

Unitarity and Universality  
from Kaons with KLOE

*Paolo Franzini*

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# History

1947: Rochester and Butler,  $K^0 \rightarrow \pi^+ \pi^-$ ,  $K^+ \rightarrow \pi^+ \pi^0$

1953: Gell-Mann (Nisijima) strangeness

1958: Gell-Mann and Feynman Universal V-A interaction.  
 $\Lambda \rightarrow p e \nu \sim 1.6\%$ ,  $\Sigma^- \rightarrow n e^- \nu \sim 5.6\%$

1958: Columbia-Pisa: should see 6 events each, none found

1963: Cabibbo,  $s - d$  mixing.  $\sin \theta_c = .26$

1970: GIM, 2 quark family mixing

1973: Kobayashi and Maskawa, 3 quark family mixing

– – : GWS, the standard model

2004: PDG,  $|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 < 1??$



# Standard Model

$$\frac{g}{\sqrt{2}} W_{\alpha}^{+} (\bar{\mathbf{U}}_L \mathbf{V}_{CKM} \gamma^{\alpha} \mathbf{D}_L + \bar{e}_L \gamma^{\alpha} \nu_{eL} + \bar{\mu}_L \gamma^{\alpha} \nu_{\mu L} + \bar{\tau}_L \gamma^{\alpha} \nu_{\tau L}) + \text{h.c.}$$

There is just one gauge coupling  $g$ ,  $G_F = g^2 / (4\sqrt{2} M_W^2)$

The quarks are mixed by a unitary matrix  $\mathbf{V}$

Quarks and leptons have the same weak charge

At the opposite extreme:

$$g_q, g_e, g_{\mu}, g_{\tau}$$

$\mathbf{V}$  is not unitary

(But probabilities are conserved!)

$$\mathbf{U} = \begin{pmatrix} u \\ c \\ t \end{pmatrix}, \quad \mathbf{D} = \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$

$$\mathbf{V} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{ud} & V_{us} & V_{ub} \\ V_{ud} & V_{us} & V_{ub} \end{pmatrix}$$



# Decay Width

From  $\langle f | J_\alpha J^\alpha | i \rangle \Rightarrow$

$$\Gamma(|\Delta S|=1) = |V_{us}|^2 \times G^2 \times F_1(\text{masses}) \times F_2(\text{emi}) \times F_3(\text{si}).$$

1.  $F_1$ : kinematics, spinor algebra, phase space, OK
2.  $F_2$ : Cannot turn off electromagnetism. Lengthy but necessary. Few % correction, to 0.1% accuracy
3.  $F_3$ : Strong interactions are hardest to compute, can have large effects, choose best processes, lattice to the rescue!



# F<sub>2</sub> – Radiation must be included

$m(\gamma) = 0 \Rightarrow$  charge particles radiate.

Take for example  $K_S \rightarrow \pi^+ \pi^-$ . In the real world the em interaction

gives  $\Gamma_1 \propto \left| \text{diagram 1} + \text{diagram 2} \right|^2 \rightarrow \infty \text{ for } \omega \rightarrow 0.$

The infinity is cancelled by an opposite sign contribution:

$$\Gamma_2 \propto \left| \text{diagram 3} + \text{diagram 4} + \text{diagram 5} + \text{diagram 6} \right|^2.$$

$\Gamma_1 + \Gamma_2$  is finite but contains a correction of  $\mathcal{O}\left(\frac{\alpha}{\pi} \times \log \frac{\omega_0}{m}\right)$ .  
 $\omega_0$  is finite and experimental acceptance dependent.



## Radiation inclusive branching ratios

Inclusion (or exclusion) of photons does affect event acceptance. It is modern practice to give results fully inclusive of radiation up to the kinematic limit. This requires correct accounting of radiation in the Monte Carlo detector simulation program.

In the KLOE MC simulation, Geanfi, radiation is included at the event generation level, event by event. In the following, even if not explicitly stated, BRs are totally inclusive of radiation.

$BR(K \rightarrow f(\gamma))$  stands for  $BR(K \rightarrow f, f + \gamma, 0 < \omega < \omega_{\max})$

(I will also just use  $BR(K \rightarrow f)$ )



## $F_3 - SI - \text{Hadrons vs Quarks}$

Nuclear  $\beta$ -decay. Look at  $0^+ \rightarrow 0^+$  transitions. **No axial current.**  
**CVC helps with nuclear matrix elements.**

$|\Delta S| = 1$  decays. Corrections to

$$\langle f | \bar{u} \gamma_\alpha s | i \rangle$$

appear only at 2<sup>nd</sup> order in  $m_s - m_d$ . A-G theorem.

Kaon semileptonic decays.  $K \rightarrow \pi \ell \nu \Rightarrow 0^- \rightarrow 0^-$ .

Only vector part of current contributes to

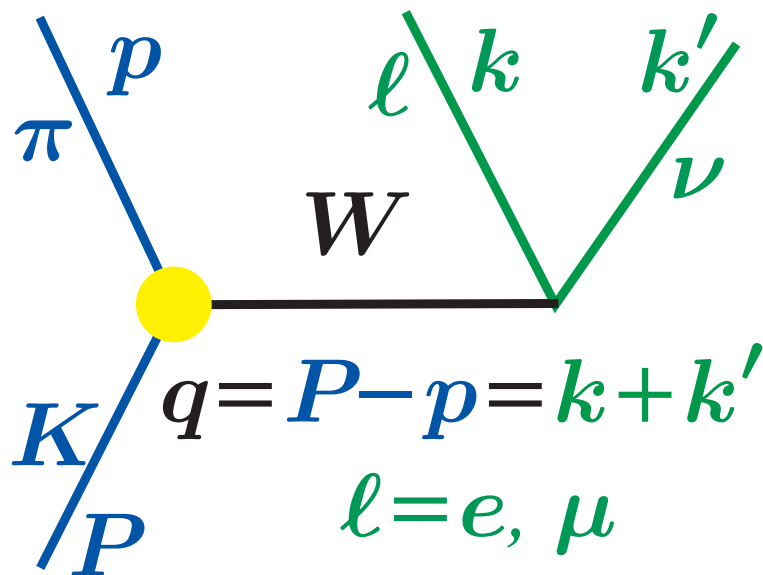
$$\langle \pi | \bar{u}_L \gamma_\alpha s_L | K \rangle$$

$$\langle \pi | \bar{u} \gamma_\alpha s | K \rangle = f_+(t)(P + p)_\alpha + f_-(t)(P - p)_\alpha$$





# Vector Current is protected



$$\langle \pi | J_\alpha | K \rangle = f_+(0) \times \left( (P + p)_\alpha \tilde{f}_+(t) + (P - p)_\alpha (\tilde{f}_0(t) - \tilde{f}_+(t)) \frac{\Delta}{t} \right)$$

$$\Delta = M^2 - m^2$$

Form factor because  $K$ ,  $\pi$  are not point like

Scalar and vector FFs equal at  $t=0$  (by construction)

Factor out  $f_+(0)$ .  $f_+(0) < 1$  because  $K \neq \pi$

$f_+(0) - 1 \ll 1$  (0.04) – small SU(3) breaking

$$\tilde{f}_+(0) = \tilde{f}_0(0) = 1$$



# Decay Width

$$\begin{aligned}
 & \Downarrow \\
 \Gamma(K_{\ell 3}(\gamma)) &= \frac{C_K^2 G_F^2 M_K^5}{768 \pi^3} S_{EW} |V_{us}|^2 f_+(0)^2 I_{K\ell} \\
 & \quad \color{red}{\Downarrow} \quad \color{red}{\Downarrow} \quad \color{red}{\Downarrow} \quad \color{red}{\Downarrow} \\
 & \quad \color{red}{192} \quad \color{purple}{128} \quad \times \left( 1 + \delta_K^{SU(2)} + \delta_{K\ell}^{EM} \right) \\
 & \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \color{red}{\Uparrow} \quad \color{red}{\Uparrow}
 \end{aligned}$$

$C_K^2 = 1(1/2)$  for  $K^0(K^\pm)$ ;  $S_{EW} = 1.0232(3)$ , EW corr, Sirlin

$\Downarrow$  Measure:  $\tau$ , BR,  $\tilde{f}_+$ ,  $\tilde{f}_0$   $\Downarrow\Uparrow$  Compute:  $\chi$ PT, Quark M, Lattice

Define

$f_+(0)$  as

$f_+^{K^0 \rightarrow \pi^\pm}(0)$

	$\delta_K^{SU(2)}(\%)$	$\delta_{K\ell}^{EM}(\%)$
$K_{e3}^0$	0	+0.57(15)
$K_{e3}^+$	2.36(22)	+0.08(15)
$K_{\mu 3}^0$	0	+0.80(15)
$K_{\mu 3}^+$	2.36(22)	+0.05(15)



# Phase space integral, $K_{e3}$

$$\rho(E_e, E_\nu) \propto \sum_{\text{spins}} |\mathfrak{M}|^2 \propto G^2 C_K^2 M^4 \left[ \frac{E_e}{M} \frac{E_\nu}{M} + \frac{\vec{k} \cdot \vec{k}'}{M^2} \right]$$

$$I_{K\ell} \propto \iint \rho(E_e, E_\nu) dE_e dE_\nu$$

$$\Gamma = \frac{G^2 C_K^2 M_K^5}{768 \pi^3} \underbrace{\int_{z_m}^{z_M} dz f(t)^2 \int_{x_m(z)}^{x_M(z)} 24 ((z+x-1)(1-x) - \alpha) dx}_{I_{K\ell}}$$

With 768 in the denominator,  $f(t)=1$ , all final state masses zero:

$$I_{K\ell} = 1.0$$

768 corresponds to 192 in muon decay rate.  $I_{K\ell}$  is dimensionless.

$$x = \frac{2E_e}{M}, \quad y = \frac{2E_\nu}{M}, \quad z = \frac{2E_\pi}{M}, \quad \alpha = \frac{m}{M}$$

End of introduction



# DAPHNE

An  $e^+e^-$  collider operated (mostly) at  $W = M_\phi$

$$e^+e^- \rightarrow \phi \rightarrow \begin{cases} K_S + K_L & 34.0\% \\ K^+ + K^- & 49.3\% \end{cases}$$

$\Rightarrow$  Tagged, monochromatic, pure  $K_S$ ,  $K_L$ ,  $K^+$ ,  
 $K^-$  beams

$$\text{BR}(K_S \rightarrow \pi \ell \nu) = (7.04 + 4.69) \times 10^{-4} \quad \text{“difficult”}$$

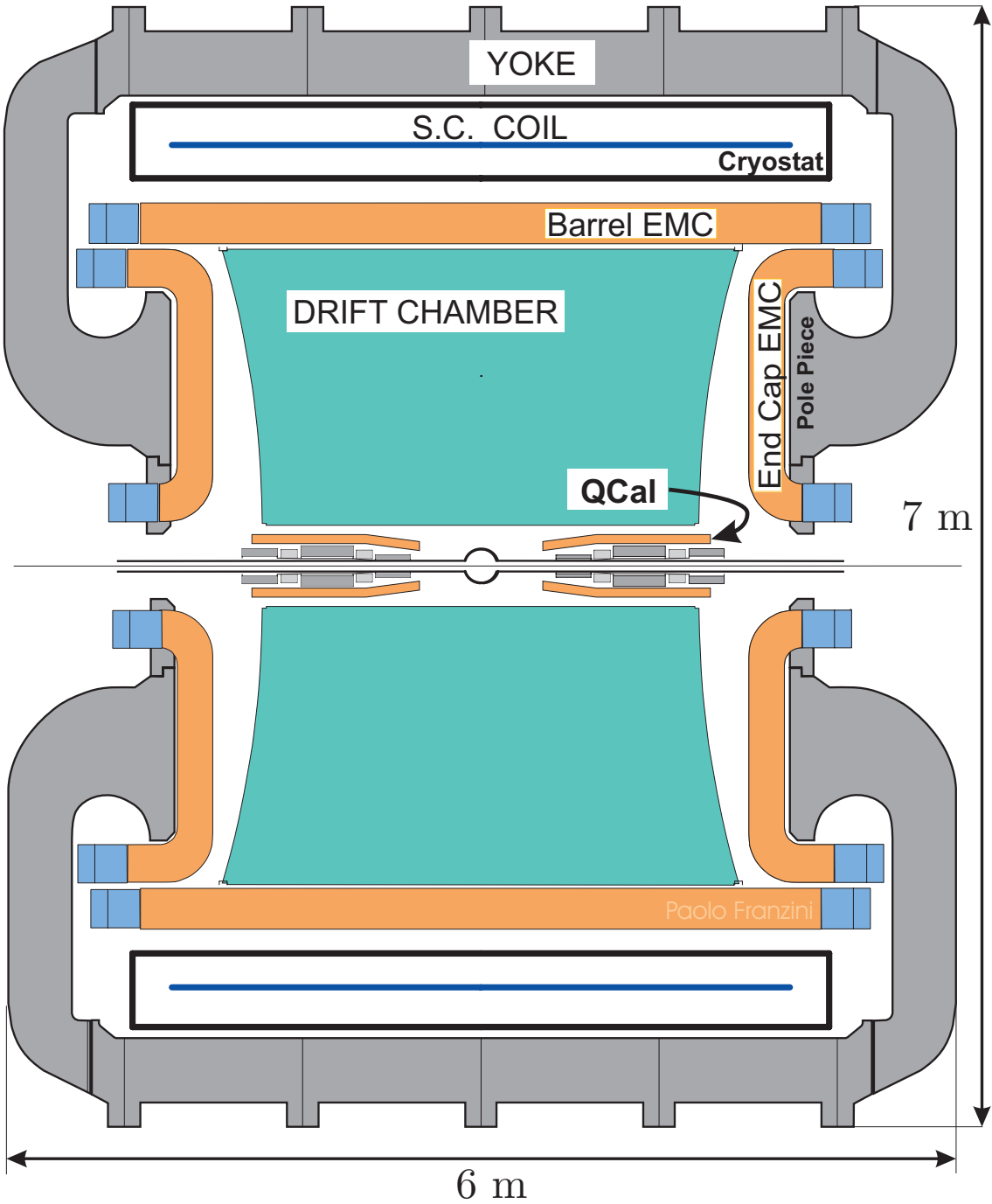
$$\text{BR}(K^\pm \rightarrow \pi \ell \nu) = (4.98 + 3.32) \times 10^{-2} \quad \text{“so so”}$$

$$\text{BR}(K_L \rightarrow \pi \ell \nu) = (4.06 + 2.71) \times 10^{-1} \quad \text{“best”}$$

