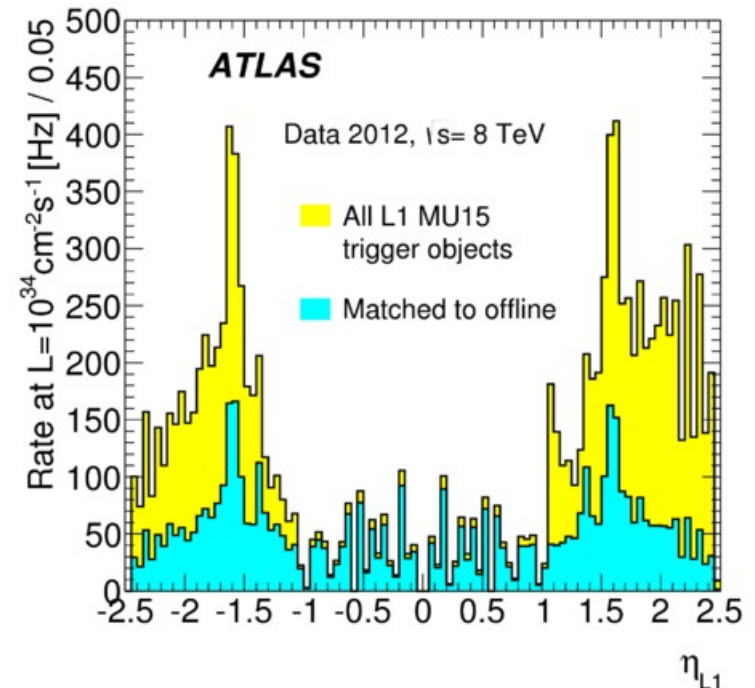
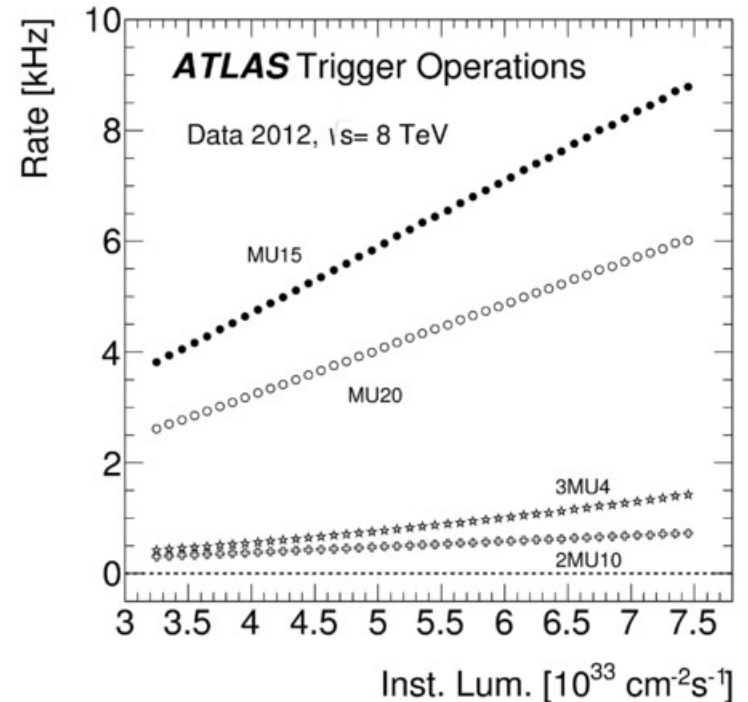
L1 Muon Trigger

- Level-1 muon Trigger
- Barrel:
 - low-pt: (4-10 GeV) two-stations
 - High-pt: (11-20 GeV) three stations
- Endcap:
 - two-stations (4 GeV)
 - three stations (6-20 GeV)
 From Run-2 additional coincidence on TGC-Inner
- Barrel (Roma-1):
 - “coincidence matrix” ASIC performs programable space and time coincidence between pivot and confirm planes
 - η and ϕ matrices are combined by the “Pad” board



