Introduction to Particle Physics - Chapter 13 -Higgs boson discovery



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Chapter summary:

- Higgs coupling with fermions and gauge bosons
- Branching ratios
- Higgs search at LEP
- Higgs search at LHC
- Higgs discovery at LHC
- Higgs properties
- A few words about Supersymmetry
- A few words about Dark Matter

Higgs couplings with W and Z

• Electroweak Lagrangian that is invariant for a local gauge transformation:

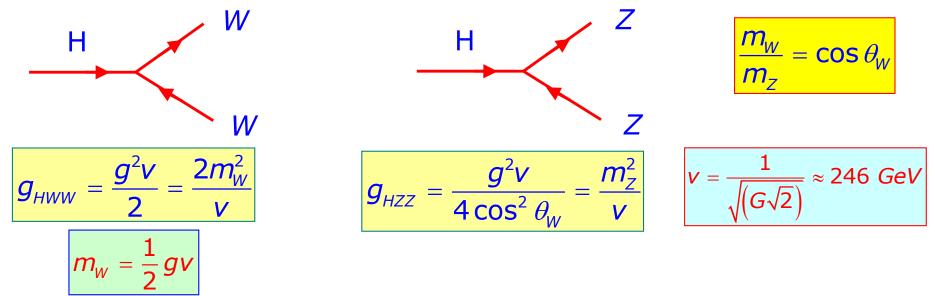
$$L = \overline{\Psi}_{L} \gamma^{\mu} \left[i \partial_{\mu} - g \vec{I} \cdot \vec{W}_{\mu}(x) - \frac{g'}{2} Y \cdot B_{\mu} \right] \Psi_{L} + \overline{\Psi}_{R} \gamma^{\mu} \left[i \partial_{\mu} - \frac{g'}{2} Y \cdot B_{\mu} \right] \Psi_{R} + L_{free}(\vec{W}, B)$$

- Replacing in the Lagrangian the field ϕ obtained after the spontaneous symmetry breaking

$$\varphi(x) = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + H(x) \end{pmatrix}$$

we get the Higgs couplings with the gauge bosons:

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Higgs couplings with the fermions

• After the spontaneous symmetry breaking we insert in the Lagrangian the field

$$\varphi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + H \end{pmatrix}$$

and we get:

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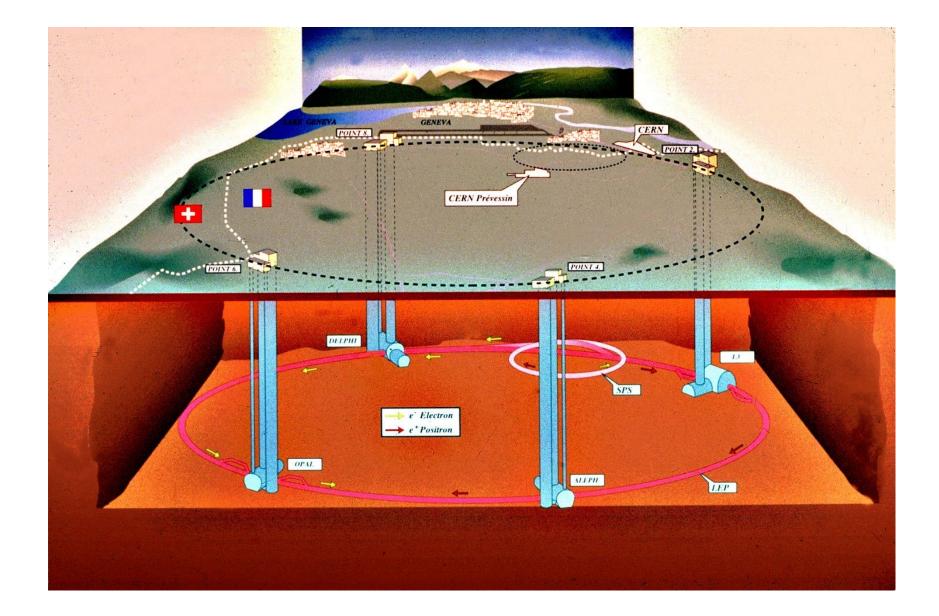
$$L = -\frac{g_e v}{\sqrt{2}} \left[\bar{e}_R e_L + \bar{e}_L e_R \right] - \frac{g_e}{\sqrt{2}} \left[\bar{e}_R e_L + \bar{e}_L e_R \right] H$$

$$\uparrow$$
Mass term
Higgs coupling with the electron

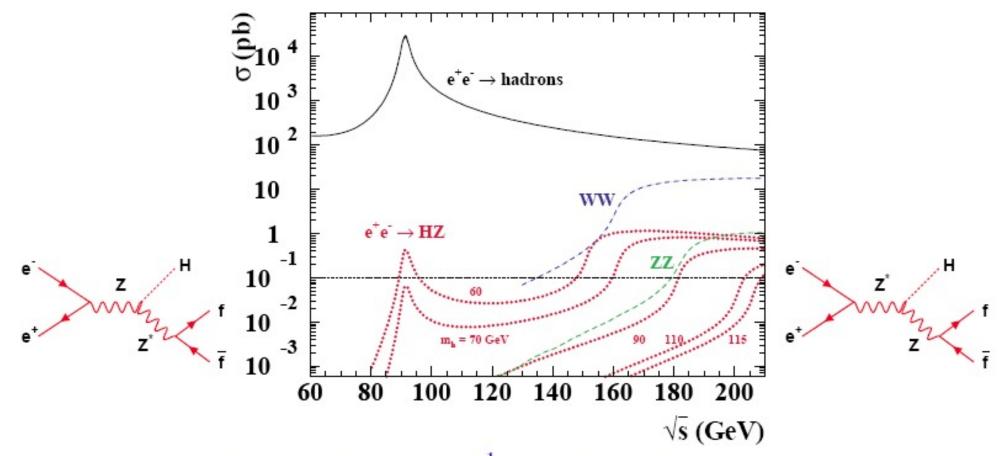
$$\mathsf{m}_e = rac{\boldsymbol{g}_e \cdot \boldsymbol{v}}{\sqrt{2}}$$

$$\Box = -m_e \bar{e}e - \left(\frac{m_e}{v}\right) \bar{e}eH$$

N.B. the coupling constant is proportional to the fermion mass



Higgs production cross-section at LEP

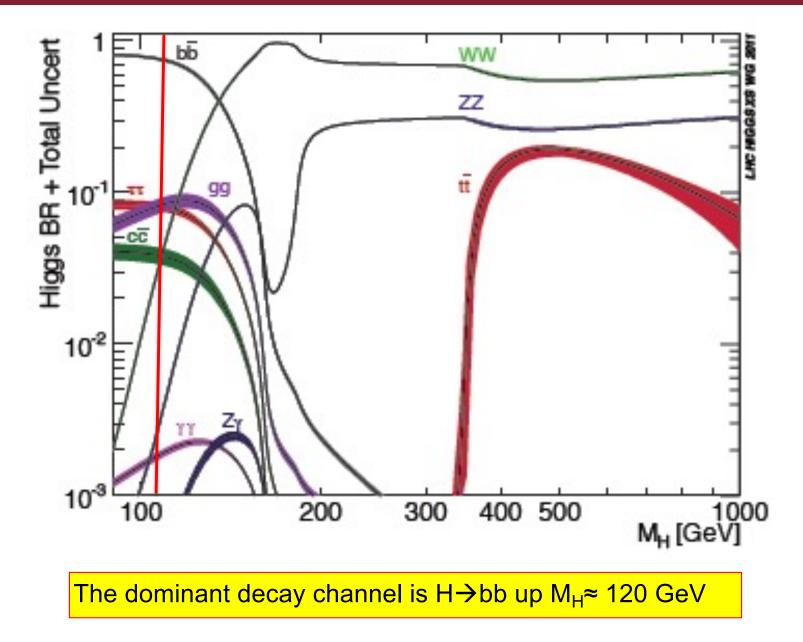


With a luminosity of about 100pb^{-1} and reasonable detection efficiency, sensitive to a cross section of O(0.1) pb.