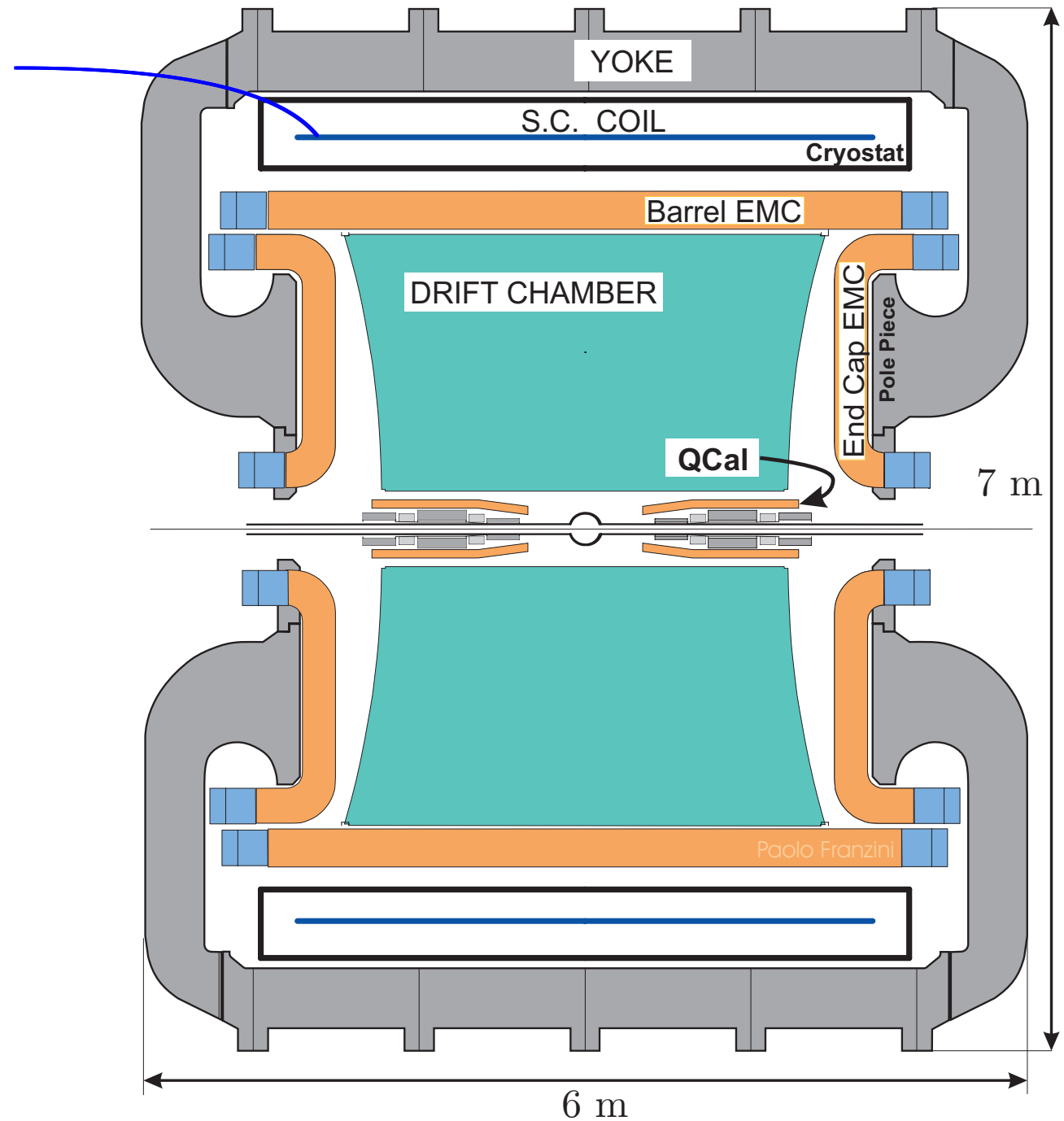
$e \quad \mu$



# KLOE

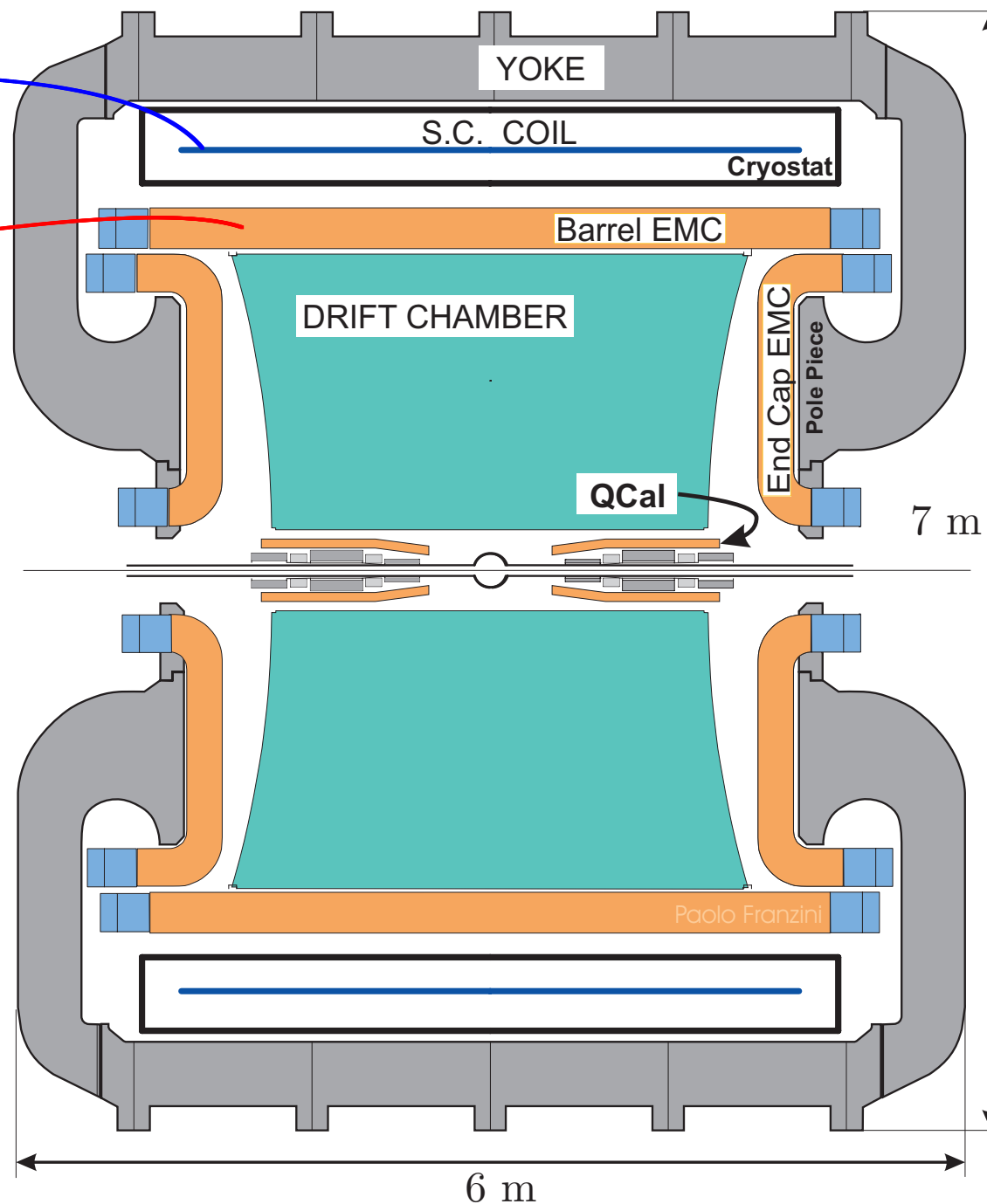


Magnet  
SC Coil,  $B=0.6$  T



Magnet  
SC Coil,  $B=0.6$  T

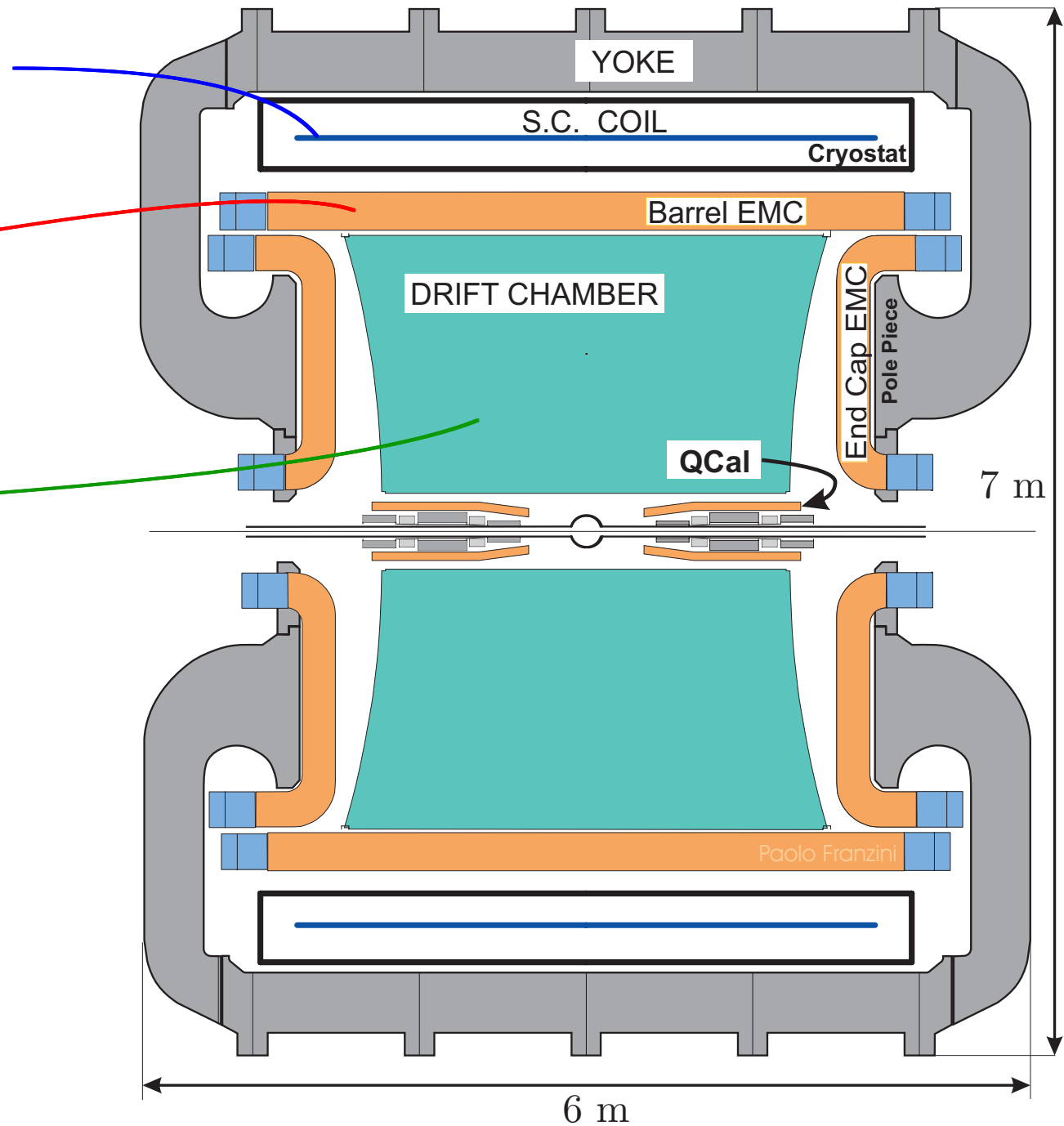
EM Calor.  
Pb-scint fiber  
4880 pm, 2440 cells



Magnet  
SC Coil,  $B=0.6$  T

EM Calor.  
Pb-scint fiber  
4880 pm, 2440 cells

Drift Ch.  
12582 sense wires  
52140 tot wires  
Carbon fiber walls



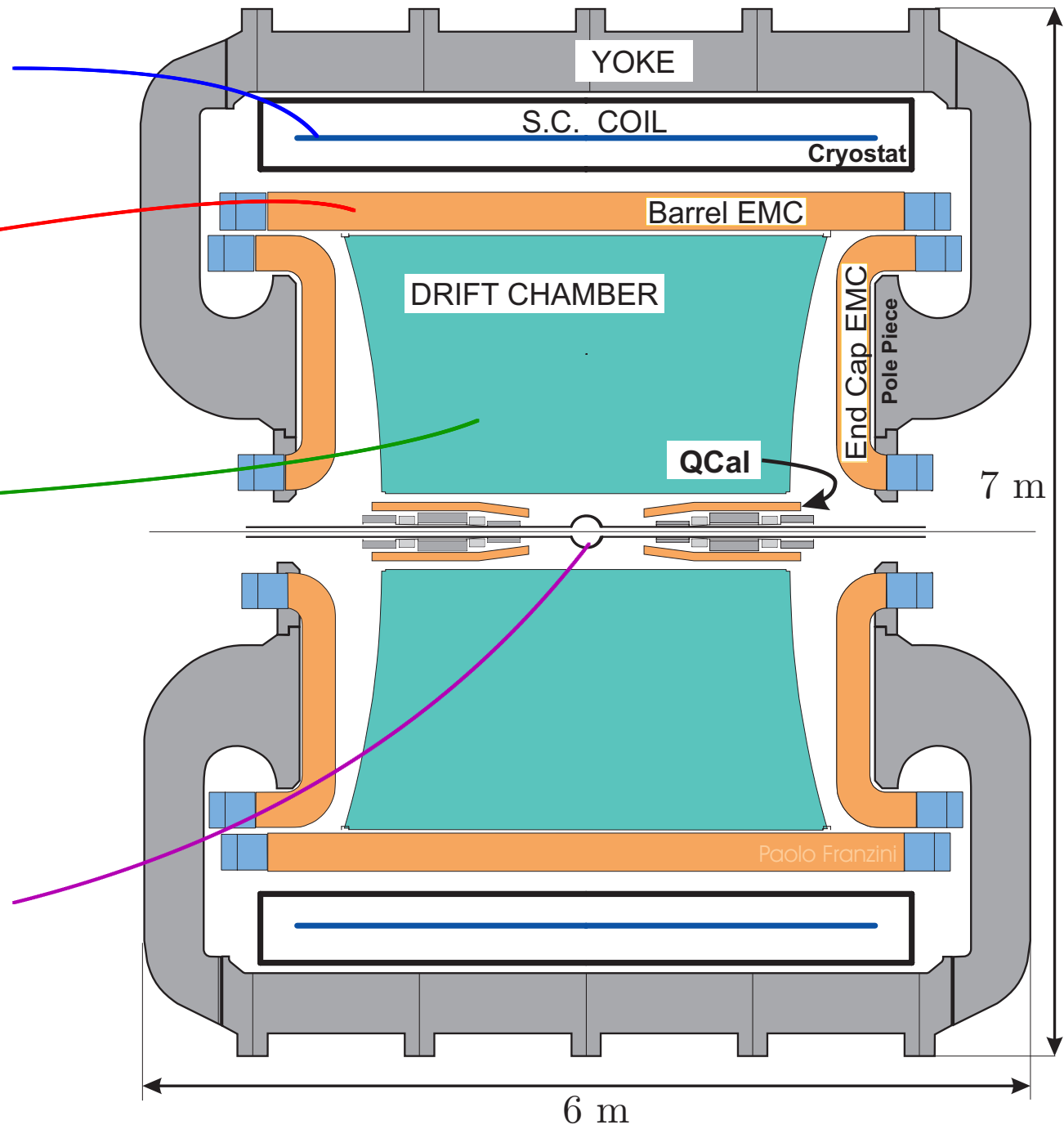


Magnet  
SC Coil,  $B=0.6$  T

EM Calor.  
Pb-scint fiber  
4880 pm, 2440 cells

Drift Ch.  
12582 sense wires  
52140 tot wires  
Carbon fiber walls

Al-Be beam pipe  
 $r=10$  cm, 0.5 mm  
thick

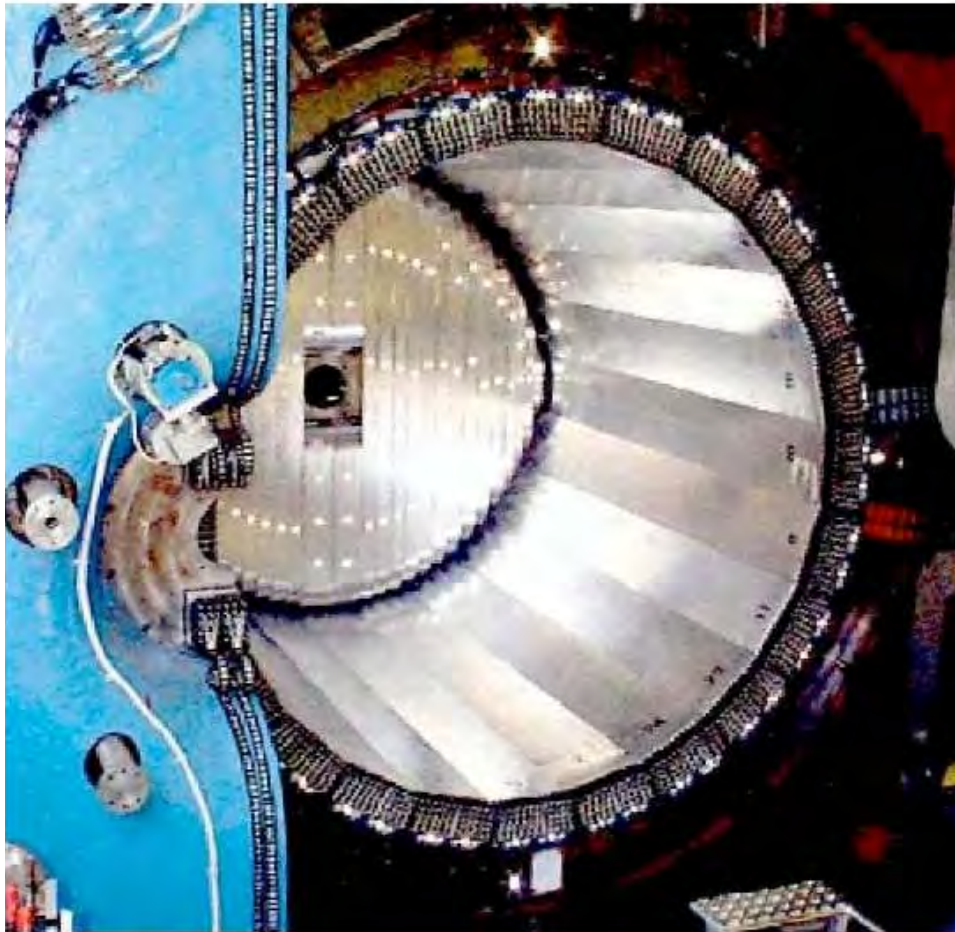


$$\sigma_E/E = 5.7\%/\sqrt{E} \text{ (GeV)}$$

$$\sigma_t = 54/\sqrt{E} \text{ (GeV)}$$

$$\oplus 100\text{-}150 \text{ ps}$$

## Drift chamber



## EM Calorimeter

$$\sigma(p_{\perp})/p_{\perp} = 0.4\%$$

$$\sigma_{x,y} = 150 \mu\text{m}; \sigma_z = 2 \text{ mm}$$



# KLOE precision measurements

$$M(K^0) = 497,583 \pm 5 \pm 20 \text{ keV}$$

$$M(\eta) = 547,874 \pm 7 \pm 29 \text{ keV}$$

Using  $M(\phi)$  measurements based on electron anomaly  
 $a_e = 0.0011596521859 \pm 0.00000000000038$ , Novosibirsk

## KLOE calibration

KLOE is continuously calibrated with data

Time, energy and position scales are calibrated every 1-2 hours

Machine energy, collision point,  $p(\phi)$  in lab

Together with  $\mathcal{L}$  and  $\int \mathcal{L} dt$ , all provided to DAΦNE



# Analysis

Prefilter. Reconstruction  $\sim 2$  hours after data taking

Initial classification

Massive Monte Carlo production. MC includes run by run  
background simulation

Tracking efficiency measured from data

Photon efficiency measured from data

Signal shapes and backgrounds from control samples

Continuous improvement of MC with feedback from  
physics analysis

Particle ID by time of flight

Particle ID by shape of energy release in EMC cells



# Measuring BRs

We only measure “absolute” BR, *i.e.*  $\Gamma_i/\Gamma$ , not ratios ( $\Gamma_i/\Gamma_j$ ) of partial rates. We have three ways to do it:

1. Measure single BR's
2. Measure (almost) all BR
3. Measure all partial rates:  $\Rightarrow$  BRs &  $\Gamma = 1/\tau$

We only use tagged kaons. Trigger efficiency **must not depend** on decay mode of **tagged** kaon. For every event, we verify the trigger was due to the **tagging** kaon.