Trigger: rates

- L1 rate allocated by ATLAS for single muons is ~ 20 kHz (out of 100 kHz total)
- In Run-1 $p_T > 15$ GeV threshold, well within 10kHz
- In Run-2 expected ~ 20 kHz with $p_T > 20$ GeV (factor ~ 2 for increased luminosity)
- Most rate from the Endcaps, mainly charged tracks (protons) from secondary interactions downstream of the IP
- Partly reduced in Run-2 with further coincidence with inner plane
- Dimuon triggers 2MU10 (plus 2MU4, 2MU6 for B physics)



Efficiency and thresholds

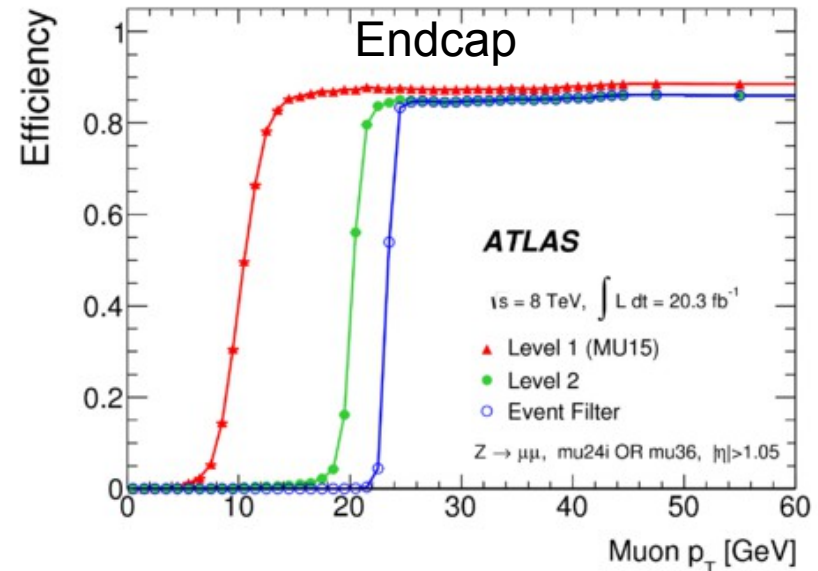
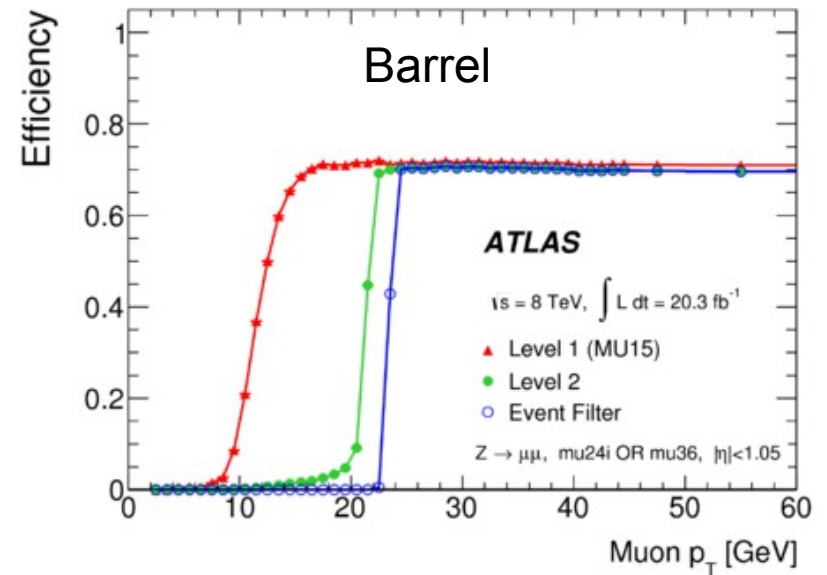
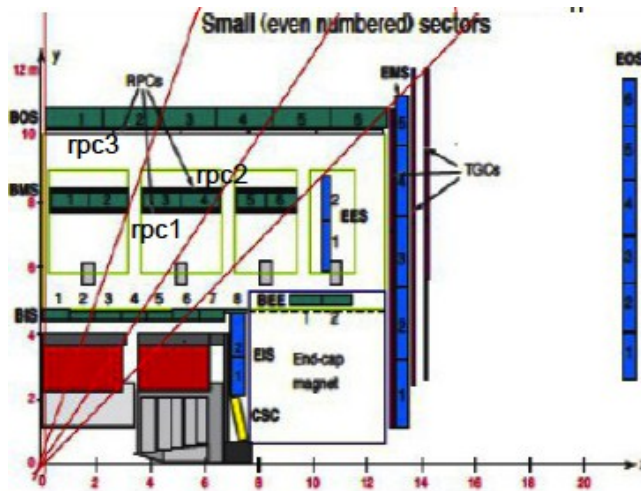
Efficiency “turn-on” curves for