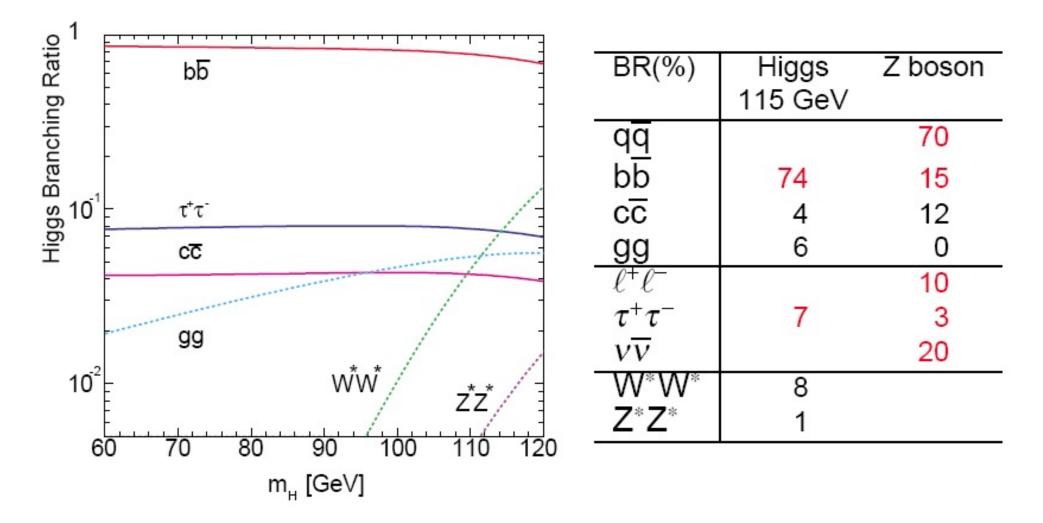
Need LEP2 to produce $m_{\rm H} \ge 65$ GeV. Reach $m_{\rm H} \le \sqrt{s} - M_{\rm Z}$

Must take into account many background processes

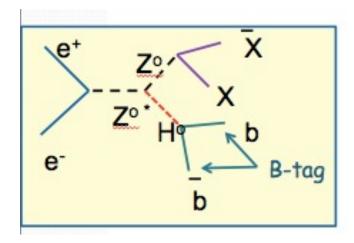
Higgs Branching Ratios



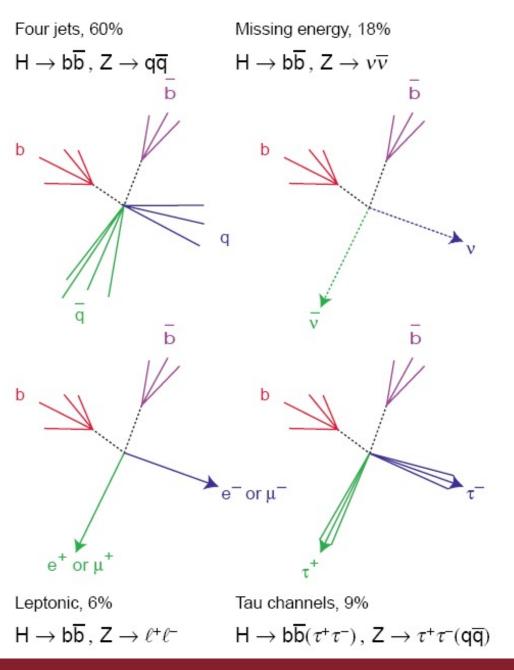
"Higgs couples to mass"



HZ decays

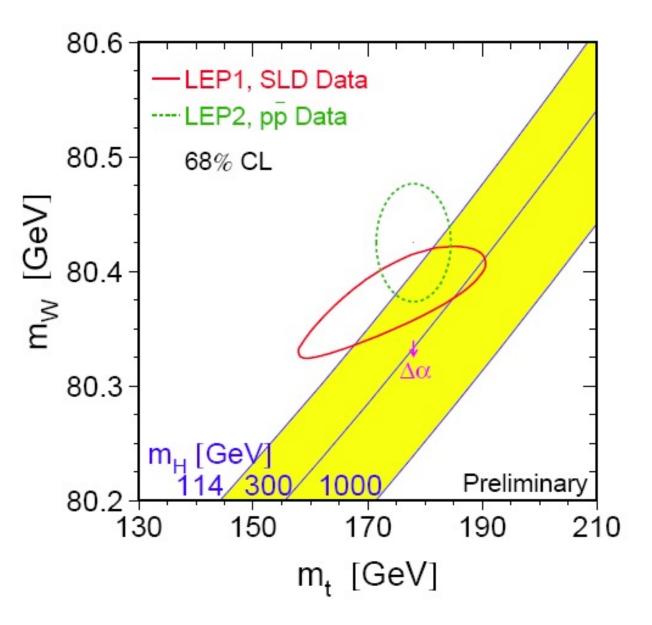


- Higgs boson decays into a b quark pair
- According to the Z decays we have different event topologies:
 - Z → e⁺e⁻; μ⁺μ⁻ (small B.R. but very little background, it was the golden channel)
 - Z → vv (good compromise between B.R. and background)
 - $Z \rightarrow q\overline{q}$ (High QCD background)



Global electroweak fit and Higgs mass

Fit to data from LEP, SLD, Tevatron... Electroweak variables depend on m_t^2 and $\log m_{\rm H}$ through radiative corrections Consistency between predicted top and W mass (Z pole) and direct measurements Preference for low Higgs mass.

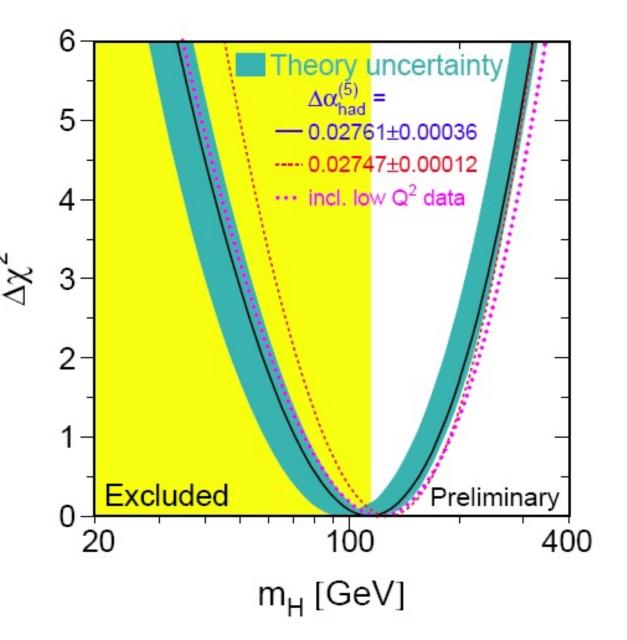


Electroweak fits $m_{\rm H} < 237 \text{ GeV} (95\% \text{ CL})$

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Theory: self consistency of SM to GUT scale $\approx 10^{16}$ GeV $130 < m_{\rm H} < 190$ GeV. $m_{\rm H}$ higher - theory non-perturbative,

*m*_H lower - vacuum unstable.