There are differences between  $K_S$ ,  $K_L$  and  $K^\pm$ :

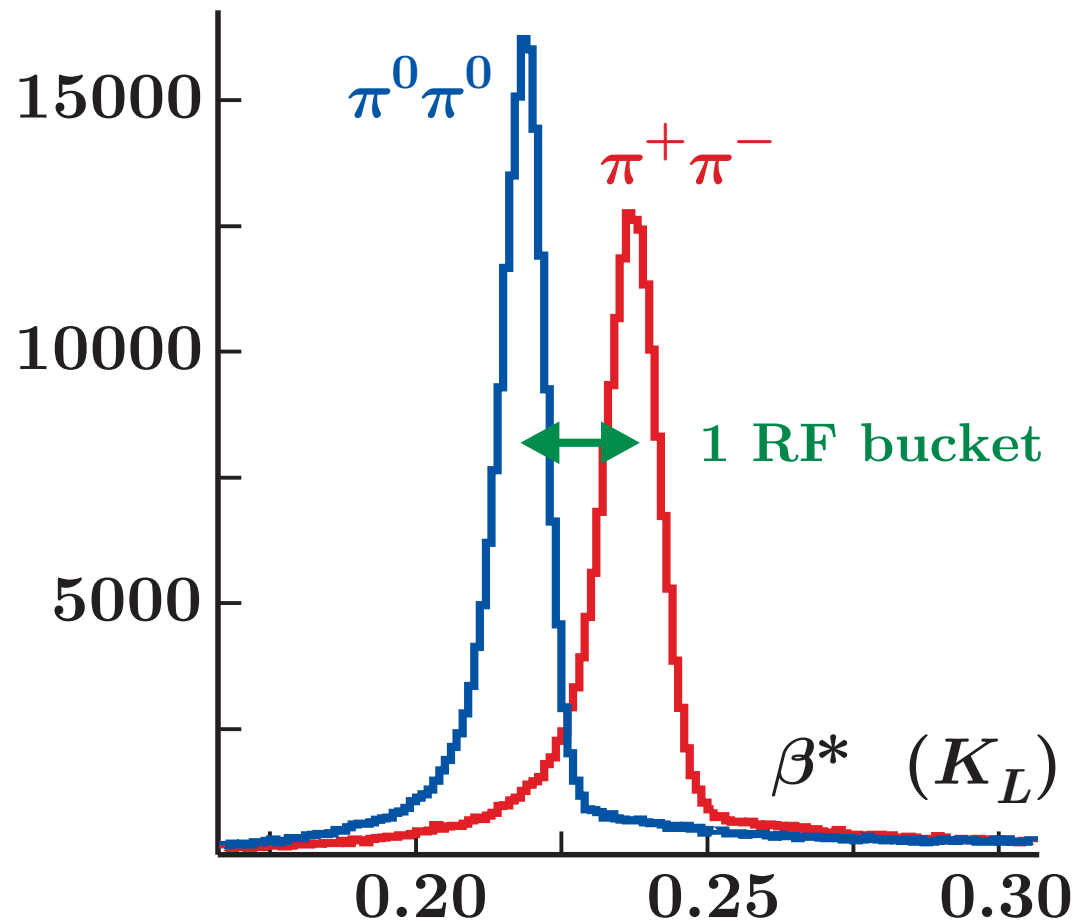
1. 1 vs 2 charged decay products
2. Tag by 1 or 2 body decays
3. Tag by time of flight



# The $K_S$ beam

First level, crude  $K_S$  tag by  $K_L$  TOF

$\sim 50\%$  of the  $K_L$ -mesons reach the calorimeter where most interact. Since  $\beta(K^0) \sim 0.2$ , we identify kaon interactions by TOF.



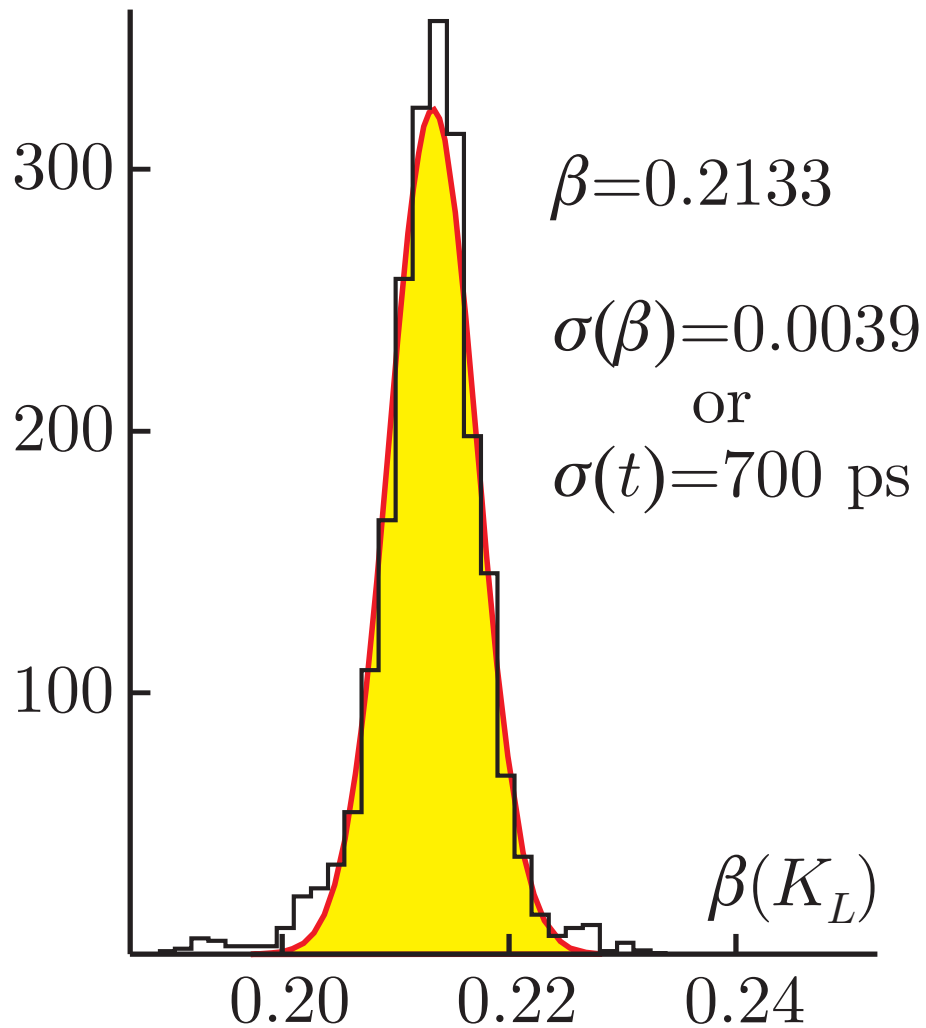
$K_L$  interacting in the calorimeter give an ideal  $K_S$  tag, almost independent of  $K_S$  decay mode. The value of  $T_0$  here is obtained from the first particle to reach the calorimeter.

$\pi, \mu, \gamma$  or  $e$





$$K_S \rightarrow \pi^\pm e^\mp \nu$$

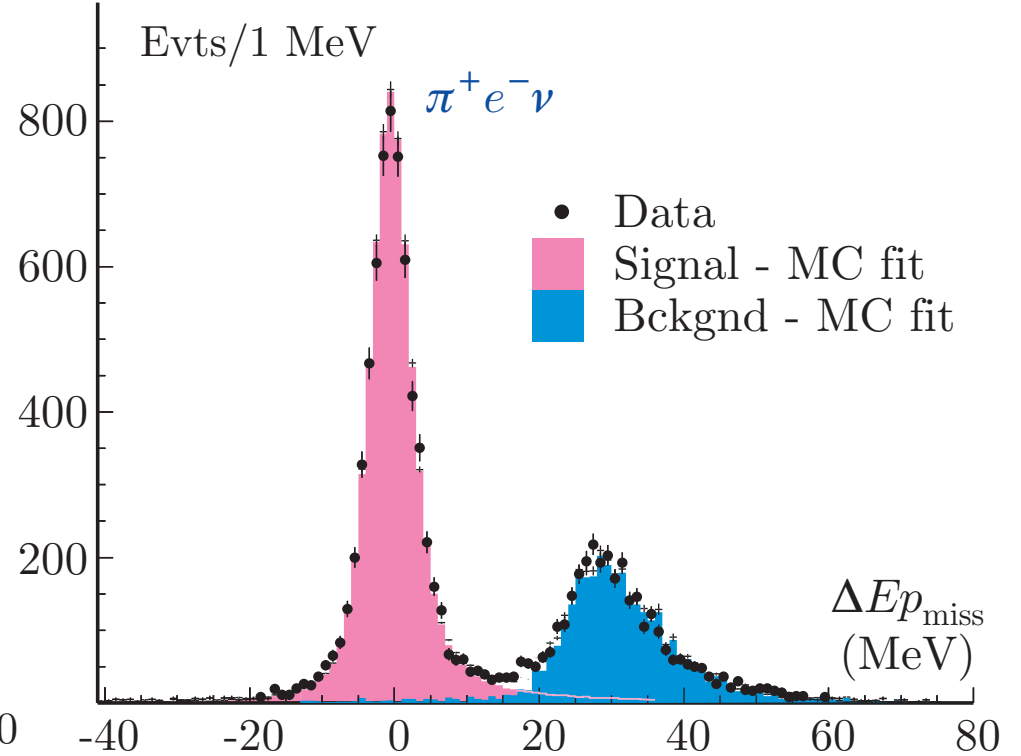
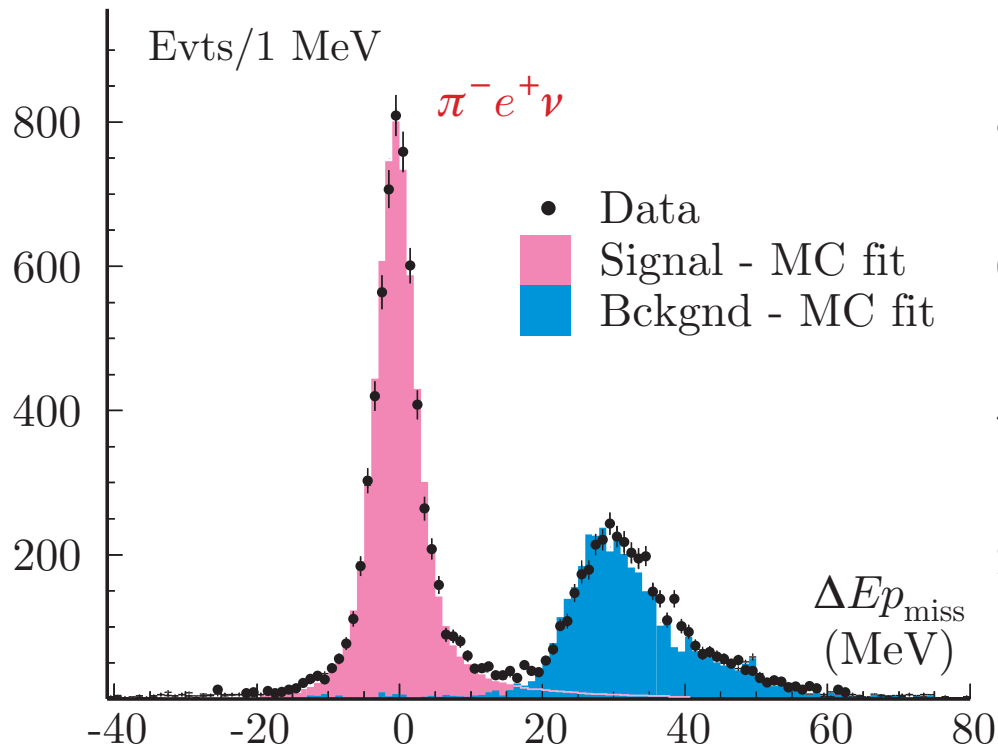


$\delta W(\text{DA}\Phi\text{NE}) = 1 \text{ MeV}$   
 gives  $\delta\beta = 0.004$

After event reconstruction very clean signal tags the presence of  $K_S$ . Search for two track decays near IP. Reject  $K_S \rightarrow \pi^+ \pi^-$ . Use TOF difference for two tracks taken as  $\pi - e$  or  $e - \pi$ . Provides electron ID and charge of lepton. Use  $|E_{\text{miss}} - p_{\text{miss}}|$  to isolate  $K_S, e3$  events.



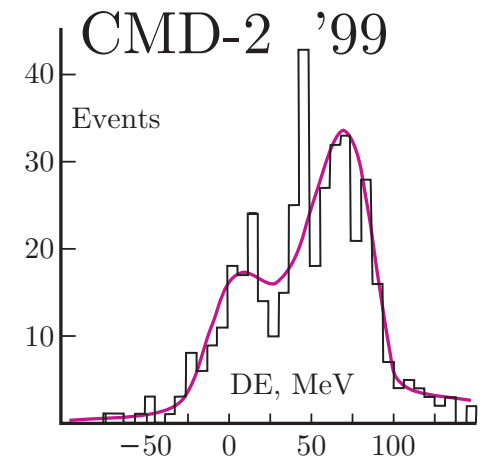
# K<sub>S</sub> → πeν, 2005



~13,000 signal events

$$\text{BR}(K_S \rightarrow \pi e \nu) = (7.046 \pm 0.091) \times 10^{-4}$$

|V<sub>us</sub>| to 0.7%





# Dominant $K_L$ decay modes

We know (KLOE, PDG, etc)

$$\begin{aligned}\Sigma_B &= \text{BR}(K_{Le3}) + \text{BR}(K_{L\mu3}) + \text{BR}(K_L \rightarrow 3\pi^0) \\ &+ \text{BR}(K_L \rightarrow \pi^+\pi^-\pi^0) = 1 - 0.0034(\pm 0.00004)\end{aligned}$$

We assume

$$\tau(K_L) = 51.7 \text{ ns}$$

We measure:

$$\Sigma_B = 1.0104 - 0.0034$$

GOOD!

We can solve for :  $\begin{cases} 4 \text{ BRs} \\ K_L \text{ lifetime} \end{cases}$



# A step at a time

Detect  $K_S \rightarrow \pi^+ \pi^-$ ,  $\Rightarrow$  tagged  $K_L$  beam

Find: a)  $K_L \rightarrow 2$  charged particles: b)  $K_L \rightarrow 3\pi^0$

Classify a) as  $e\pi\nu$ ,  $\mu\pi\nu$ ,  $\pi^+ \pi^- \pi^0$ ,  $\pi^+ \pi^-$

Determine “tag bias” corrections

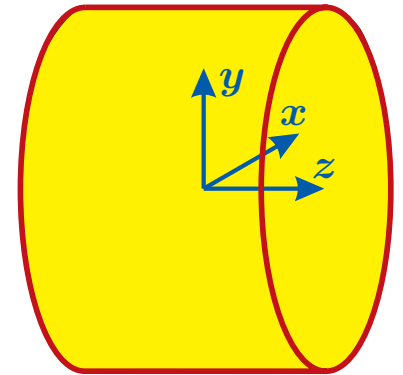
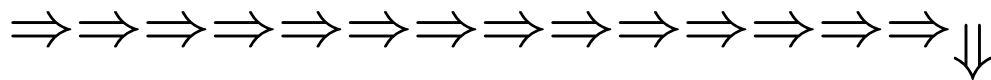
Determine FV acceptance & all other efficiencies

Count in FV. FV:  $35 < x, y < 150$  cm,  $|z| < 120$  cm

All the above requires knowledge of  $\tau$ . Use old value  $\tau^0 = 51.7$  ns.