- L1 : $p_T > 15$ GeV
- HLT : $p_T > 24$ GeV

Barrel ~70%:
due to large acceptance holes for coils support and atlas structures and calorimeter services

Endcap: ~90%

Holes
In barrel
3-station
trigger

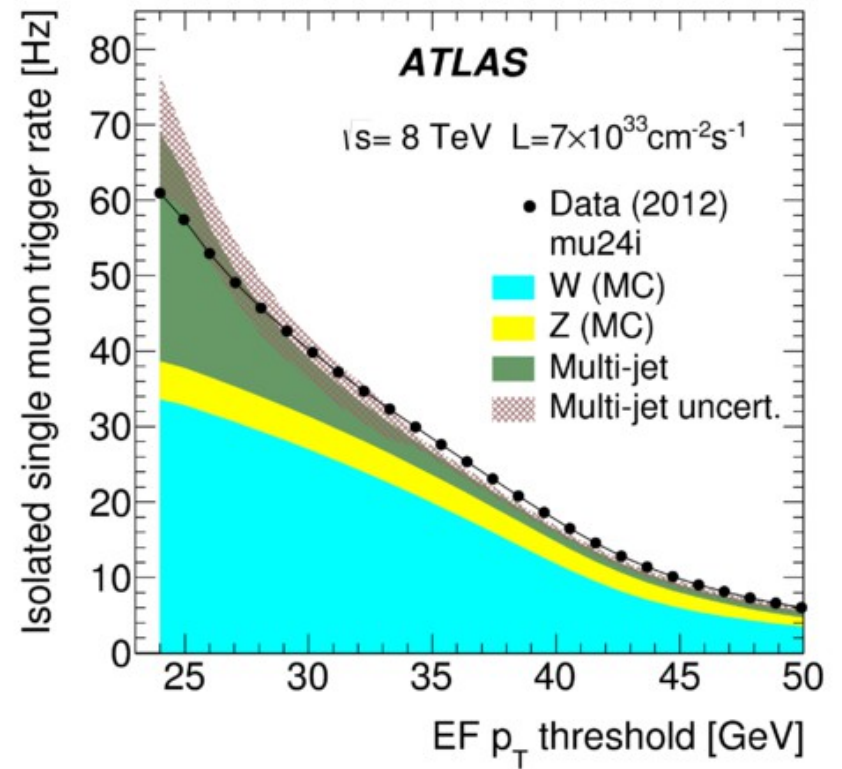


Summary

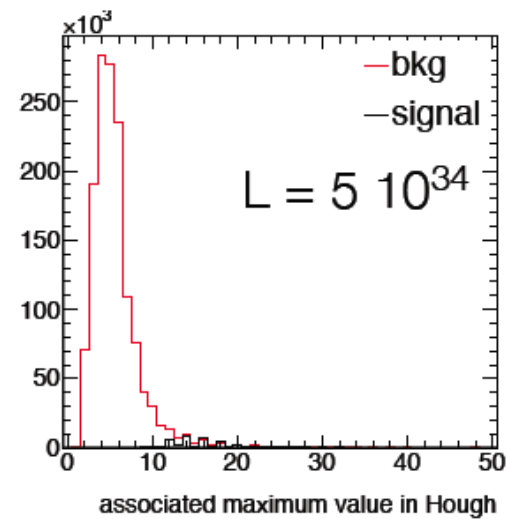
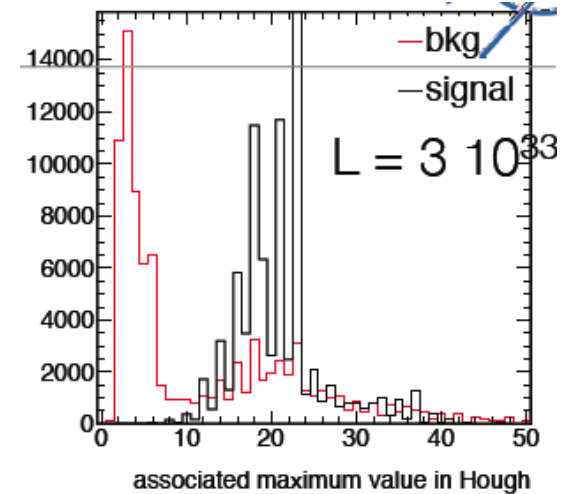
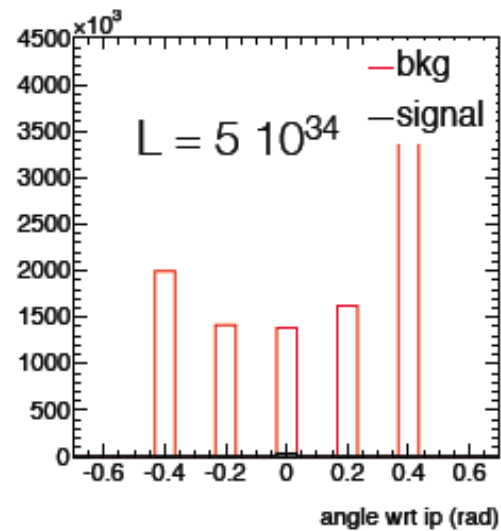
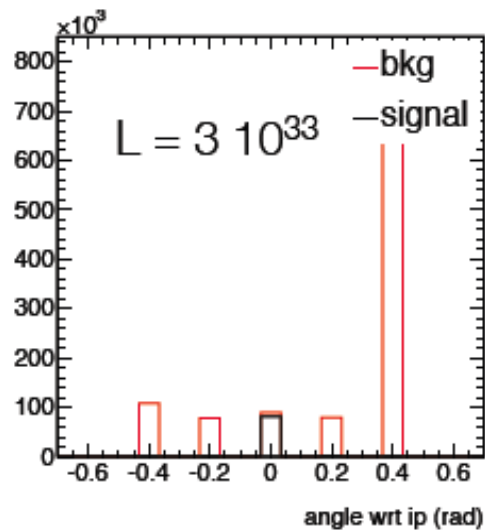
- Overall design
- Track reconstruction
- Performance measurements
- Trigger
- Outlook