The Large Hadron Collider LHC

Installed in 26.7 km LEP tunnel

Depth of 70-140 m

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LHC Control Room

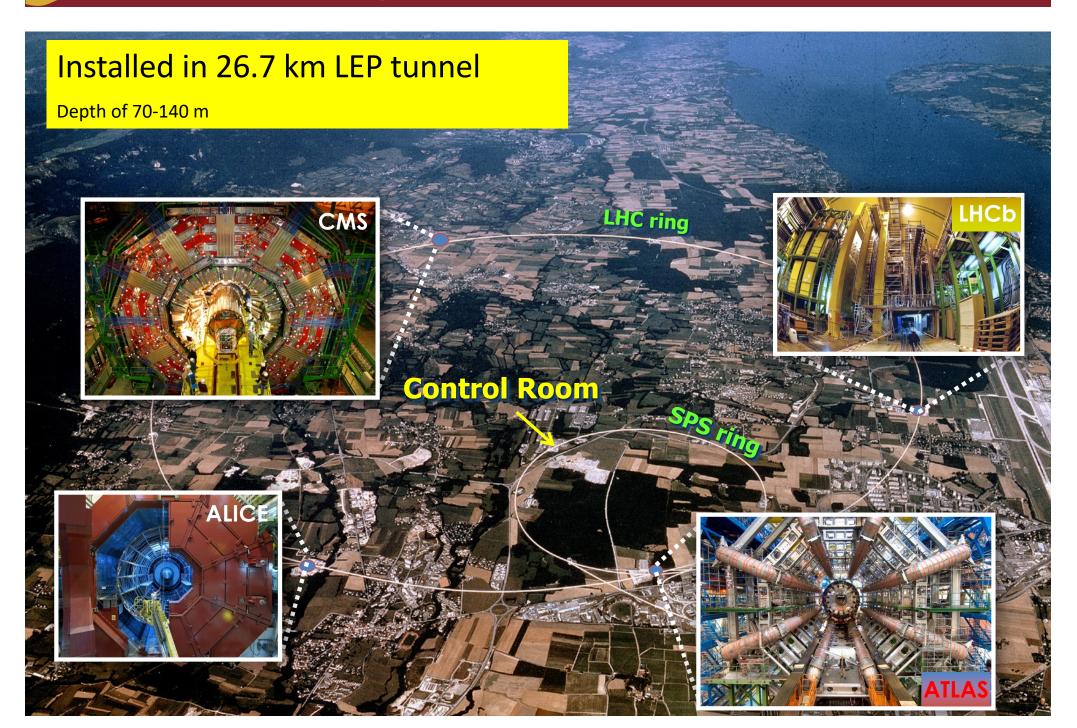


CERN Meyrin

LHC ring

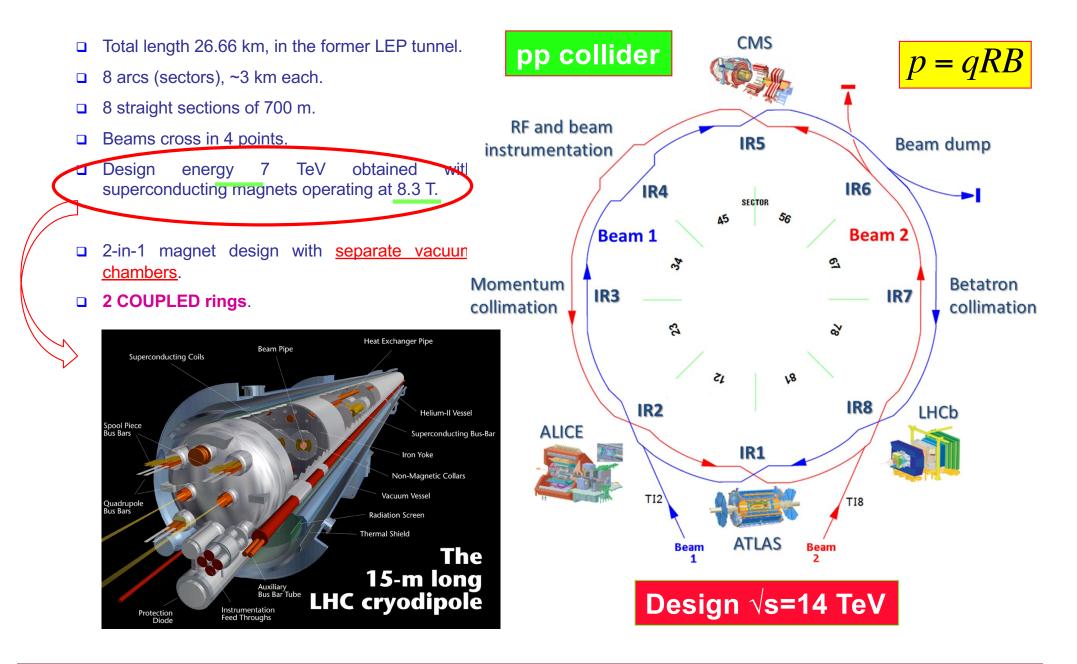


The Large Hadron Collider LHC

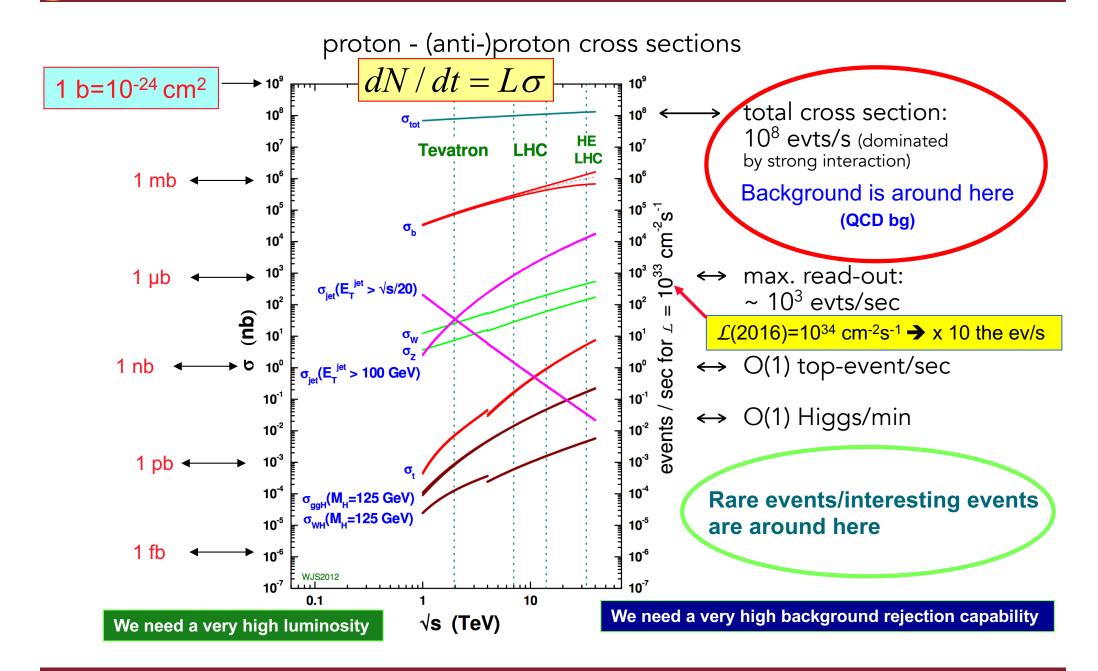




LHC ring layout



Production cross-section at LHC



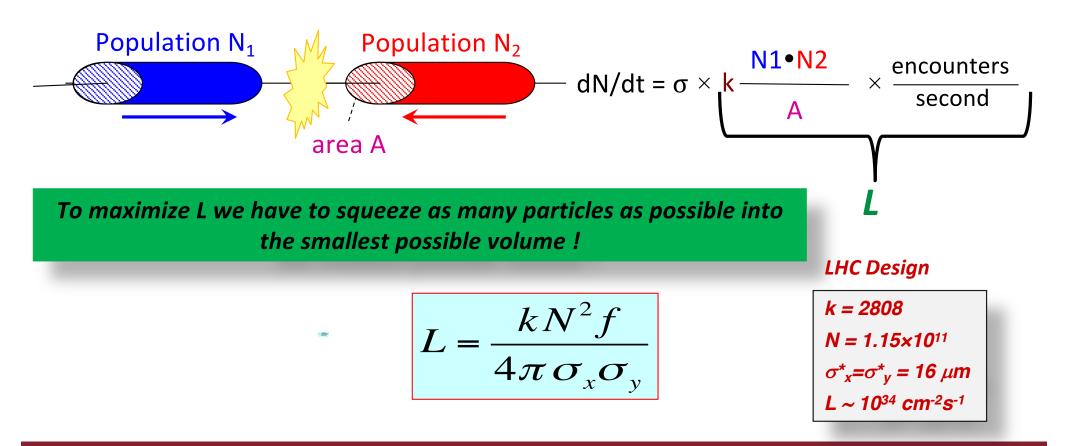


Collider Luminosity

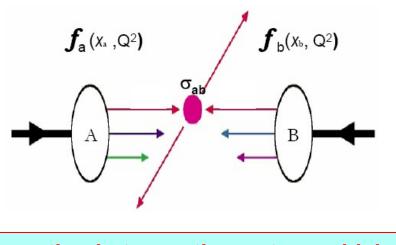
The key parameter for the experiments is the event rate dN/dt. For a physics process with <u>cross-section σ </u> it is proprotional to the collider <u>Luminosity L</u>:

 $dN/dt = L\sigma$

unit of L : (surface \times time)⁻¹

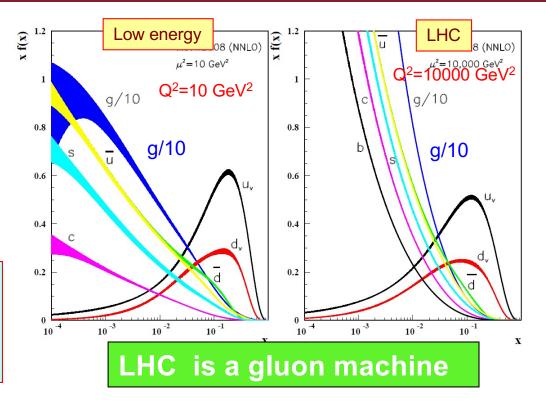


LHC: parton-parton interactions



Interaction between the partons which constitute the hadrons:

not well defined parton energy but energy distribution \rightarrow pdf



PDFs are parameterizations of the partonic content of the proton:

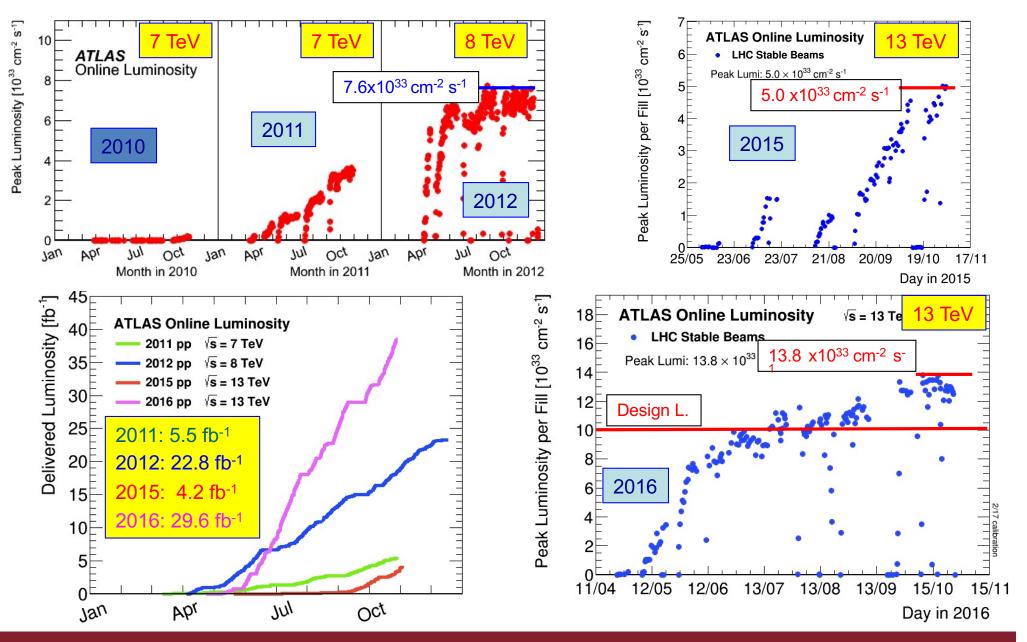
at Hadron Colliders cross-section calculations are a convolution of the cross-section at parton level and PDFs

$$\sigma_{X} = \sum_{a,b} \int_{0}^{1} dx_{a} dx_{b} f(x_{a}, flav_{a}, Q^{2}) f(x_{b}, flav_{b}, Q^{2}) \sigma_{ab \to X}(x_{a}, x_{b}, Q^{2})$$