In KLOE,

$$\mathcal{A}_{FV} = \mathcal{A}_{FV}^0 \times \left( 1 + 0.0128 \text{ ns}^{-1} (\tau^0 - \tau) \right)$$



# Solve for BRs and lifetime

$$\text{BR}(K_{Le3}) = 0.4009 \pm 0.0015$$

$$\text{BR}(K_{L\mu3}) = 0.2699 \pm 0.0014$$

$$\text{BR}(K_L \rightarrow 3\pi^0) = 0.1996 \pm 0.0020$$

$$\text{BR}(K_L \rightarrow \pi^+\pi^-) = 0.1261 \pm 0.0011$$

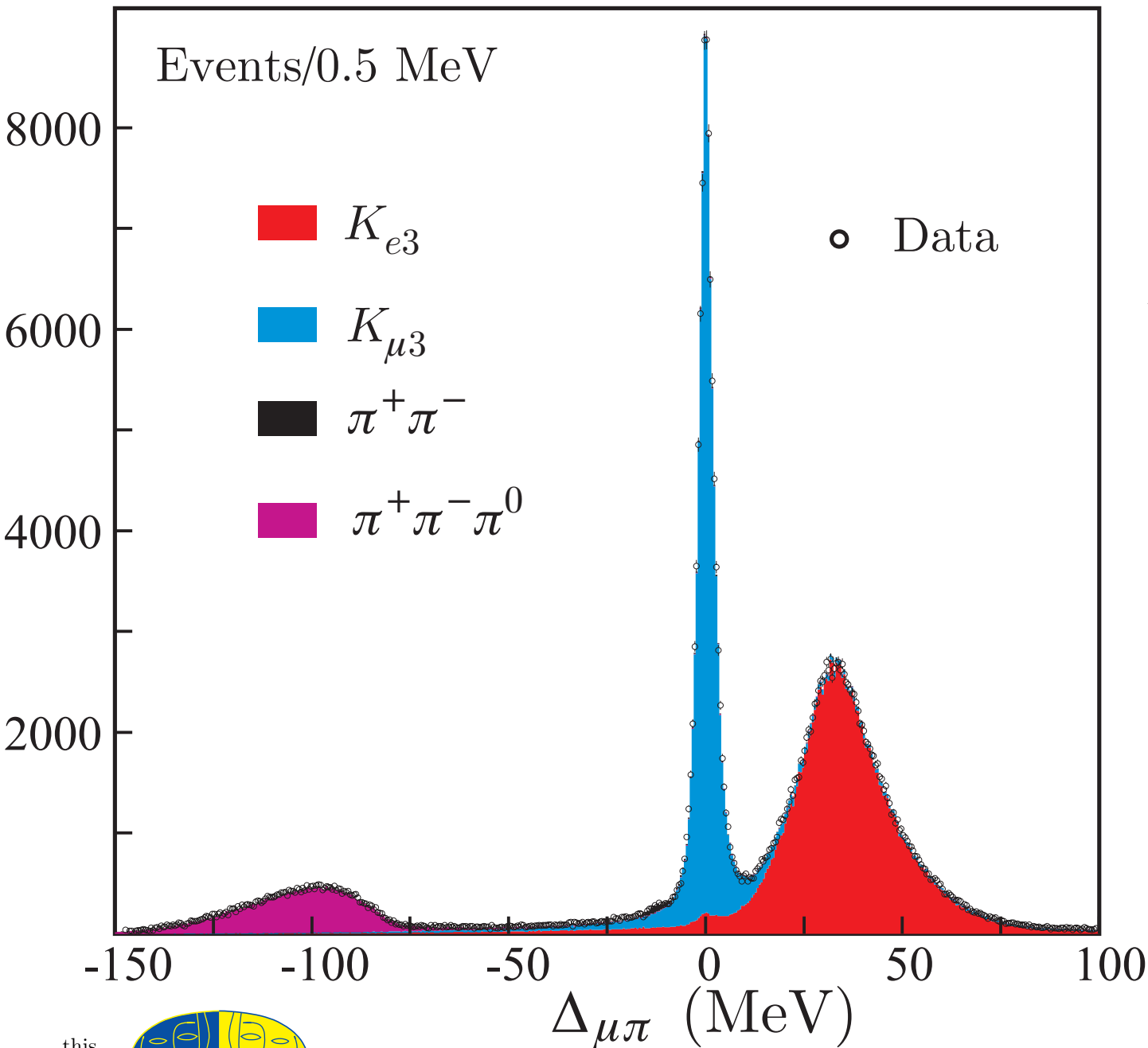
$$\tau(K_L) = 50.72 \pm 0.37 \text{ ns}$$

There are of course correlations

$$\begin{pmatrix} -0.25 & -0.56 & -0.07 & 0.25 \\ & -0.43 & -0.20 & 0.33 \\ & & -0.39 & -0.21 \\ & & & -0.38 \end{pmatrix}$$



# $K_L \rightarrow$ two track events



For  $K_L \rightarrow \pi^\pm \ell^\mp \nu$ :

$E_{\text{miss}} \neq 0$

$\mathbf{p}_{\text{miss}} \neq 0$

$m_\nu = 0 \rightarrow$

$\Delta = E_{\text{miss}} - |\mathbf{p}_{\text{miss}}| = 0$

$\Delta_{\mu\pi} =$  smallest of

$\Delta_{\mu^+ \pi^-}, \Delta_{\pi^+ \mu^-}$

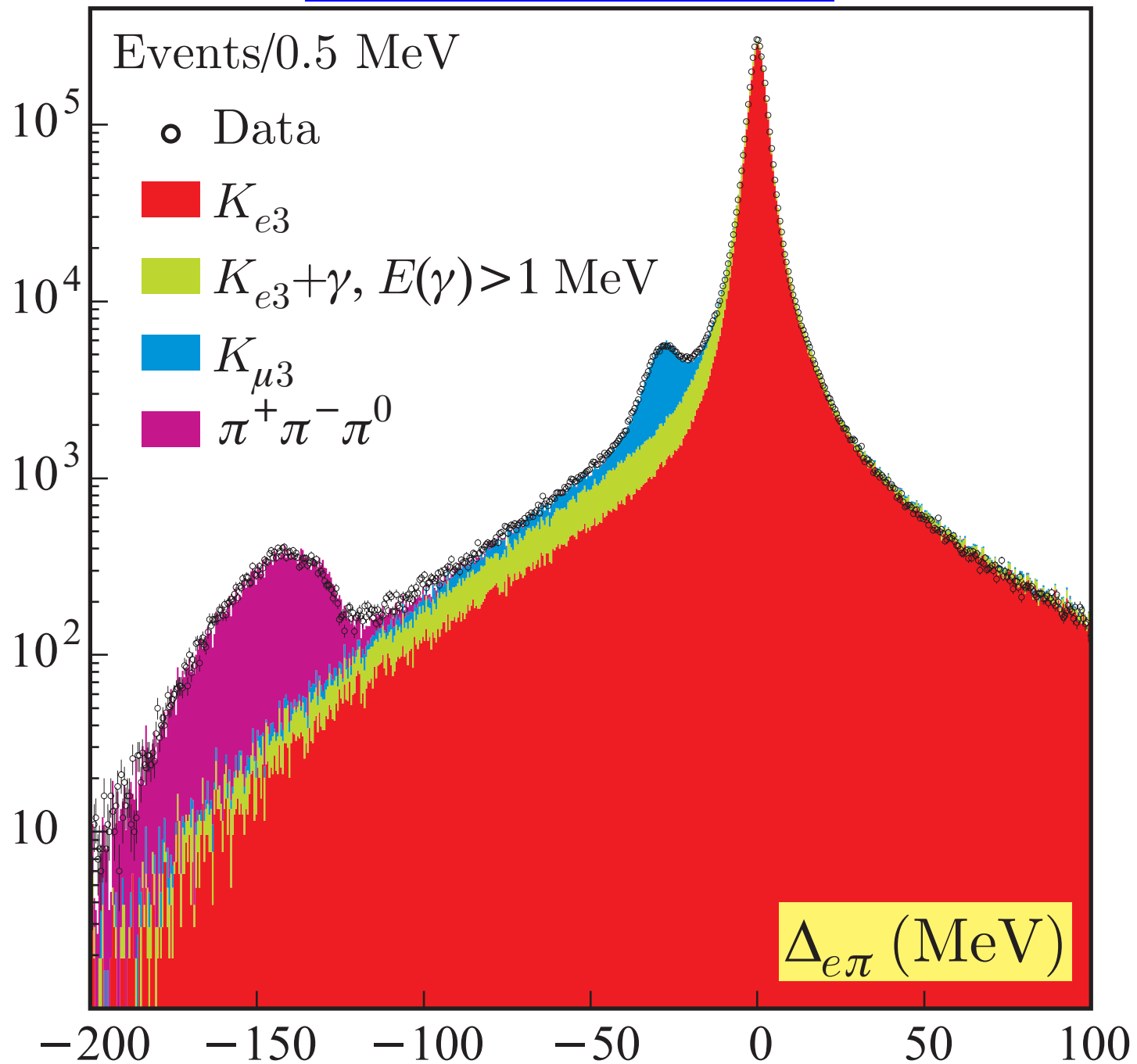
Fit to MC shapes

Count, BRs

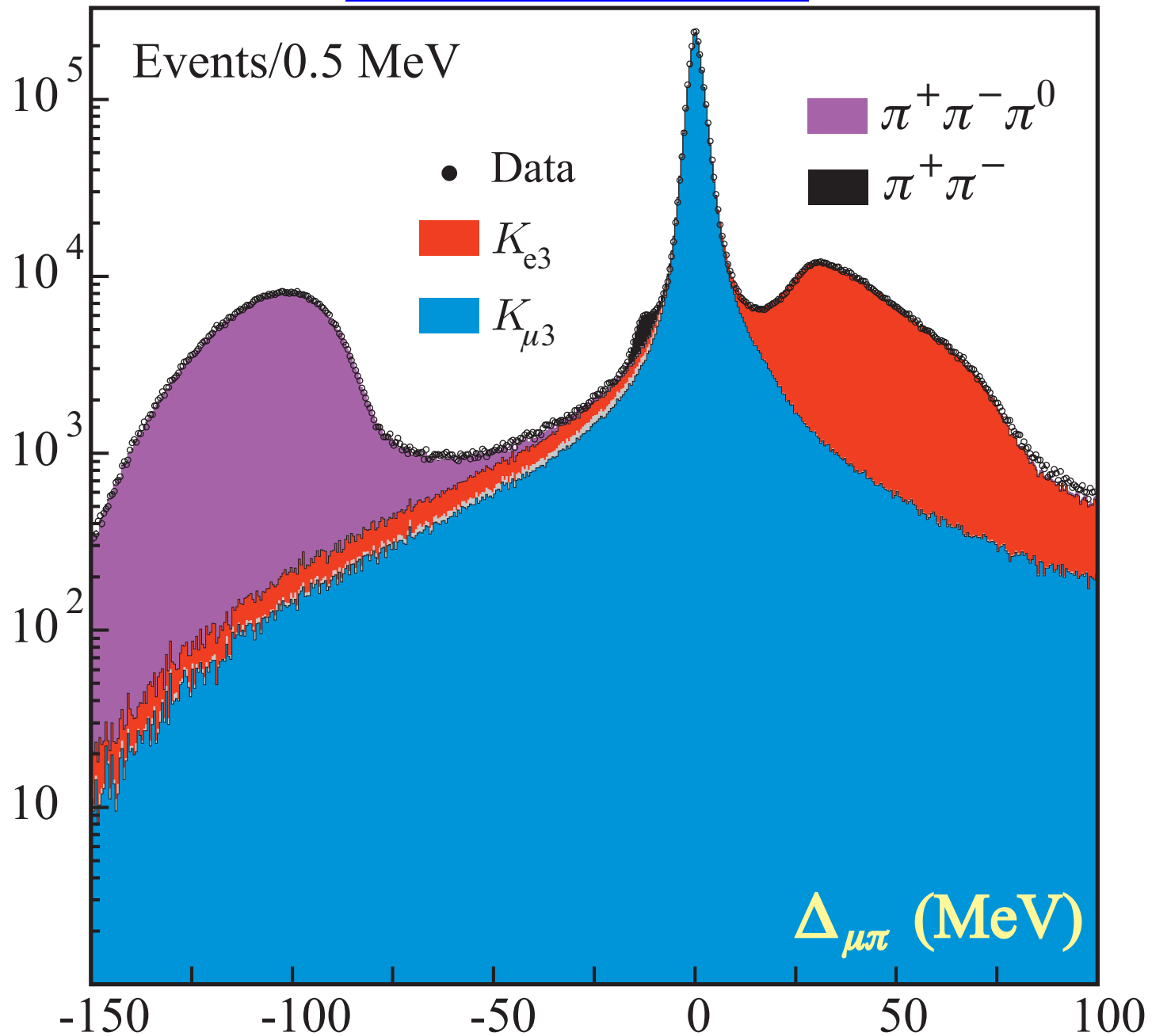
14 sub-samples



# Select electrons



# Select muons



$$K_L \rightarrow \pi^0 \pi^0 \pi^0$$

Count tagged  $K_L$  decaying to  $3\gamma$ ,  $4\gamma$ ,  $5\gamma$ ,  $6\gamma$ ,  $7\gamma$

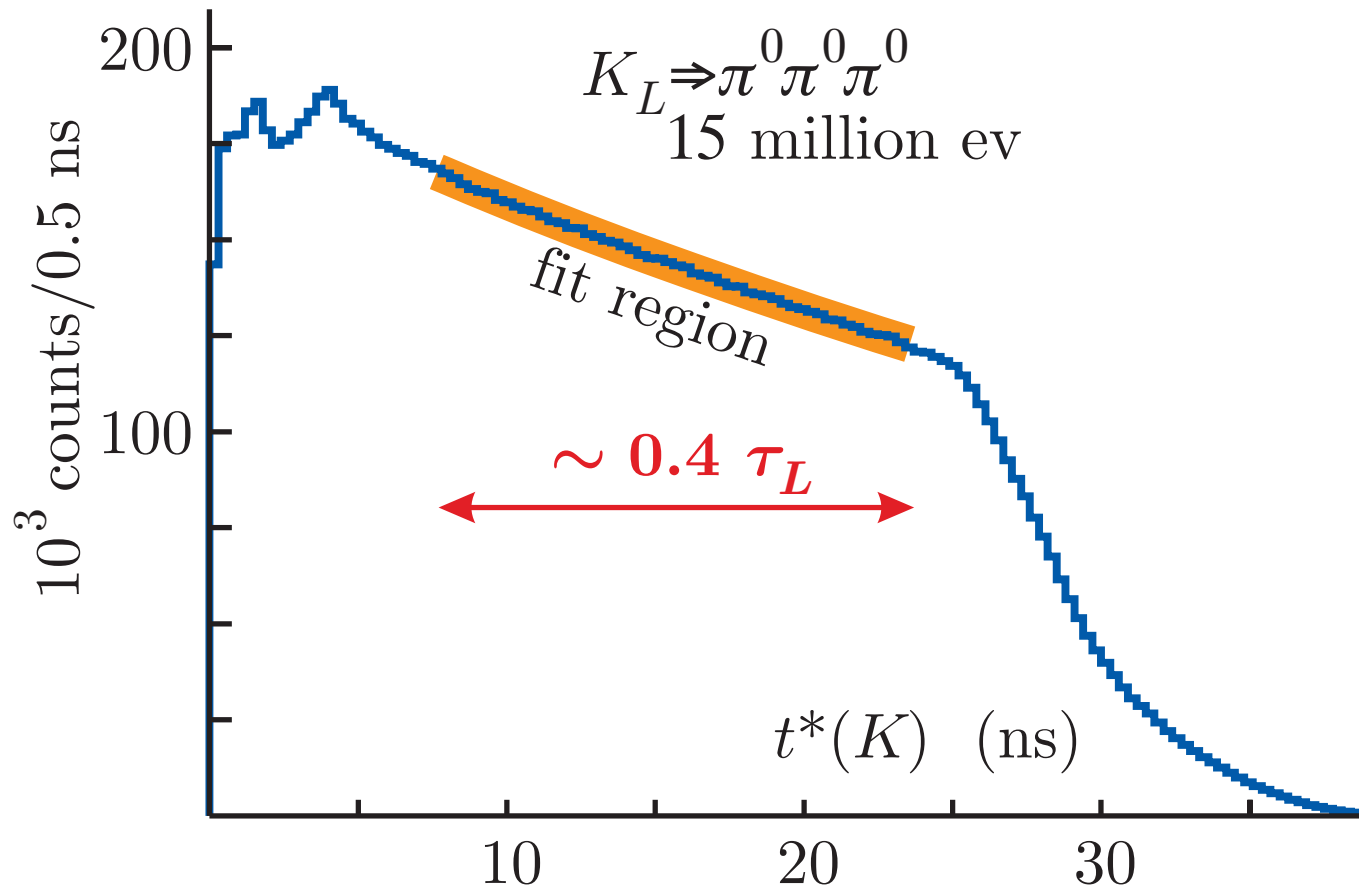
Efficiency  $\mathcal{O}(100\%)$

MC corrections

Background  $\mathcal{O}(0.001)$



# Time dependence of $K_L \rightarrow \pi^0 \pi^0 \pi^0$



$$\tau(K_L) = 50.92 \pm 0.30 \text{ ns}$$





# All KLOE $K_L$ results

Param.	Value	Correlation coefficients						
$BR(K_{Le3})$	0.4008(15)							
$BR(K_{L\mu3})$	0.2699(14)	-0.31						
$BR(3\pi^0)$	0.1996(20)	-0.55	-0.41					
$BR(\pi^+\pi^-)$	0.1261(11)	-0.01	-0.14	-0.47				
$BR(\pi^+\pi^-)$	$1.96(2) \times 10^{-3}$	-0.15	0.50	-0.21	-0.07			
$BR(\pi^0\pi^0)$	$8.49(9) \times 10^{-4}$	-0.15	0.48	-0.20	-0.07	0.97		
$BR(\gamma\gamma)$	$5.57(8) \times 10^{-4}$	-0.37	-0.28	0.68	-0.32	-0.14	-0.13	
$\tau_L$	50.84(23) ns	0.16	0.22	-0.14	-0.26	0.11	0.11	-0.09

We can do it!