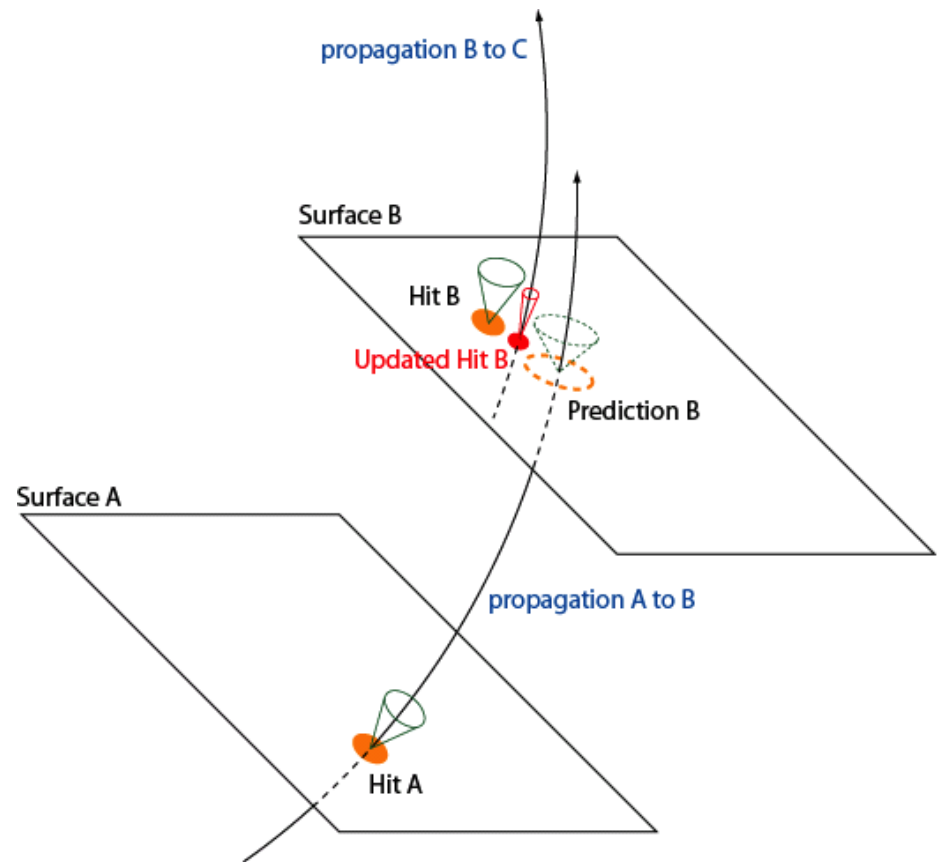
Backup slides