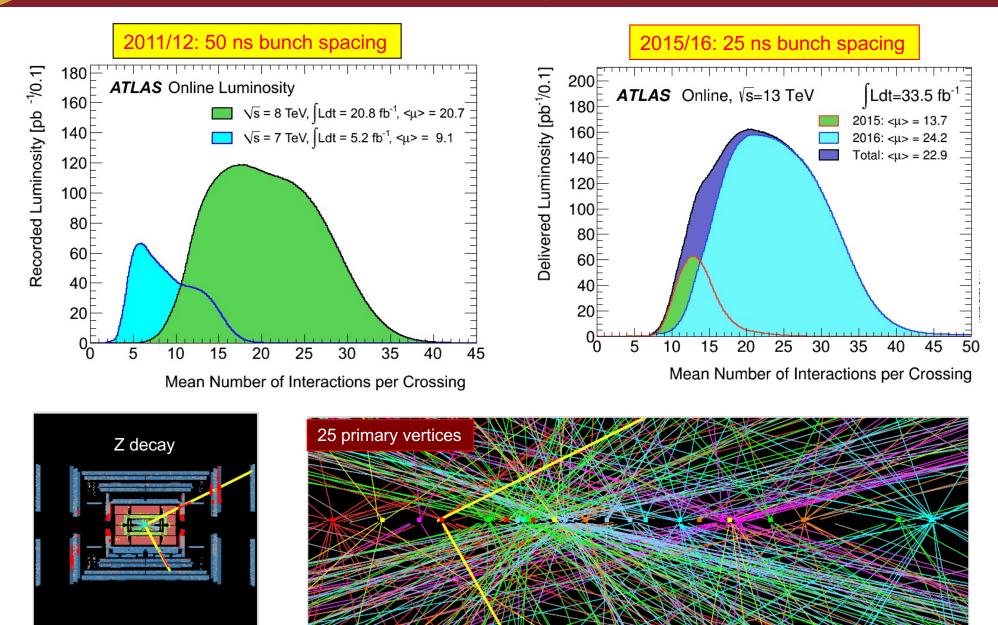
Sum over initial partonic states a,b Parton Density Function hard scattering cross-section

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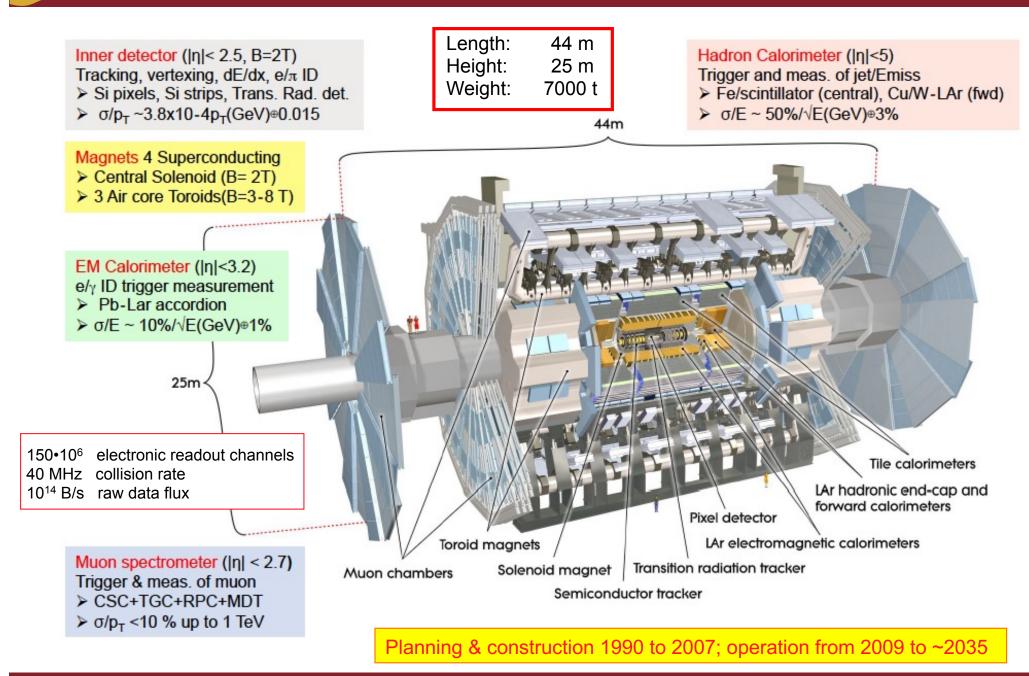
^{1/3} Luminosity: 2011, 2012, 2015 and 2016