When PDG does it, combining inconsistent results from different experiments it can be a disaster!!! Till '06 they could not input correlations.

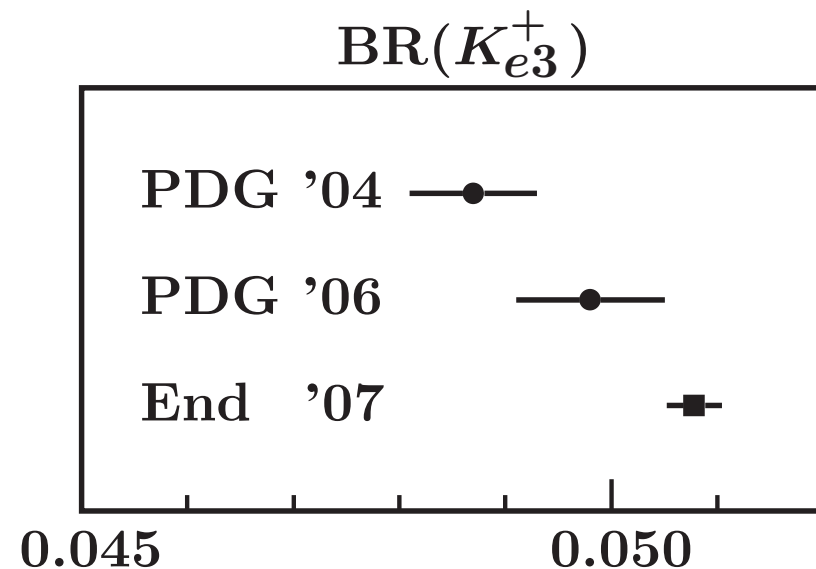
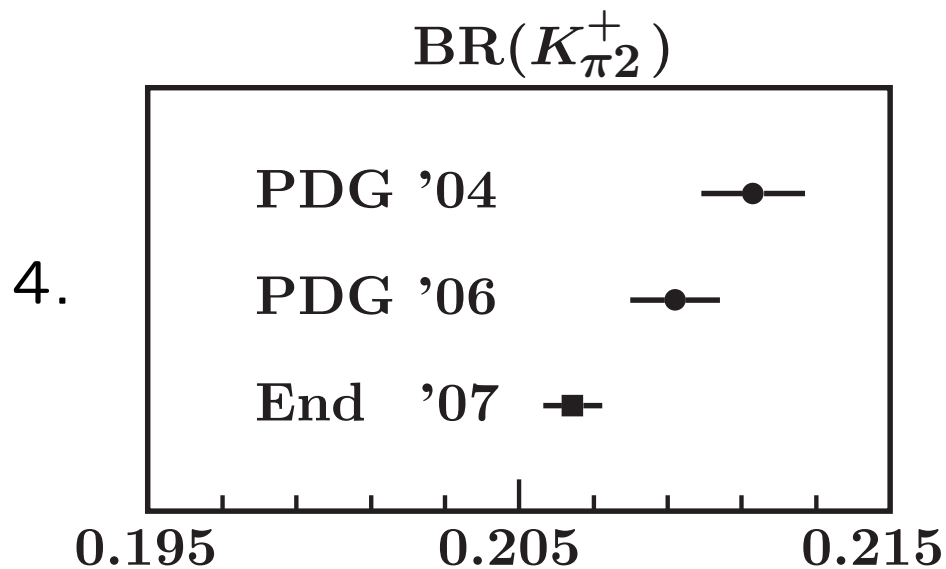


# PDG Example

1.  $BR(K \rightarrow 3\pi^0)$ . NA31 '95:  $0.2105 \pm 0.0028$  . PDG'04:  $0.2105 \pm 0.0021$ . '08:  $0.1951 \pm 0.0008$ .  $5.3\sigma$  to  $6.8\sigma$  off!

2.  $3\pi^0/K_{e3}$  NA31 '95 and PDG '04  $6.4\sigma$  to  $7.6\sigma$  off!

3.  $BR(K^\pm \rightarrow \pi^+\pi^0)$ . Chiang '72:  $0.2118 \pm 0.0028$ . PDG'04:  $0.2113 \pm 0.0014$ . '08:  $.2064 \pm 0.0008$ .  $1.9\sigma$  to  $3.2\sigma$  off!



## PDG masses

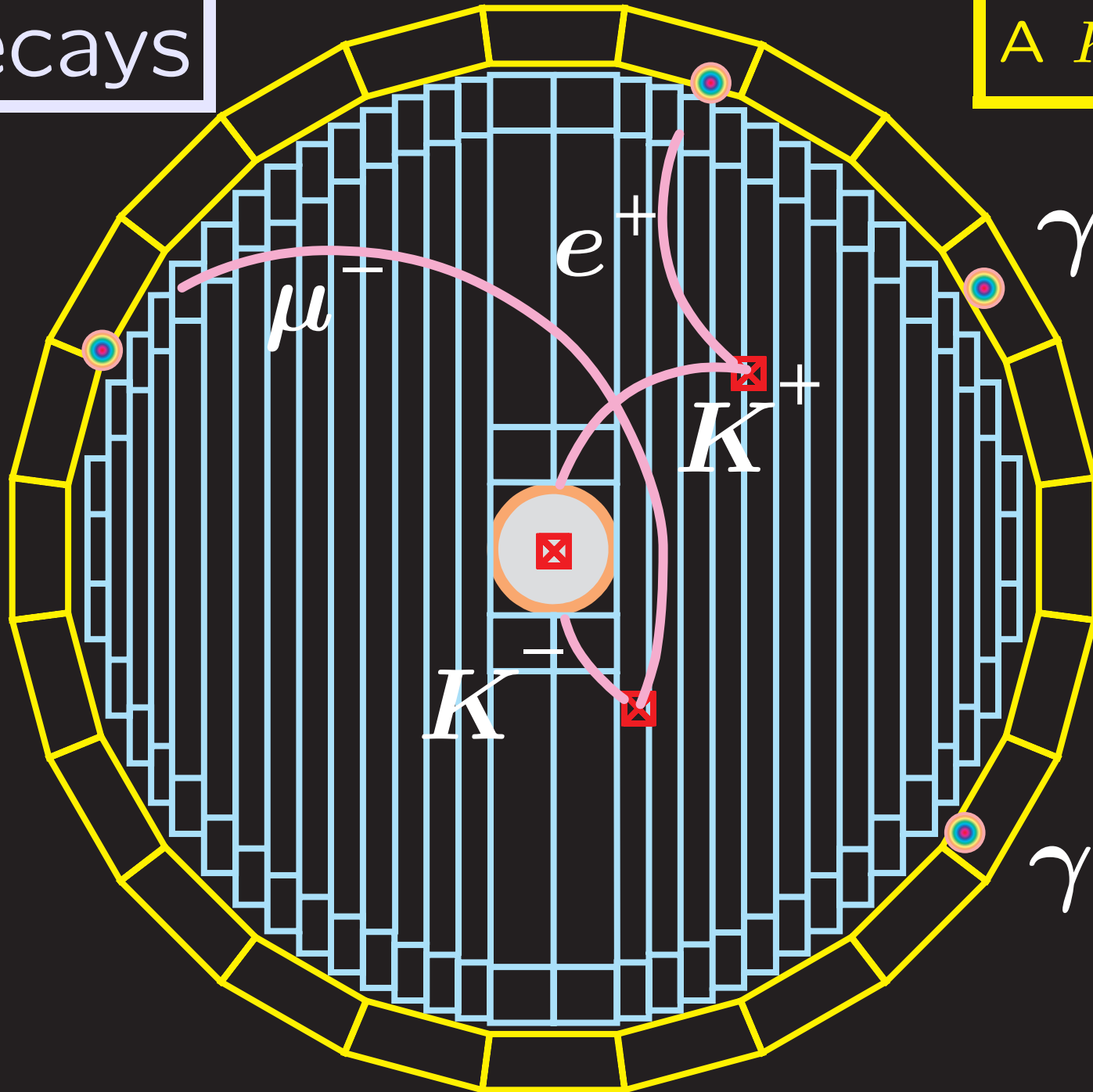
5.  $\eta$  mass. The PDG procedure should give  $m(\eta)=547.53\pm 0.26$  MeV,  $CL=10^{-15}$ . They include 3 old measurements which should give  $m(\eta)=547.51\pm 0.13$  MeV,  $CL=3 \times 10^{-13}$ . They give  $m(\eta)=547.51\pm 0.18$  MeV,  $CL=0.001$ ???

6.  $K^+$  mass. PDG procedure  $\Rightarrow m(K^+)=493.677\pm 0.012$  from 6 values. From only last two  $m(K^+)=493.686\pm 0.019$ . They give  $493.677\pm 0.013$ . Also CL is 0.00035 ( $2.1 \times 10^{-6}$ ), they give  $CL=0.001$ .

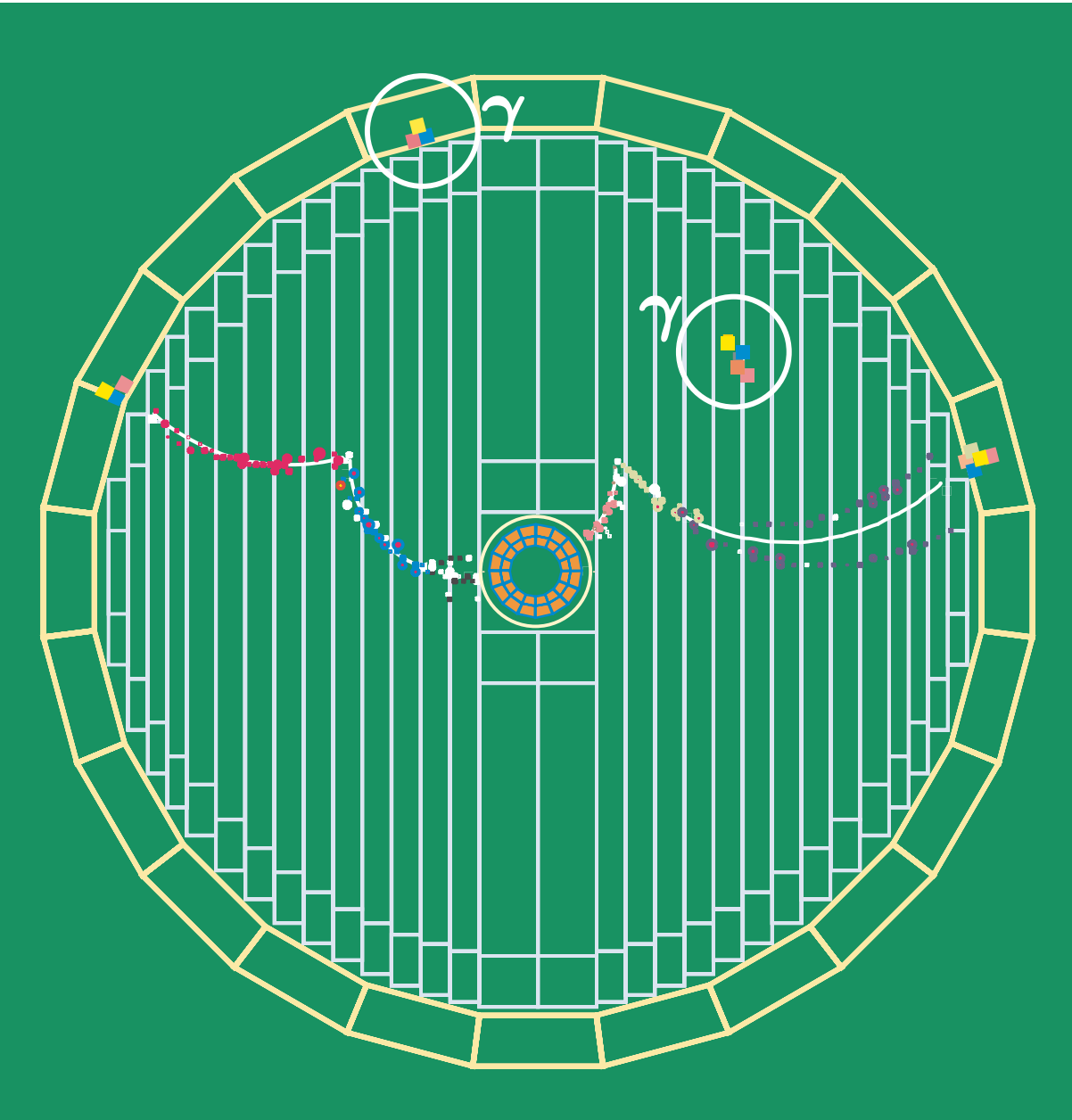


# $K^\pm$ decays

# A $K_{e3}^\pm$ decay



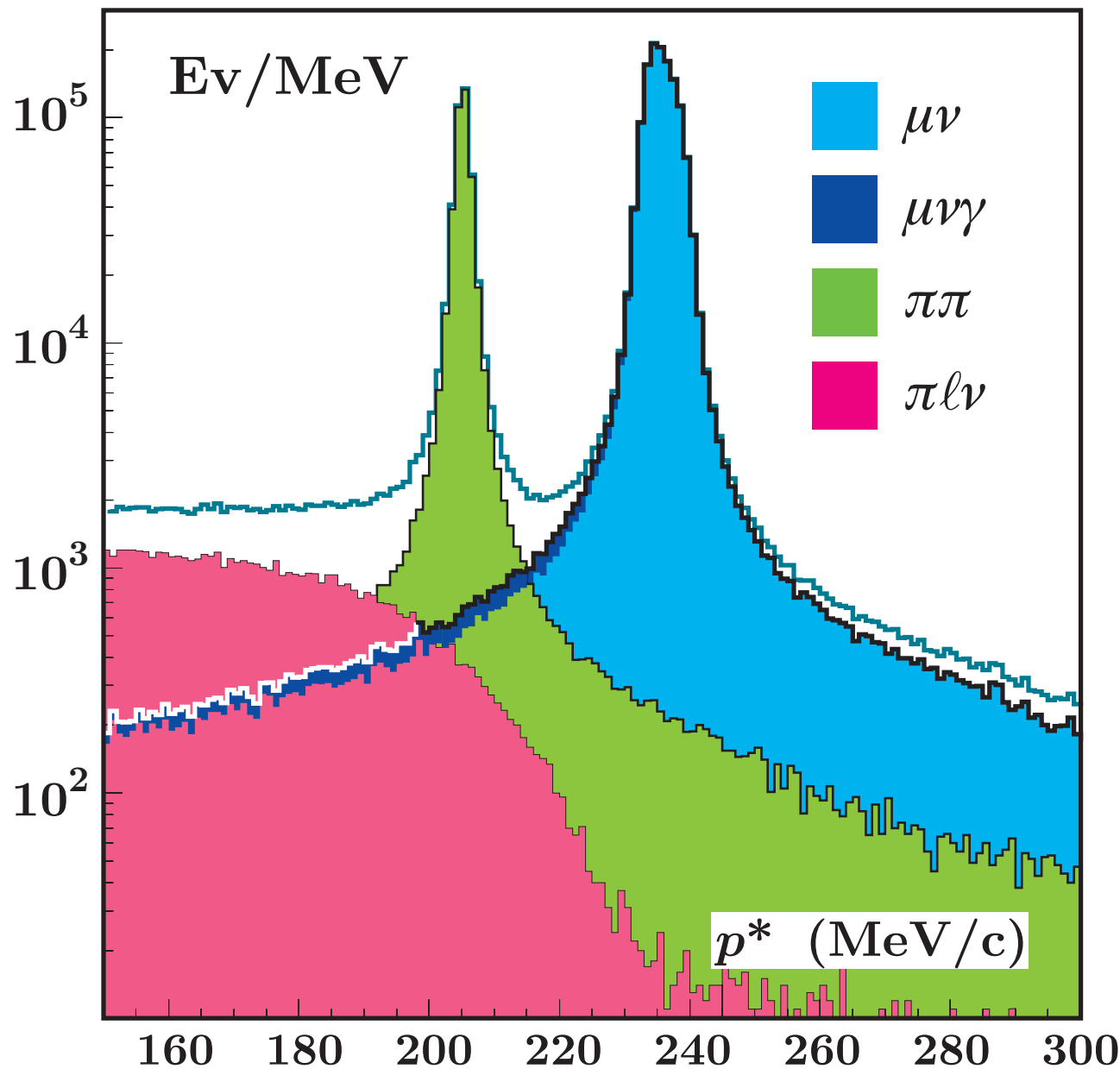
# $K_{\ell 3}^{\pm}$ decays



Tag	–	Decay
$K^+ \rightarrow \mu^+ \nu$	–	$K^- \rightarrow \pi^0 e^- \bar{\nu}$
$K^+ \rightarrow \pi^+ \pi^0$	–	$K^- \rightarrow \pi^0 \mu^- \bar{\nu}$
$K^- \rightarrow \mu^- \bar{\nu}$	–	$K^+ \rightarrow \pi^0 e^+ \nu$
$K^- \rightarrow \pi^- \pi^0$	–	$K^+ \rightarrow \pi^0 \mu^+ \nu$



# Tagging decays and $K^\pm \rightarrow \mu\nu$ decays



Compute CM momentum using  $m(\pi)$ . The  $\mu\nu$  decay peak is distorted but quite recognizable for tagging.