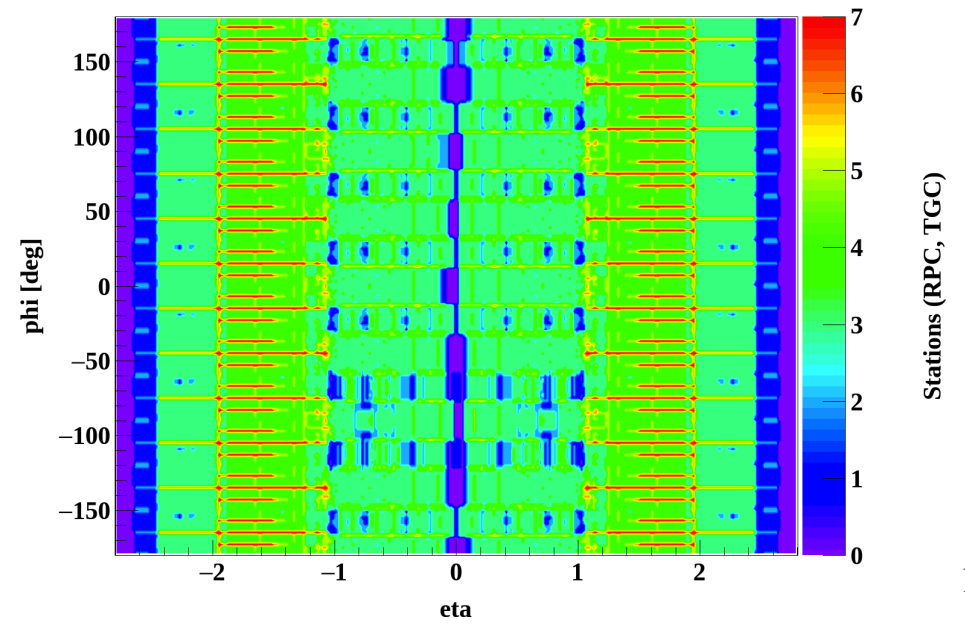
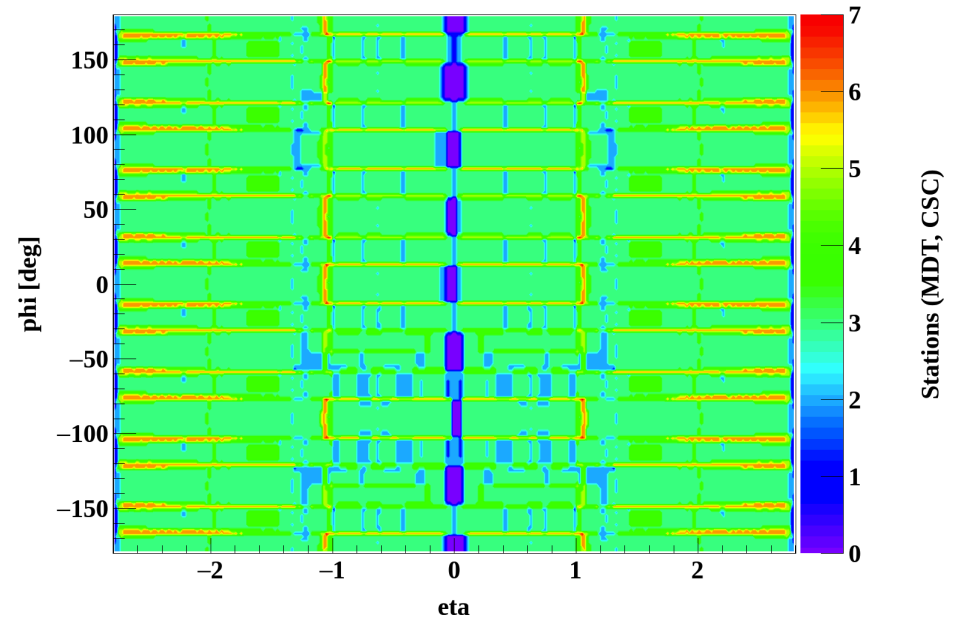
HLT : rates