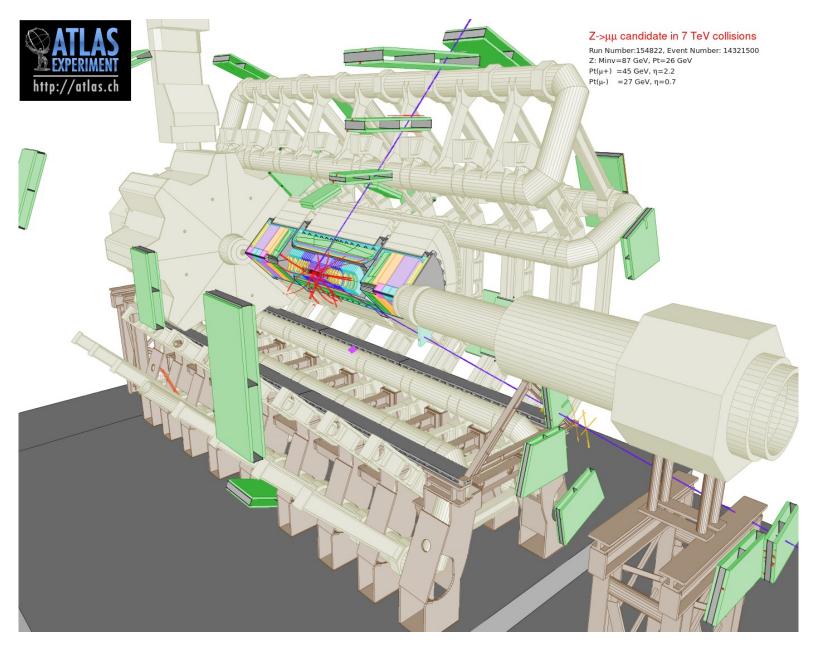
Pile-up



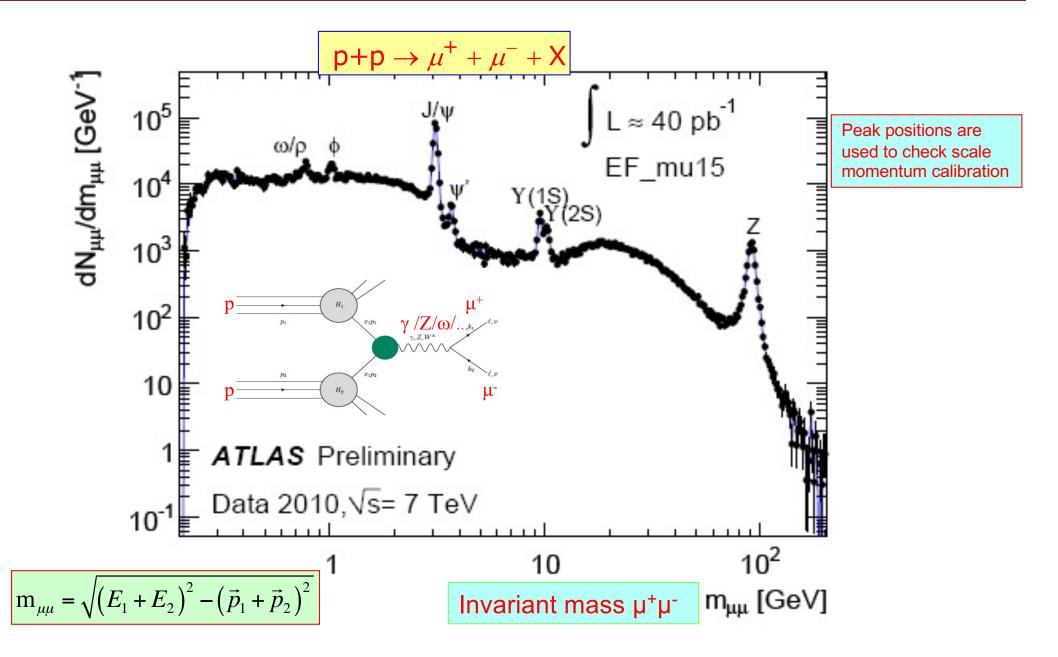
The ATLAS detector



ATLAS: $Z \rightarrow \mu^+ \mu^-$ candidate



ATLAS: 50 years in one slide

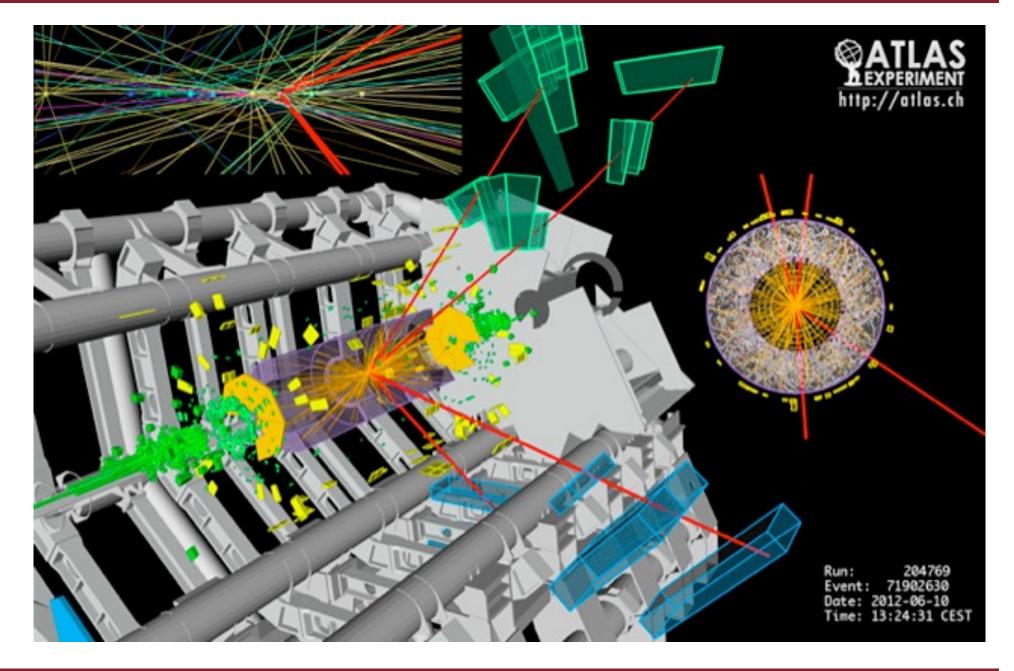


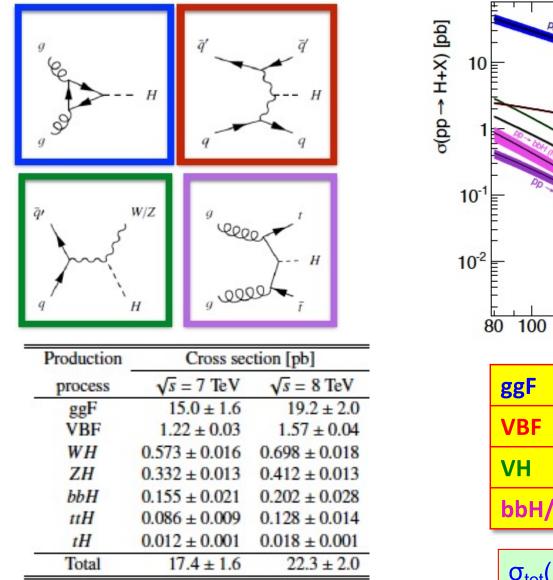
The Higgs boson

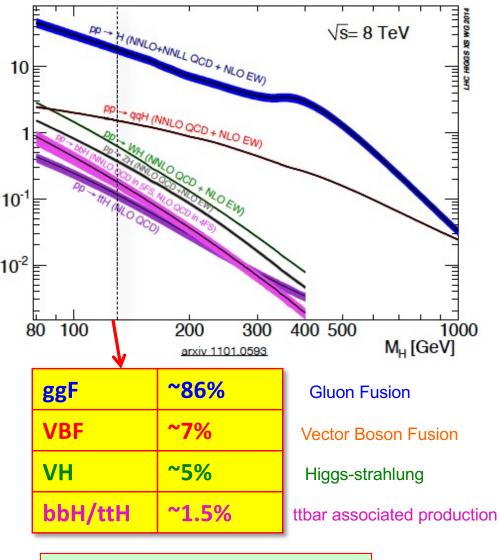


The first Higgs seen in the ATLAS Experiment

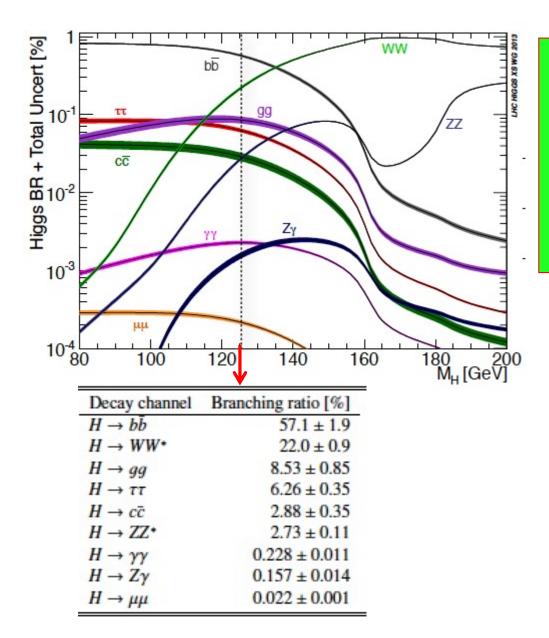








$$\sigma_{tot}(13 \text{ TeV}) = \sim 2\sigma_{tot}(8 \text{ TeV})$$



Higgs couplings (tree level)

$$g_{HWW} = gm_{W}$$

$$g_{HZZ} = g \frac{m_{Z}}{2\cos \theta_{W}}$$

$$g_{Hee} = g \frac{m_{e}}{2m_{W}}$$

Most used decay modes are built using isolated leptons (e,µ), photons and missing energy

H→ZZ*→4I(<u>e,µ</u>)	~0.013%
Н→үү	~0.23%
Н→π	~6.3%
H→WW→IvIv	~1.1%

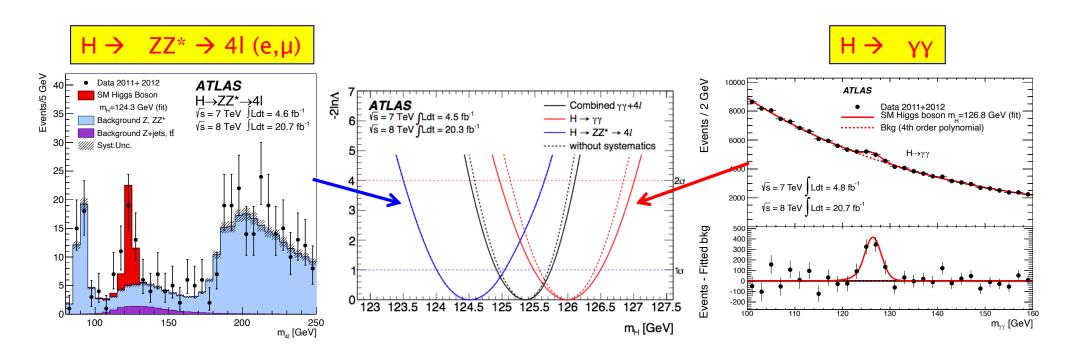
Higgs Boson Mass Measurement

Precise measurement of m_H from channels with the best mass resolution: H→γγ and H→ZZ*→4I (e,µ) (but B.R.≈0.25% only)

Dominant

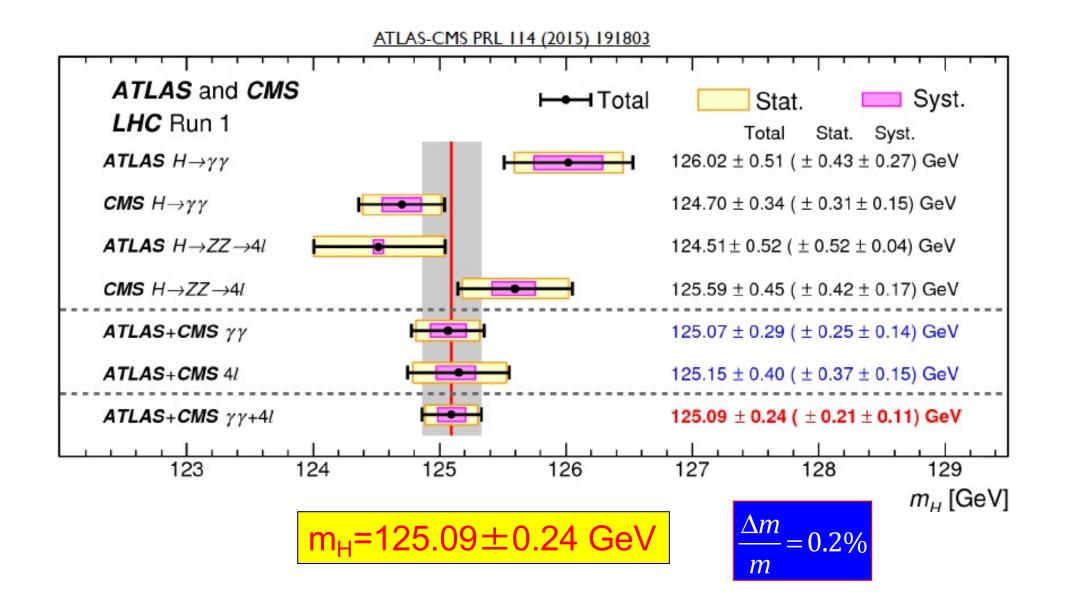
uncertainties:

photon energy scale $(H \rightarrow \gamma \gamma)$, statistics $(H \rightarrow 4I)$