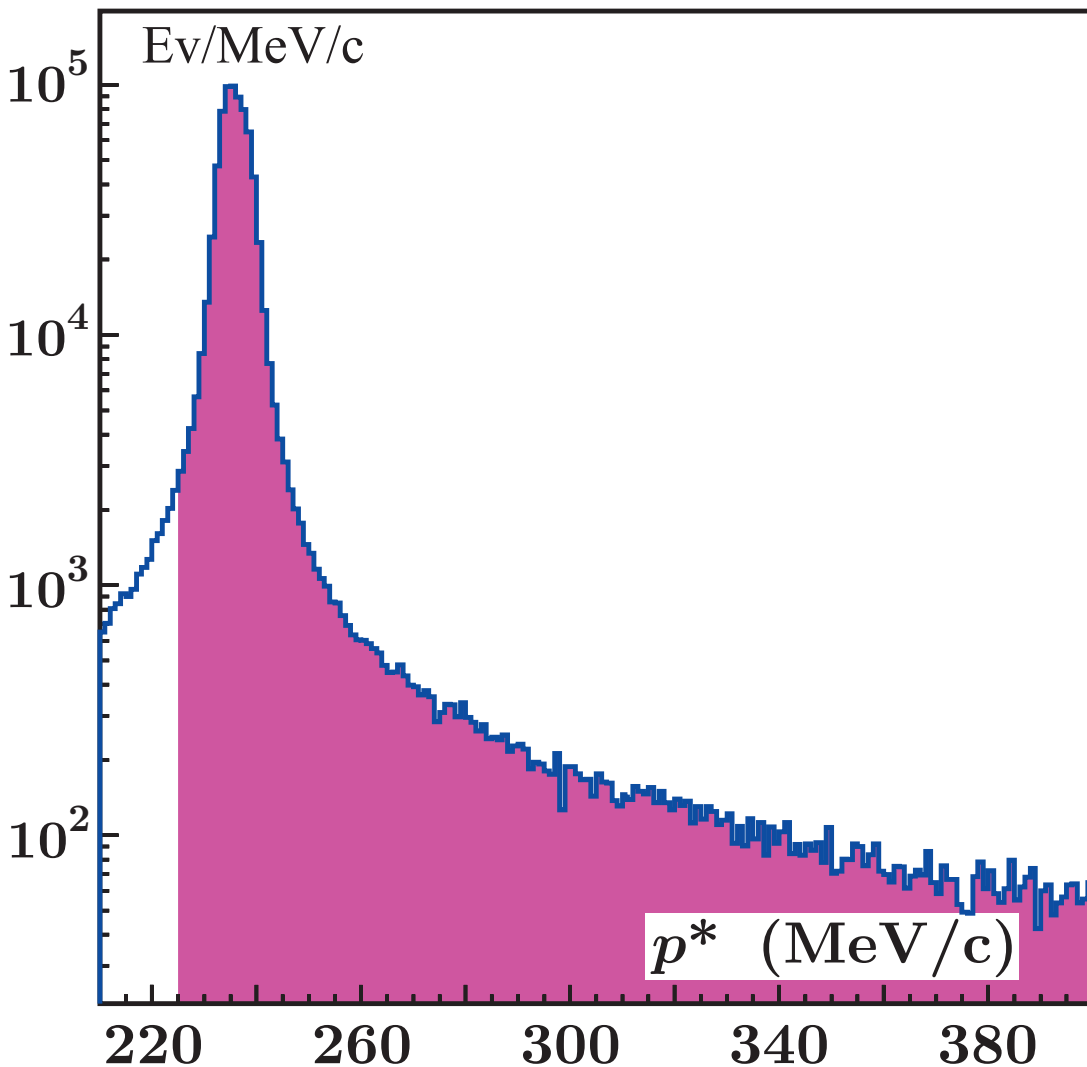
We also get

$\text{BR}(K \rightarrow \mu\nu(\gamma))$  and

$\text{BR}(K \rightarrow \pi\pi^0(\gamma))$

Note the radiative tail!



$$K \rightarrow \mu\nu(\gamma) \text{ BR}$$


All background decays, contain a  $\pi^0$ . The background shape is obtained from a control sample

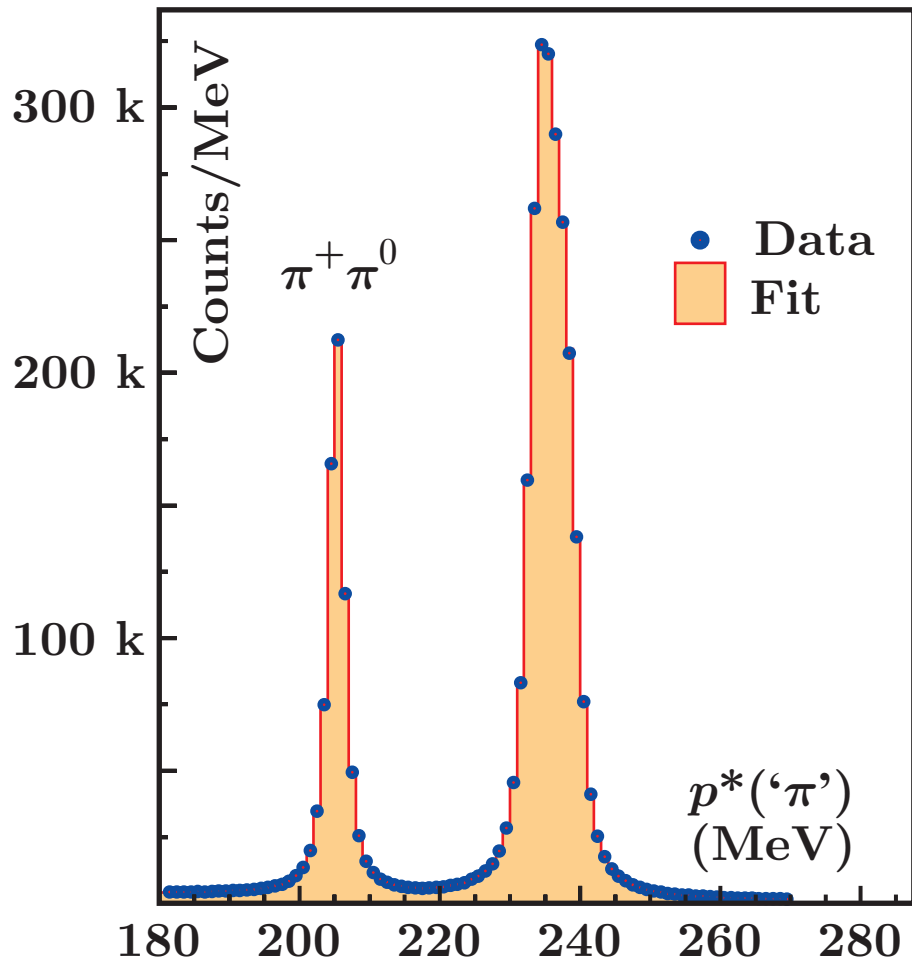
$K_{\mu 2}$  events are counted only in the shaded area, after background subtraction,  $\sim 2\%$ .

$$\text{BR}(K^+ \rightarrow \mu^+ \nu(\gamma)) = 0.6366 \pm 0.0009 \pm 0.0015.$$

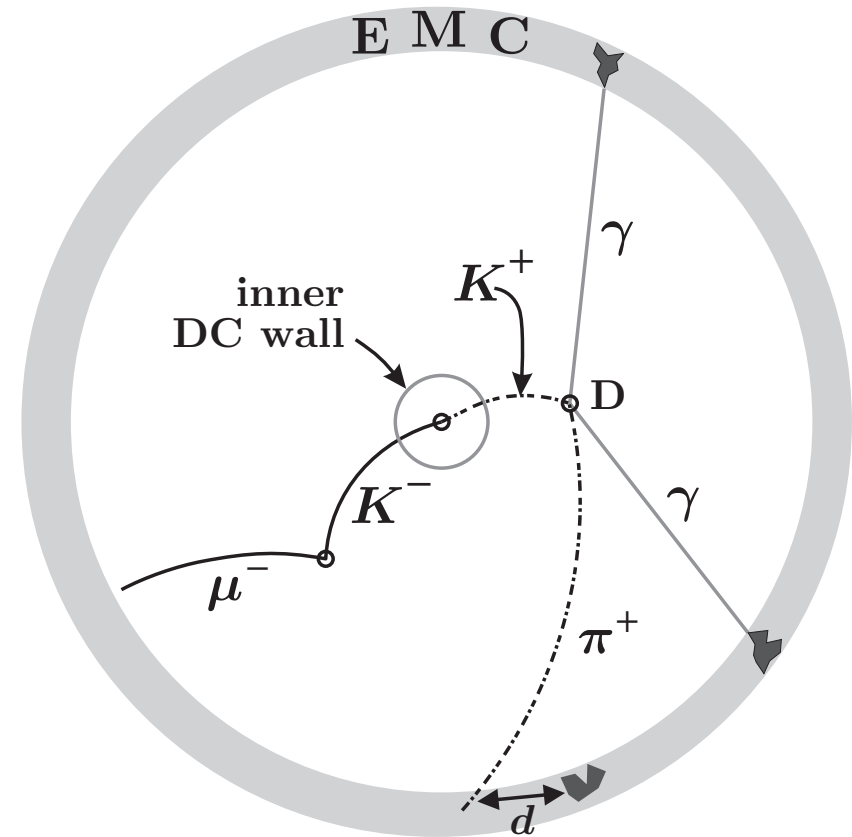


# K<sup>±</sup> → π<sup>+</sup>π<sup>0</sup> decays

No π<sup>0</sup> requirement!



Measuring the  $K \rightarrow \pi$  detection efficiency



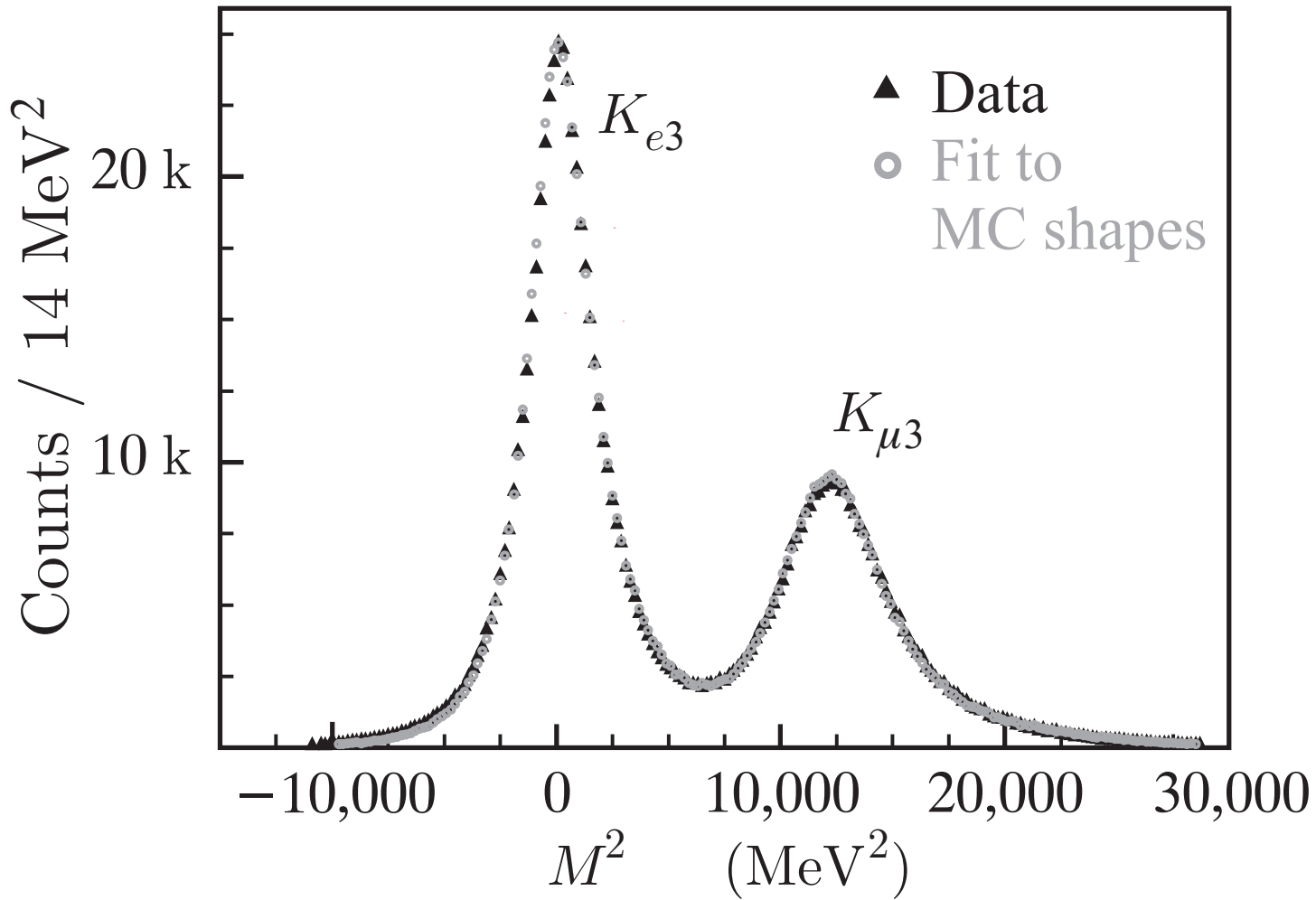
$$\text{BR}(K^+ \rightarrow \pi^+ \pi^0(\gamma)) = 0.2065 \pm 0.0005_{\text{stat}} \pm 0.0008_{\text{sys}}$$

Average of  $K_{\mu 2}^-$  and  $K_{\pi 2}^-$  tags. Normalization for many  $K^\pm$  BRs





# $K_{\ell 3}^{\pm}$ decays



$K \rightarrow \mu\nu$  and  
 $K \rightarrow \pi^{\pm}\pi^0$  tag.

Verify  $\pi^0 \rightarrow 2\gamma$ .

Lepton mass  
from  $p$  and  
TOF.



$$K^\pm \rightarrow \ell^\pm \pi^0 \nu \text{ BRs}$$

BR \ Tag	$K_{\mu 2}^+$	$K_{\pi 2}^+$	$K_{\mu 2}^-$	$K_{\pi 2}^-$
BR( $K_{e3}$ )	0.0495(7)	0.0493(10)	0.0497(8)	0.0502(10)
BR( $K_{\mu 3}$ )	0.0322(6)	0.0322(9)	0.0323(5)	0.0327(9)