Hough transform, selecting only pointing segments







Momentum is determined by measurement of **track curvature** $\kappa = 1/\rho$ in B field:

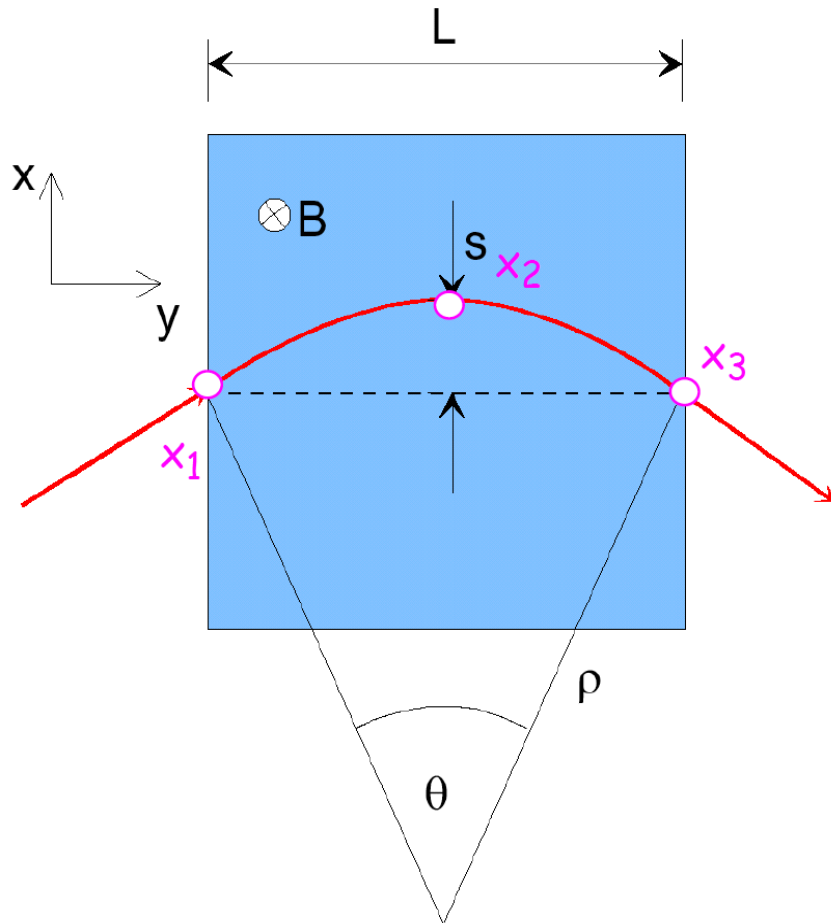
Measure **sagitta** s of the track. For the momentum component transverse to B field:

$$p_T = qB\rho$$

Units: $p_T[\text{GeV}] = 0.3B[\text{T}]\rho[\text{m}]$

$$\frac{L/2}{\rho} = \sin\frac{\theta}{2} \approx \frac{\theta}{2} \text{ (for small } \theta) \Rightarrow \theta \approx \frac{L}{\rho} = \frac{0.3B \cdot L}{p_T}$$

$$s = \rho\left(1 - \cos\frac{\theta}{2}\right) \approx \rho\left(1 - \left(1 - \frac{1}{2}\frac{\theta^2}{4}\right)\right) = \rho\frac{\theta^2}{8} \approx \frac{0.3L^2B}{8 p_T}$$



For the simple case of **three measurements**:

$$s = x_2 - (x_1 + x_3)/2 \Rightarrow ds = dx_2 - dx_1/2 - dx_3/2$$

with $\sigma_x \approx dx_i$ uncorrelated error of single measurement:

$$\sigma_s^2 = \sigma_x^2 + \frac{\sigma_x^2}{4} \cdot 2 = \frac{3}{2}\sigma_x^2$$