Combined mass: m_H =125.36±0.37 (stat) ±0.18 (syst) GeV

Combined Higgs Boson Mass

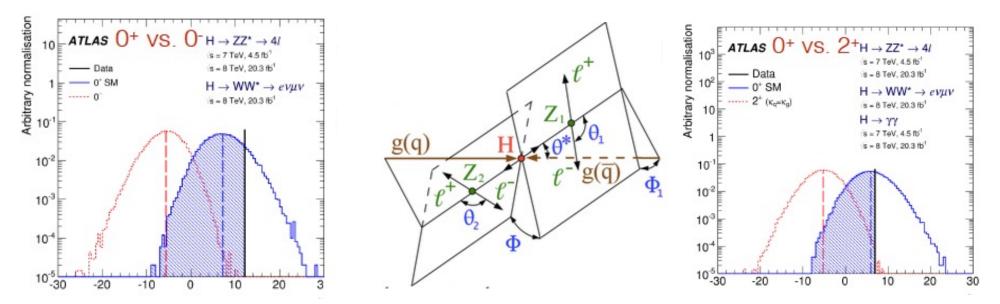


Higgs Boson Spin Measurement

Test SM (0+) against various models

- Spin-2 Higgs
- Spin-0 odd (BSM Higgs)
- (Spin-1 ruled out by observation of $H \rightarrow \gamma \gamma$ decays)

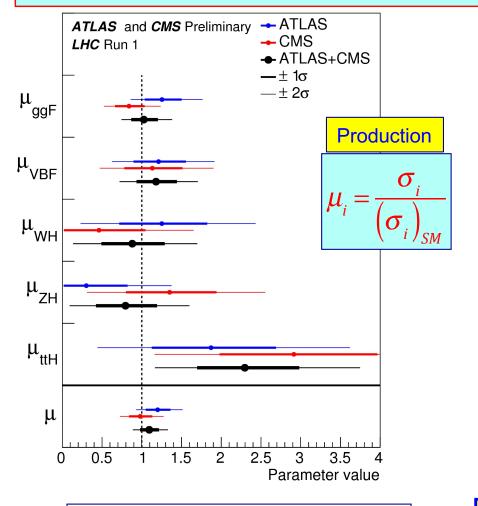
In all tested cases non-SM models rejected at >99% CL



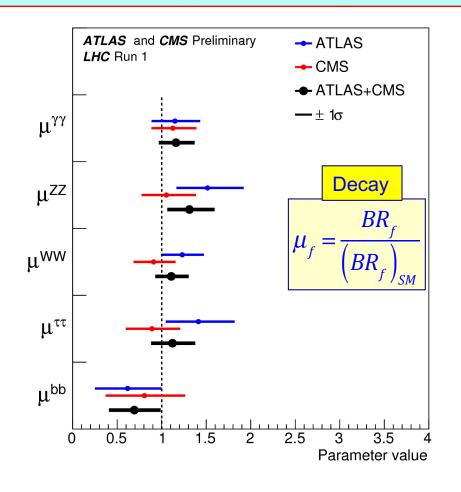
(Multivariate analysis (MVA) based on angular variables)

Comparison with SM expectations

Measure the ratio between observed rate and SM Higgs boson expectation

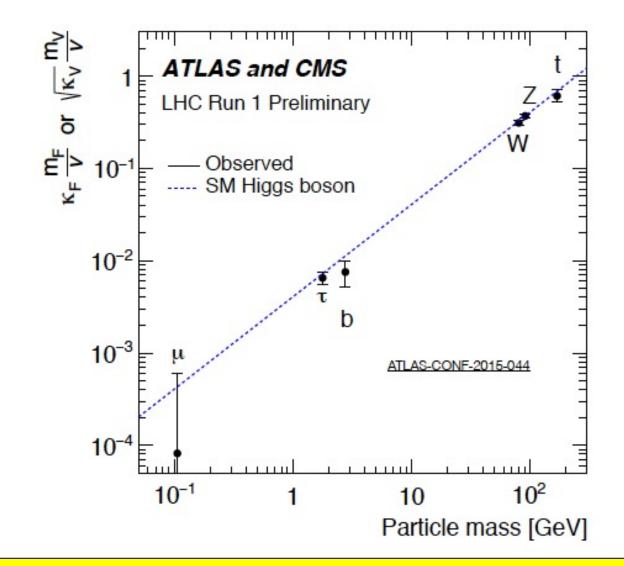


(µ on production modes have been combined assuming SM BR for the decay)



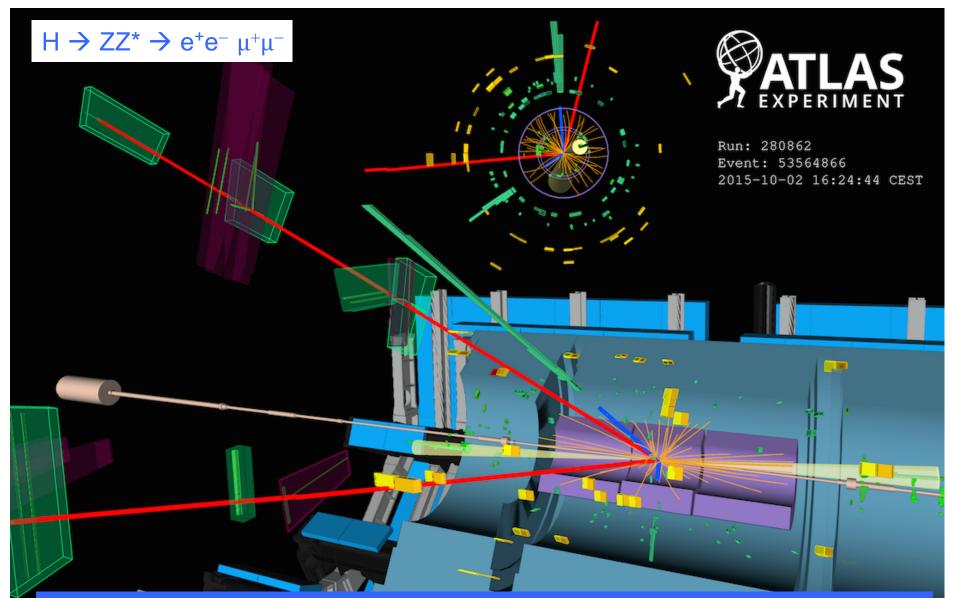
Results are SM like (all $\mu_s \sim 1$)

Higgs Boson Couplings



Coupling strengths scale with mass just as predicted by the SM

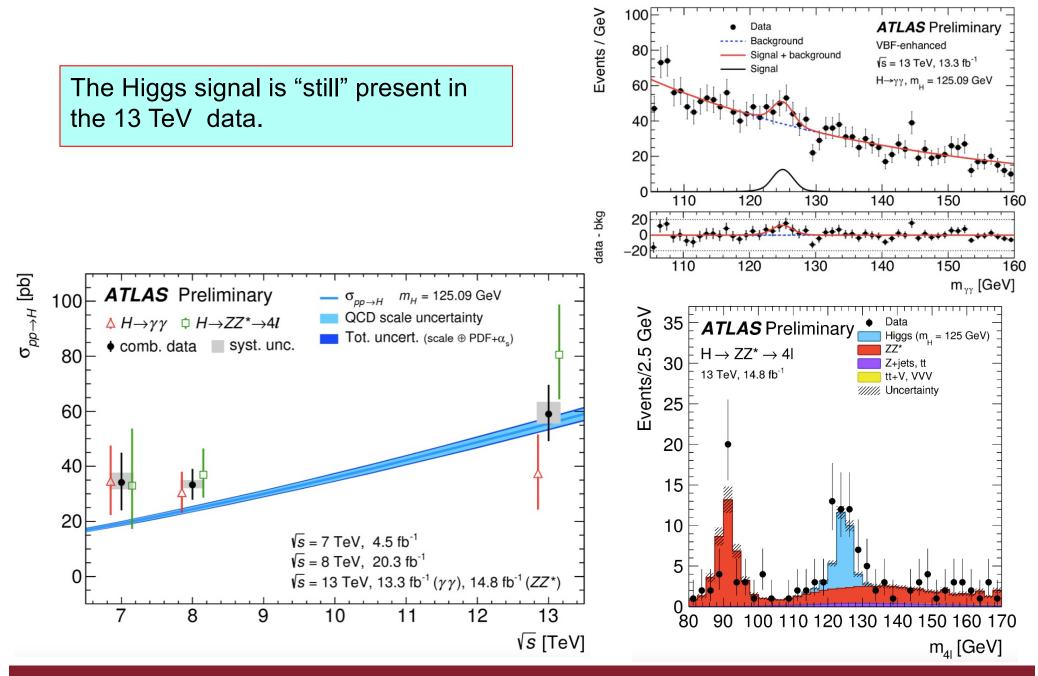
First look at 13 TeV data for H(125)



The electrons have a transverse momentum of 111 and 16 GeV, the muons 18 and 17 GeV, the jets 118 and 54 GeV. The invariant mass of the four lepton system is 129 GeV, the di-electron invariant mass is 91 GeV, the di-muon invariant mass is 29 GeV, The di-jet invariant mass is 2 TeV.