8 measurements, very good consistency:

$$\text{BR}(K^\pm \rightarrow e^\pm \pi^0 \nu) = 0.0497 \pm 0.0005$$

$$\text{BR}(K^\pm \rightarrow \mu^\pm \pi^0 \nu) = 0.0323 \pm 0.0004$$



# All KLOE $K^\pm$ results

Constrained fit:

$$\text{BR}(K^\pm \rightarrow \mu^\pm \nu) = 0.63765 \pm 0.0013$$

$$\text{BR}(K^\pm \rightarrow \pi^\pm \pi^0) = 0.20680 \pm 0.0009$$

$$\text{BR}(K^\pm \rightarrow \pi^\pm \pi^+ \pi^-) = 0.05534 \pm 0.0009 \text{ (input PDG '04)}$$

$$\text{BR}(K^\pm \rightarrow e^\pm \pi^0 \nu) = 0.04978 \pm 0.0005$$

$$\text{BR}(K^\pm \rightarrow \mu^\pm \pi^0 \nu) = 0.03242 \pm 0.0004$$

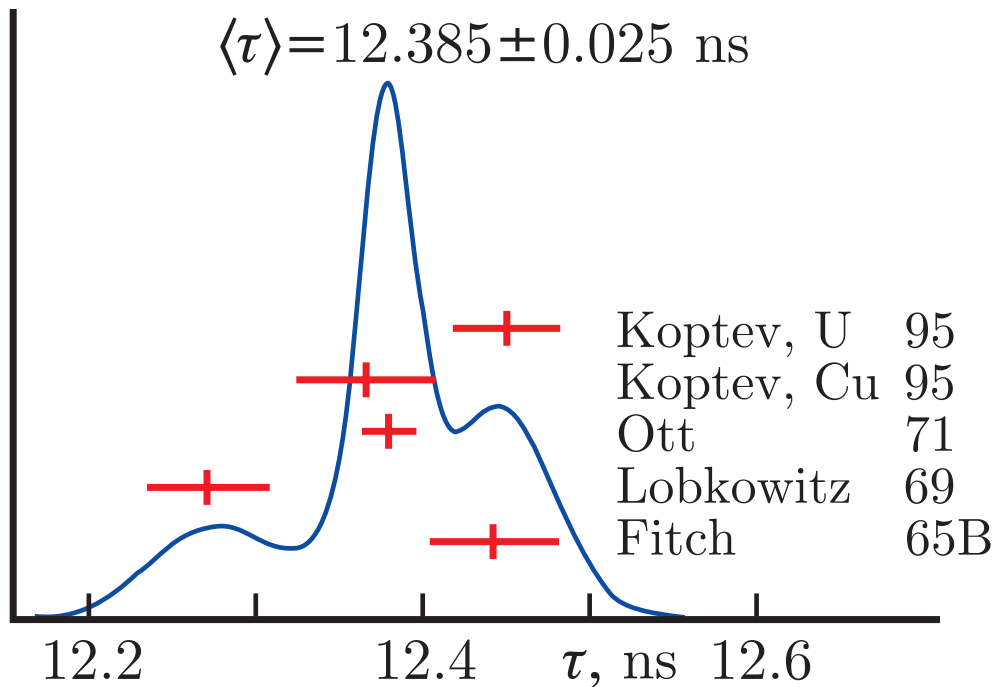
$$\text{BR}(K^\pm \rightarrow \pi^\pm \pi^0 \pi^0) = 0.01765 \pm 0.0002$$

$$\text{sum} = 1.0$$

$$\chi^2/\text{dof} = 0.78/1, \text{ CL}=38\%$$



# Must check $\tau(K^\pm)$



Ott et al.: error is 0.032,  
not 0.016

Koptev:  $\tau = 12.422 \pm 0.048$

Lobkowitz is way off

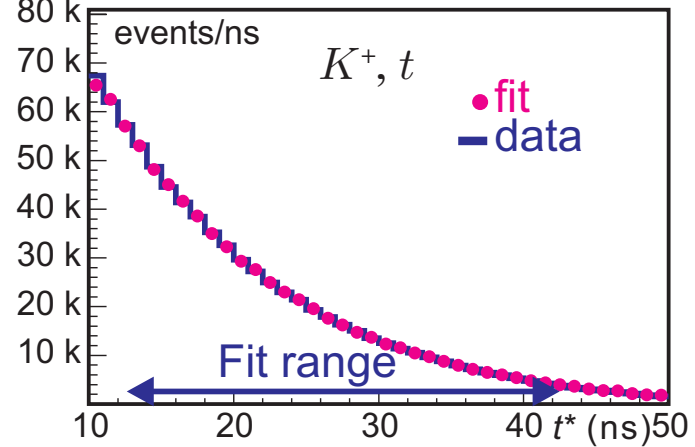
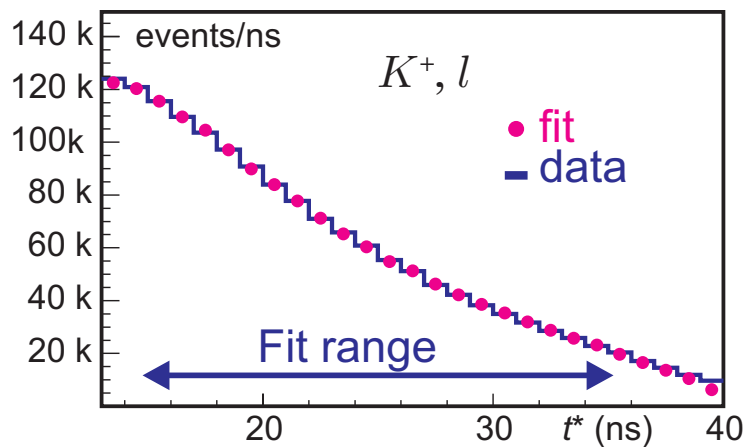
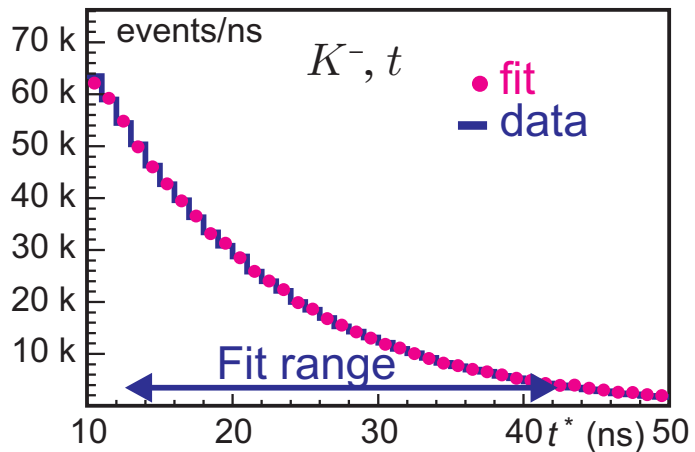
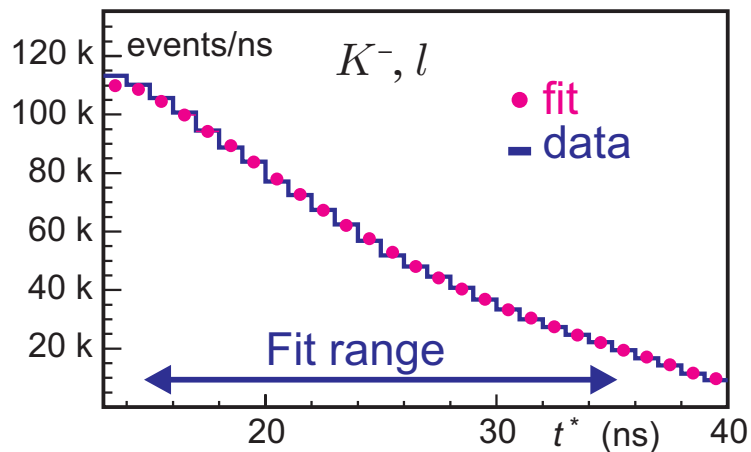
$\delta\tau = 0.0119$  ns  $\Rightarrow$  0.025 ns

Better remeasure!

$\tau = 12.409 \pm 0.022$  ns



# KLOE $\tau(K^\pm)$



4 measurement  
of  $\tau -$

Consistent -

Can average

Use only our result

$$\tau(K^\pm) = 12.347 \pm 0.030, \quad 0.25\%$$

$$\tau(K^+)/\tau(K^\pm) = 1.004 \pm 0.004$$

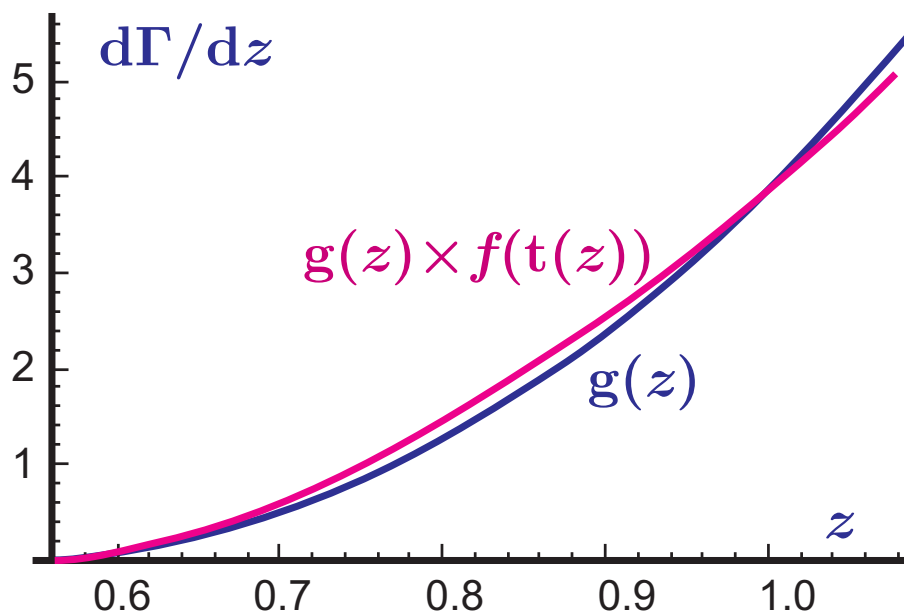
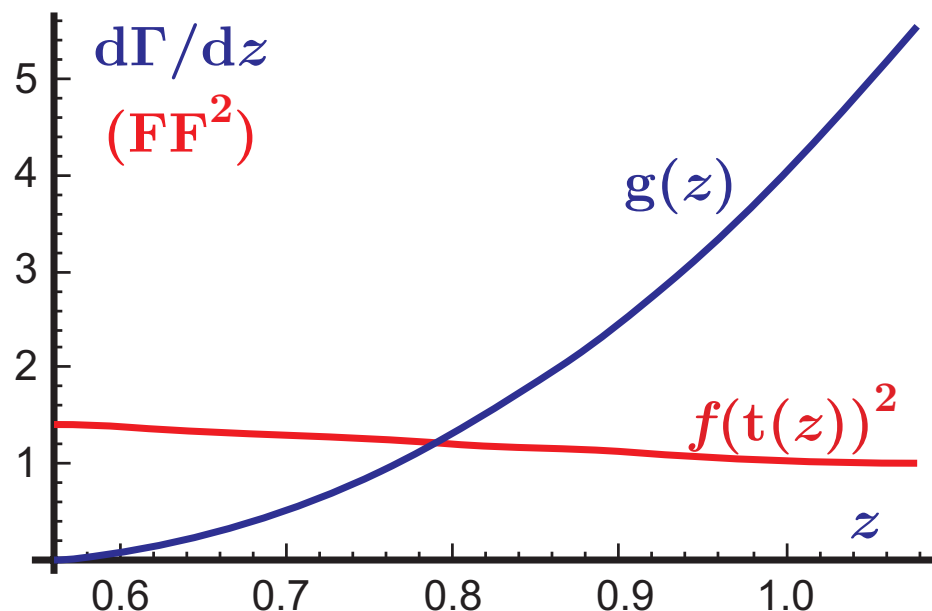


# The form factors

$\tilde{f}_{+,0}(t)?? \Rightarrow$  Phase Space Integrals  $I_{K\ell}$

$$t = (P - p)^2 = M^2 + m^2 - 2M E_\pi, \quad 0 < t < M^2 + m^2 \text{ for } K_{e3}$$

No need to fit Dalitz Plot. All info in  $E_\pi$  spectrum.



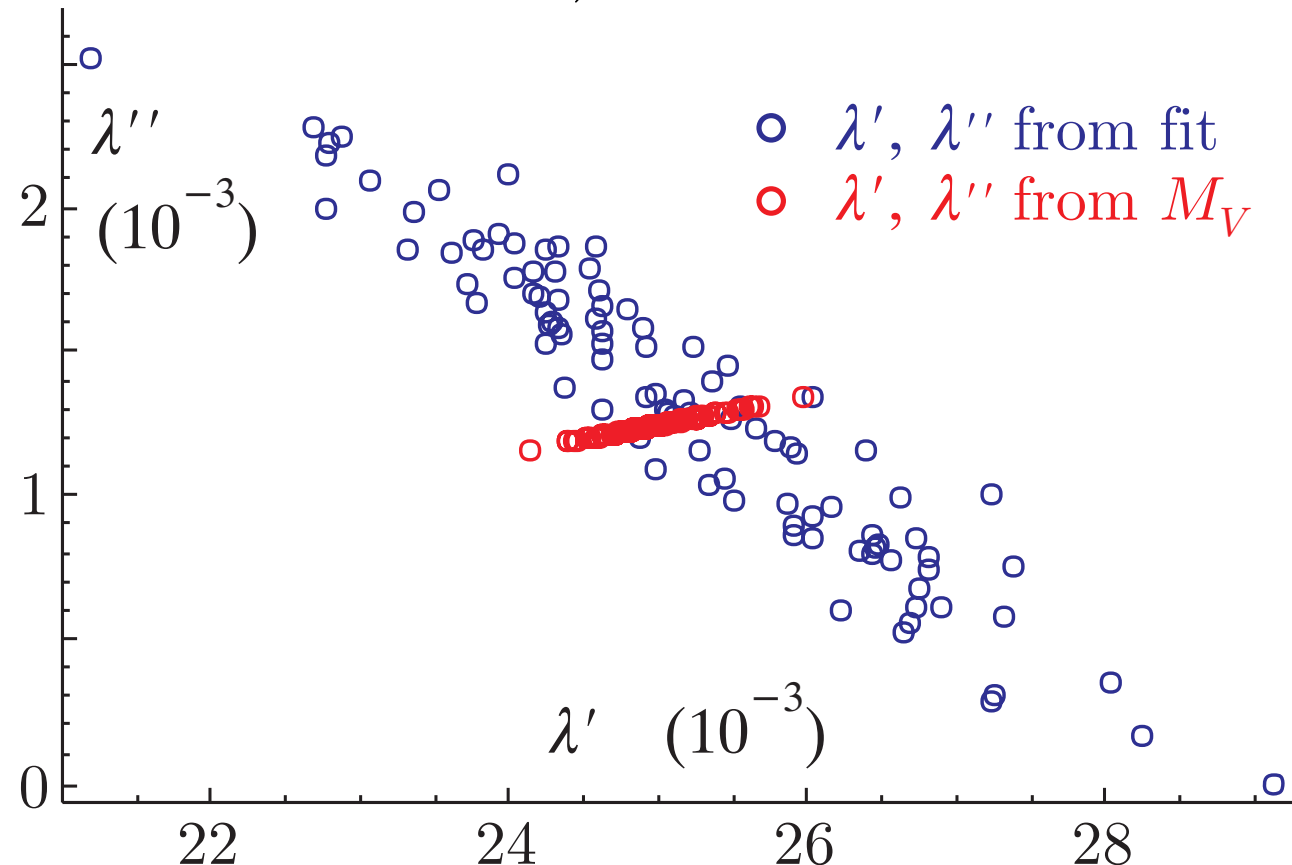
$$z = \frac{2 E_\pi}{M}$$



# FF: quadratic, pole, disp.

$$\tilde{f}_+(t) = 1 + \lambda' \frac{t}{m} + \frac{\lambda''}{2} \left( \frac{t}{m} \right)^2 \quad \text{or} \quad \tilde{f}_+(t) = \frac{M_V^2}{t - M_V^2}$$

$\lambda'$ ,  $\lambda''$  are 95% correlated, errors  $\times 3-4$



$$\lambda' = (m/M_V)^2, \lambda'' = 2\lambda'^2$$

Pole and quad give different  $I$ , KTeV, '04

KTeV unlucky, but should have understood.

Pole was right!!



## Common practice

Quadratic  $FF_+$  and linear  $FF_0$ : three parameters  $\lambda'_+$ ,  $\lambda''_+$ ,  $\lambda_0$

$$\begin{pmatrix} 1.79^2 & 3.51 & -1.98 \\ 3.51 & 3.15^2 & -4.04 \\ -1.98 & -4.04 & 1.37^2 \end{pmatrix} \times \frac{1}{N}$$

This systematically over-estimates  $\lambda_0$  by  $\sim 20\%$ . Errors “look” acceptable.

$$\begin{pmatrix} \lambda'_0 & \lambda''_0 & \lambda'_+ & \lambda''_+ \\ 63.9^2 & -1200 & -923 & 197 \\ -1200 & 18.8^2 & 272 & -59 \\ -923 & 272 & 14.8^2 & -49 \\ 197 & -59 & -48 & 3.4^2 \end{pmatrix} \times \frac{1}{N}$$

Errors are unacceptable, correlations reach  $-99.96\%!!$

NEED HELP!





## Dispersive approach

Use a dispersion relation for  $\log(\tilde{f}_+(t))$ , twice subtracted at  $t=0$ .  
 $\tilde{f}_+(t) = \exp[(t/m^2)(\lambda_+ + H(t))]$ .  $\lambda_+$  slope at  $t=0$  to be measured.  $H(t)$  is obtained from  $K\pi$   $p$ -wave phase, dominated by  $K^*(892)$ .  $\tilde{f}_+$  expands to

$$\tilde{f}_+ = 1 + \lambda_+ \frac{t}{m^2} + \frac{\lambda_+^2 + 0.00058}{2} \left(\frac{t}{m^2}\right)^2 + \frac{\lambda_+^3 + 3 \times 0.00058 \times \lambda_+ + 0.00003}{6} \left(\frac{t}{m^2}\right)^3$$

The Callan-Treiman relation relates the scalar FF at  $t = \Delta_{K\pi}$  to the ratio of the pseudoscalar decay constants  $f_K/f_\pi$ .

$$\tilde{f}_0(\Delta_{K\pi}) = \frac{f_K}{f_\pi} \frac{1}{f_+(0)} + \Delta_{CT}$$

where  $\Delta_{K\pi} = M^2 - m^2$ , the so called Callan-Treiman point, and  $\Delta_{CT}$  is a correction which at NLO in  $\chi$ PT is  $\sim -3.5 \times 10^{-3}$ .



Use dispersion relation subtracted at  $t=0$  and  $t = \Delta_{K\pi}$ . Use  $s$ -wave  $K\pi$  phases to compute the equivalent function  $G(t)$ .  $\tilde{f}_0(t)$  expands to

$$\tilde{f}_+ = 1 + \lambda_+ \frac{t}{m^2} + \frac{\lambda_+^2 + 0.00042}{2} \left(\frac{t}{m^2}\right)^2 + \frac{\lambda_+^3 + 3 \times 0.00042 \times \lambda_+ + 0.000027}{6} \left(\frac{t}{m^2}\right)^3$$

(Use dispersion relations and  $\chi$ PT: Stern *et al.*. Also Pich *et al.*)

FFs are power series in  $t$  but with only one free parameter each!

Fit for  $\lambda_{+,0}$  and get “phase space integrals”

$I(K_{e3}^0)$	$I(K_{\mu 3}^0)$	$I(K_{e3}^+)$	$I(K_{\mu 3}^+)$
0.6191(14)	0.4105(18)	0.6365(14)	0.4224(18)

0.2 to 0.4 %.  $\sim 1/3$  of KTeV error.



[Γ]

$$f(0)|V_{us}|$$

We are almost there. Let us begin with  $f(0)|V_{us}|$ .  $f(0)$ , by definition, belongs to  $K^0 \rightarrow \pi^\pm$ . If we know the I-spin corrections and the long distance EM corrections all decays must give the same value for  $f(0)|V_{us}|$ .

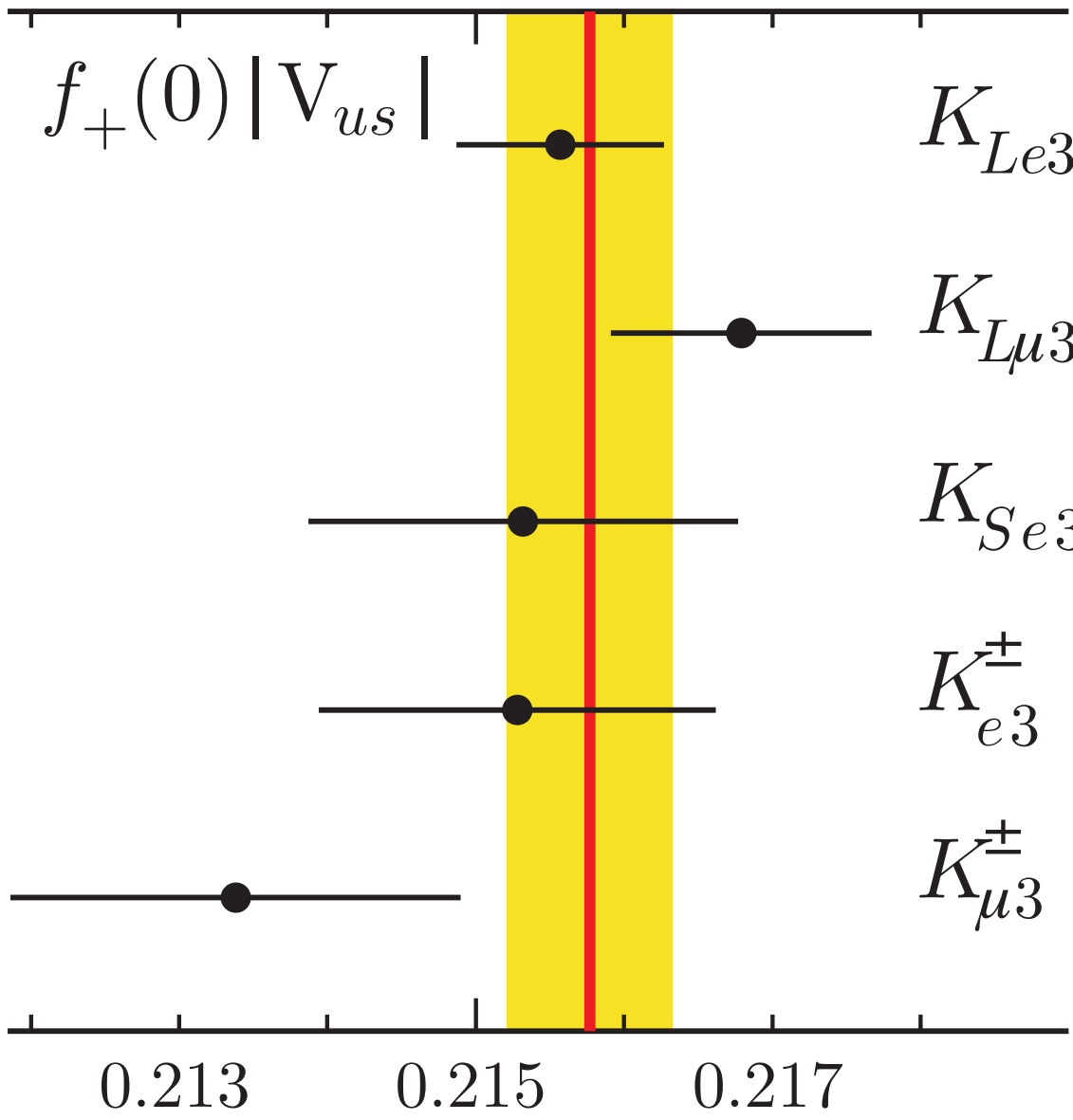
This is in fact a check of the  $\delta \dots$  estimates and of the experiment.

We find

Channel	$f(0) V_{us} $	Correlation coefficients			
$K_{Le3}$	0.2155(7)				
$K_{L\mu3}$	0.2167(9)	0.28			
$K_{Se3}$	0.2153(14)	0.16	0.08		
$K_{e3}^\pm$	0.2152(13)	0.07	0.01	0.04	
$K_{\mu3}^\pm$	0.2132(15)	0.01	0.18	0.01	0.67



$f(0)|V_{us}|$



$\langle \rangle = -0.2157 \pm 0.0006$

$\chi^2/ndf = 7.0/4$   
(CL=13%)

$SU(2)$  corrections are verified to  $1.1 \sigma$ .



# Lepton Universality

From:

$$R_{\mu e} \equiv \frac{[f(0)V_{us}]_{\mu 3}^2}{[f(0)V_{us}]_{e 3}^2} = \frac{\Gamma_{\mu 3} I_{e 3} (1 + \delta_e)^2}{\Gamma_{e 3} I_{\mu 3} (1 + \delta_\mu)^2} \quad [\Gamma]$$

$\Rightarrow$

$$R_{\mu e} = \frac{g_e^2}{g_\mu^2} = 1.000 \pm 0.008 \quad [W_\alpha J^\alpha]$$

Competitive with pure leptonic processes,  $\tau$  decays

$$(R_{\mu e})_\tau = 1.000 \pm 0.004$$

Also with  $\pi \rightarrow \ell \nu$  results

$$(R_{\mu e})_\pi = 1.0042 \pm 0.0033$$



## $f(0)$ and $|V_{us}|$

The accepted value was  $|f(0)| = 0.961 \pm 0.008$ , Roos & Leutwyler, 1984. Recently Lattice results have become convincing. We take  $f(0) = 0.9644 \pm 0.0049$ , RBC and UKQCD, 2007. Then

$$|V_{us}| = 0.2237 \pm 0.0013$$

$$|V_{us}|^2 = 0.05002 \pm 0.00057$$

## $0^+ \rightarrow 0^+$ and $|V_{ud}|^2$

$$|V_{ud}|^2 = 0.9490 \pm 0.0005$$

$$1 - |V_{us}|^2 - |V_{ud}|^2 = 0.0009 \pm 0.0008 \quad (\sim 1.1\sigma)$$

**But we can do better**



$$K \rightarrow \mu\nu / \pi \rightarrow \mu\nu$$

$$\frac{\Gamma(K_{\mu 2}(\gamma))}{\Gamma(\pi_{\mu 2}(\gamma))} = \frac{|V_{us}|^2}{|V_{ud}|^2} \frac{f_K^2}{f_\pi^2} \frac{m_K (1 - m_\mu^2/m_K^2)^2}{m_\pi (1 - m_\mu^2/m_\pi^2)^2} \times (0.9930 \pm 0.0035)$$

Cancellations in  $f_K^2/f_\pi^2$ . From lattice

Cancellations in rad. cor. From Marciano.

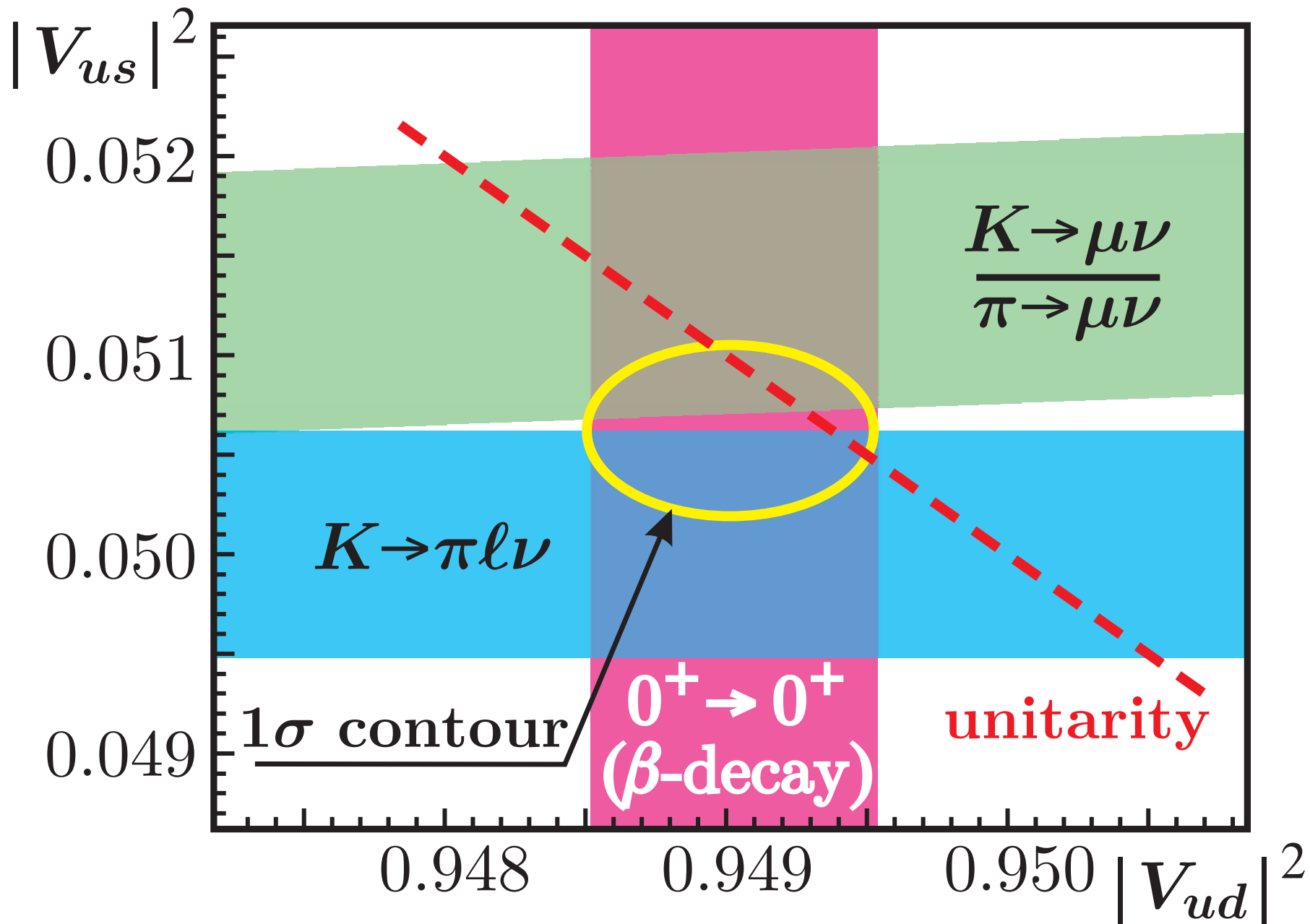
$$\text{KLOE} \Rightarrow |V_{us}/V_{ud} \times f_K/f_\pi|^2 = 0.0765 \pm 0.0005$$

$$\text{HPQCD/UKQCD} \Rightarrow f_K/f_\pi = 1.189 \pm 0.007$$

$$|V_{us}/V_{ud}|^2 = 0.0541 \pm 0.0007$$



# CKM Unitarity





## From the fit

$$1 - |V_{us}|^2 - |V_{ud}|^2 = 0.0004 \pm 0.0007 \quad (\sim 0.6\sigma)$$

$$|V_{us}| = 0.2249 \pm 0.0010$$

$$|V_{ud}| = 0.97417 \pm 0.00026$$

## Fit with unitarity as constraint

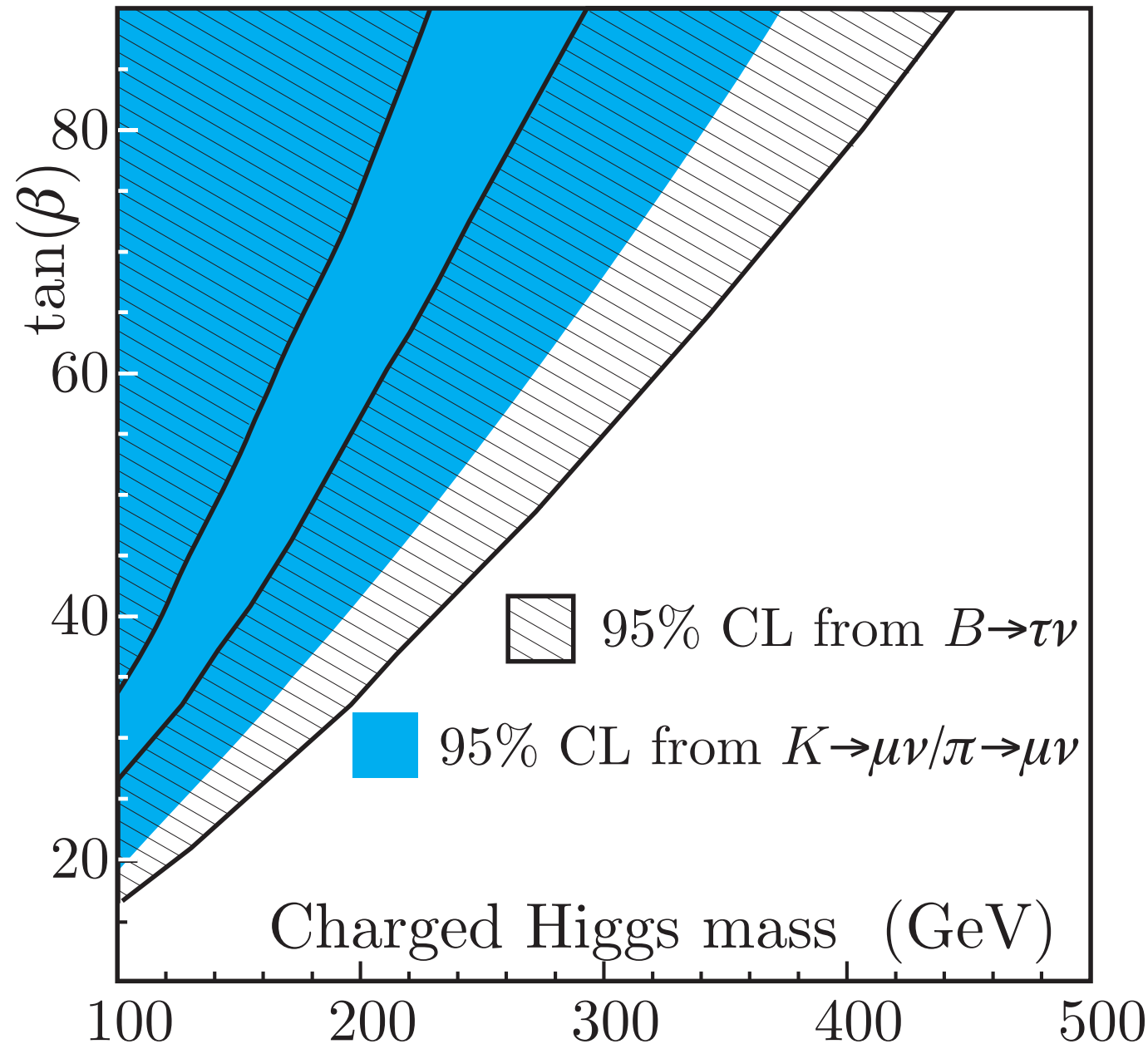
$$|V_{us}| = 0.2253 \pm 0.0007$$

$$|V_{ud}| = \sqrt{1 - |V_{us}|^2} = 0.97429 \pm 0.00017.$$

$$\theta_C = 13.02^\circ \pm 0.06^\circ$$



# tan $\beta$ – $M(\text{Higgs}^\pm)$



$$R_{\ell 23} = \left| \frac{V_{us}(K_{\mu 2})}{V_{us}(K_{\ell 3})} \times \frac{V_{ud}(0^+ \rightarrow 0^+)}{V_{ud}(\pi_{\mu 2})} \right| = 1$$

in the SM. With charged Higgs exchange  $R_{\ell 23} =$

$$\left| 1 - \frac{m_{K^+}^2}{m_{H^+}^2} \left( 1 - \frac{m_{\pi^+}^2}{m_{K^+}^2} \right) \times \frac{\tan^2 \beta}{1 + \epsilon_0 \tan \beta} \right|$$

$\epsilon_0 \approx 1/(16\pi^2)$



# Conclusions

Unitarity is verified at the 0.1% level

$$1 - |V_{us}|^2 - |V_{ud}|^2 = 0.0004 \pm 0.0007$$

Muons, electrons and quarks

carry the same weak charge to better than 0.5%

Kaon properties limit the parameter space  
for some new physics models