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First look at 13 TeV data for H(125)



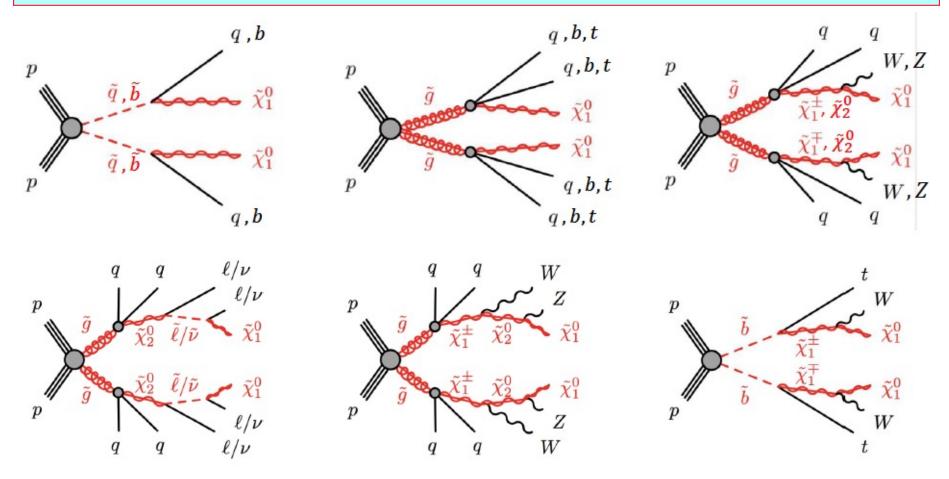
A brief introduction to SuperSymmetry

SuSy is a generalization of the SM: symmetry between fermions and bosons

- Introduces sfermions and gauginos
 - \rightarrow doubles particles content with respect to SM
- Extended Higgs sector: h, H, A, H⁺, H⁻
- PRO:
 - > Alleviates hierarchy problem ($m_h \ll m_P$)
 - has a good Dark Matter candidate (neutralino)
 - > Allows for gauge coupling unification
- CONS:
 - >Over 100 free parameters (although with some ad hoc assumptions we can reduce the number of parameters)
 - > wide range of possible experimental signatures

It was expected "something" at the TeV scale

a few diagrams with susy particles in the final state, with the decay chain

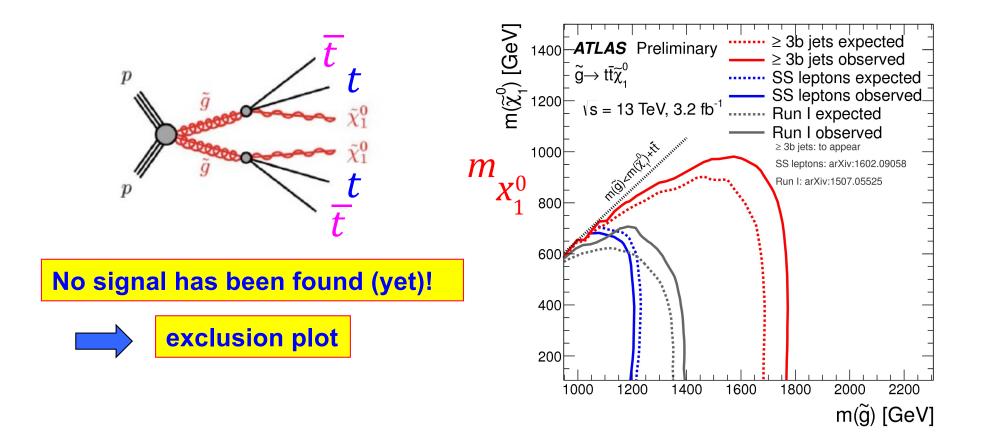


□ Lightest susy particle ($\tilde{\chi}_1^0$) escapes detection → Missing Transverse Momentum and Missing Energy

Different analysis strategies according to many different final states

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example: gaugino and neutralino mass limits



From other susy searches many exclusions limits on the parameters phase space

now there is less and less room to "manouver".

Particles masses higher and higher; cross-sections lower and lower

ATLAS susy particles: Run2 results

ATLAS SUSY Searches* - 95% CL Lower Limits Status: March 2016

ATLAS Preliminary

 $\sqrt{s} = 7, 8, 13 \text{ TeV}$

Model	e, μ, τ, γ	Jets	E _T ^{miss}	$\int \mathcal{L} dt [fb$	Mass limit	$\sqrt{s} = 7, 8 \text{ TeV}$	$\sqrt{s} = 13 \text{ TeV}$	Reference
$\begin{array}{c} & MSUGRA/CMSSM \\ \bar{q}\bar{q}, \bar{q} \rightarrow q \tilde{\chi}_{1}^{0} \\ \bar{q}\bar{q}, \bar{q} \rightarrow q \tilde{\chi}_{1}^{0} \\ \bar{q}\bar{q}, \bar{q} \rightarrow q \tilde{\chi}_{1}^{0} \\ \bar{q}\bar{q}, \bar{q} \rightarrow q \ell (\ell \ell \nu / \nu \nu) \tilde{\chi}_{1}^{0} \\ \bar{g}\bar{g}, \bar{g} \rightarrow q q \ell \tilde{\ell} + q q W^{2} \tilde{\chi}_{1}^{0} \\ \bar{g}\bar{g}, \bar{g} \rightarrow q q \ell \tilde{\ell} + q q W^{2} \tilde{\chi}_{1}^{0} \\ \bar{g}\bar{g}, \bar{g} \rightarrow q q \ell \tilde{\ell} + q q W^{2} \tilde{\chi}_{1}^{0} \\ \bar{g}\bar{g}, \bar{g} \rightarrow q q \ell \ell \ell / \nu / \nu \nu \tilde{\chi}_{1}^{0} \\ \bar{g}\bar{g}, \bar{g} \rightarrow q q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q Q \ell \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q \ell \ell + \delta \chi V \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow q \ell \ell + \delta \chi \ell + \delta \chi \ell + \delta \chi \ell + \delta \chi \ell \\ \bar{g}\bar{g}, \bar{g} \rightarrow \chi \ell + \delta \ell + \delta \chi \ell + \delta \ell + \delta \chi \ell$	0-3 $e, \mu/1-2 \tau$ 0 mono-jet 2 e, μ (off-Z) 0 1 e, μ 2 e, μ 0 1-2 $\tau + 0-1 \ell$ 2 γ γ 2 e, μ (Z) 0	2-6 jets 1-3 jets	Yes Yes Yes Yes Yes Yes Yes Yes Yes	20.3 3.2 20.3 3.2 3.3 20 3.2 20.3 20.3 2	ğ ğ <u>1.38</u> ğ 1.4	$\begin{array}{c} m(\hat{a}) \cdot m \\ m(\tilde{k}^2) = \\ 52 \ {\rm TeV} \\ {\rm I.6 \ {\rm TeV}} \\ m(\tilde{k}^2) = \\ m(\tilde{k}^2) = \\ {\rm TeV} \\ {\rm I.63 \ {\rm TeV}} \\ {\rm TeV} \\ {\rm reV} \\ {\rm m}(\tilde{k}^2) = \\ {\rm cr}({\rm NLS}) {\rm cr}({\rm nLS}) {\rm cr}({\rm cr}({\rm nLS}) {\rm cr}({\rm nLS}) {\rm cr}({\rm cr}({\rm nLS}) {\rm cr}({\rm cr}({\rm nLS}) {\rm cr}({\rm cr}$	0 GeV, m(1 st gcn. \tilde{q})=m(2 ^{md} gcn. \tilde{q}) ($\tilde{\chi}_{1}^{0}$)<5 GeV 0 GeV 0 GeV 350 GeV, m($\tilde{\chi}^{\pm}$)=0.5(m($\tilde{\chi}_{1}^{0}$)+m(\tilde{g})) 0 GeV =100 GeV	1507.05525 ATLAS-CONF-2015-062 <i>To appear</i> 1503.03290 ATLAS-CONF-2015-062 ATLAS-CONF-2015-076 1501.03555 1602.06194 1407.0603 1507.05493 1507.05493 1507.05493 1503.03290 1502.01518
$\begin{array}{c} \tilde{g}\tilde{g}, \tilde{g} \rightarrow b\bar{b}\tilde{\chi}^{0}_{1} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow t\bar{t}\tilde{\chi}^{0}_{1} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow t\bar{t}\tilde{\chi}^{0}_{1} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow b\bar{t}\tilde{\chi}^{1}_{1} \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow b\bar{t}\tilde{\chi}^{1}_{1} \end{array}$	0 0-1 <i>e</i> ,μ 0-1 <i>e</i> ,μ	3 b 3 b 3 b	Yes Yes Yes	3.3 3.3 20.1	ह हे हे 1.37	1.76 TeV m($\tilde{\chi}_1^0$)=	:800 GeV :0 GeV :300 GeV	ATLAS-CONF-2015-067 To appear 1407.0600
$\begin{array}{c} \begin{array}{c} \begin{array}{c} & & & & \\ & & & & \\ & & & & \\ & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & \\ & & & & \\ & & & & \\ & & & & \\ & & & & \\ &$	0 π 2 <i>e</i> , μ (Z) 3 <i>e</i> , μ (Z)	2 b 0-3 b 1-2 b 0-2 jets/1-2 nono-jet/c-t 1 b 1 b 6 jets + 2 b	b Yes ag Yes Yes Yes	3.2 3.2 1.7/20.3 20.3 20.3 20.3 20.3 20.3 20.3	840 GeV 51 325-540 GeV 7117-170 GeV 200-500 GeV 71 90-198 GeV 205-715 GeV 71 90-245 GeV 745-785 GeV 71 150-600 GeV 745-785 GeV 72 290-610 GeV 320-620 GeV	$m(\tilde{\chi}_{1}^{0}) = m(\tilde{\chi}_{1}^{0})$ $m(\tilde{\chi}_{1}^{\pm}) = m(\tilde{\chi}_{1}^{0}) = m(\tilde{\chi}_{1}) - n$ $m(\tilde{\chi}_{1}) - n$ $m(\tilde{\chi}_{1}^{0}) > n$	n(ℓ̃1)<85 GeV -150 GeV -200 GeV	ATLAS-CONF-2015-066 1602.09058 1209.2102, 1407.0583 08616, ATLAS-CONF-2016-0 1407.0608 1403.5222 1403.5222 1506.08616
$ \begin{array}{c} \overbrace{\begin{matrix} \overleftarrow{t}_{L,R}}{} \overbrace{\begin{matrix} \overleftarrow{t}_{L,R}}{} \overbrace{\begin{matrix} \overleftarrow{t}_{L,R}}{} \overbrace{\begin{matrix} \overleftarrow{t}_{L,R}}{} \overbrace{\begin{matrix} \overleftarrow{t}_{L,R}}{} \overbrace{\begin{matrix} \overleftarrow{t}_{L,R}}{} \overbrace{\begin{matrix} \overleftarrow{t}_{L}} \overbrace{\begin{matrix} \overleftarrow{t}_{1}}{} \\ \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \\ \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \\ \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \\ \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \end{array} \\ \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \overbrace{\begin{matrix} \overleftarrow{t}_{2}} \\ \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \overbrace{\begin{matrix} \overleftarrow{t}_{2}} \\ \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \end{array} \\ \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \overbrace{\begin{matrix} \overleftarrow{t}_{2}} \\ \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \overbrace{\begin{matrix} \overleftarrow{t}_{2}} \\ \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \end{array} \\ \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \overbrace{\begin{matrix} \overleftarrow{t}_{2}} \\ \overbrace{t}_{2}} \\ \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \overbrace{\begin{matrix} \overleftarrow{t}_{2}} \\ \overbrace{t}_{2}} \\ \overbrace{t} \end{array} \\ \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \overbrace{\begin{matrix} \overleftarrow{t}_{2}} \\ \overbrace{t} \end{array} \\ \overbrace{\begin{matrix} \overleftarrow{t}_{1}} \overbrace{\begin{matrix} \overleftarrow{t}_{2}} \\ \overbrace{t} \end{array} \\ \overbrace{t} \end{array} \\ \overbrace{t} \overbrace{t} \overbrace{t} \end{array} \\ \overbrace{t} \end{array} \\ \overbrace{t} \overbrace{t} \overbrace{t} \end{array} \\ \overbrace{t} \end{array} \\ \overbrace{t} \overbrace{t} \overbrace{t} \overbrace{t} \end{array} \\ \overbrace{t} \end{array} \\ \overbrace{t} \overbrace{t} \overbrace{t} \overbrace{t} \overbrace{t} \end{array} \\ \overbrace{t} \overbrace{t} \atop\overbrace{t} \overbrace{t} \end{array} \\ \overbrace{t} \overbrace{t} \overbrace{t} \overbrace{t} \overbrace{t} \overbrace{t} \overbrace{t} \overbrace{t}$	2 e, μ 2 e, μ 2 τ 3 e, μ 2-3 e, μ 2-3 e, μ 7/γγ e, μ, γ 4 e, μ 1 e, μ + γ	0 0 0-2 jets 0-2 b 0 -	Yes Yes Yes Yes Yes Yes Yes	20.3 20.3 20.3 20.3 20.3 20.3 20.3 20.3	7 90-335 GeV \$\$\$^1_1\$ 140-475 GeV \$	$m(\tilde{\chi}_{1}^{0}) =$ $m(\tilde{\chi}_{1}^{\pm}) = m(\tilde{\chi}_{2}^{0}), m$ $m(\tilde{\chi}_{1}^{\pm}) =$ $m(\tilde{\chi}_{1}^{\pm}) =$	$\begin{array}{l} 0 \ \mbox{GeV}, \ m(\tilde{\ell}, \bar{\nu}) {=} 0.5(m(\tilde{\chi}^+_1) {+} m(\tilde{\chi}^0_1)) \\ 0 \ \mbox{GeV}, \ m(\tilde{\tau}, \bar{\nu}) {=} 0.5(m(\tilde{\chi}^+_1) {+} m(\tilde{\chi}^0_1)) \\ (\tilde{\chi}^0_1) {=} 0, \ m(\tilde{\ell}, \bar{\nu}) {=} 0.5(m(\tilde{\chi}^+_1) {+} m(\tilde{\chi}^0_1)) \\ m(\tilde{\chi}^0_2), \ m(\tilde{\chi}^0_1) {=} 0, \ \mbox{sleptons} \ \mbox{decoupled} \\ m(\tilde{\chi}^0_1) {=} 0, \ \mbox{sleptons} \ \mbox{decoupled} \\ \tilde{\chi}^0_1 {=} 0, \ m(\tilde{\ell}, \bar{\nu}) {=} 0.5(m(\tilde{\chi}^0_2) {+} m(\tilde{\chi}^0_1)) \end{array}$	1403.5294 1403.5294 1407.0350 1402.7029 1403.5294, 1402.7029 1501.07110 1405.5086 1507.05493
$\begin{array}{c} \begin{array}{c} \text{Direct} \tilde{\chi}_{1}^{+}\tilde{\chi}_{1}^{-} \text{ prod., long-lived} \tilde{\chi}_{1}^{+}\\ \text{Direct} \tilde{\chi}_{1}^{+}\tilde{\chi}_{1}^{-} \text{ prod., long-lived} \tilde{\chi}_{1}^{+}\\ \text{Direct} \tilde{\chi}_{1}^{+}\tilde{\chi}_{1}^{-} \text{ prod., long-lived} \tilde{\chi}_{1}^{+}\\ \text{Stable, stopped} \tilde{g} \text{ R-hadron}\\ \text{Metastable } \tilde{g} \text{ R-hadron}\\ \text{GMSB, } \tilde{\chi}_{1}^{0} \rightarrow \tilde{\gamma}(\tilde{e}, \tilde{\mu}) + \tilde{\gamma}(\tilde{e}, \mu$	dE/dx trk 0 dE/dx trk		Yes Yes - - Yes - -	20.3 18.4 27.9 3.2 19.1 20.3 20.3 20.3	Image: style	$\begin{array}{c} m(\tilde{k}_{1}^{*}) - \\ m(\tilde{k}_{1}^{0}) = \\ \textbf{1.54 TeV} \\ \textbf{m}(\tilde{k}_{1}^{0}) = \\ 10 < tar \\ 1 < \tau(\tilde{k}_{1}^{0}) \\ \tau < \tau(\tilde{k}_{1}^{0}) \\ \textbf{7} < c\tau(\tilde{k}_{1}^{0}) \\ \textbf{7} < c\tau(\tilde{k}^{0}) \\ \textbf{7} < c\tau(\tilde$	$\begin{split} & m(\tilde{k}^0) \!$	1310.3675 1506.05332 1310.6584 <i>To appear</i> 1411.6795 1409.5542 1504.05162 1504.05162
$ \begin{array}{c} & \begin{array}{c} & \mbox{LFV} pp \rightarrow \tilde{v}_{\tau} + X, \tilde{v}_{\tau} \rightarrow e\mu/e\tau/\mu\tau \\ \mbox{Bilinear RPV CMSSM} \\ & \tilde{\chi}^{\dagger}_{1}\tilde{\chi}_{1}, \tilde{\chi}^{\dagger}_{1} \rightarrow W\tilde{V}^{\dagger}_{1}, \tilde{\chi}^{0}_{1} \rightarrow ee\tilde{v}_{\mu}, e\mu\tilde{v}_{e} \\ & \tilde{\chi}^{\dagger}_{1}\tilde{\chi}^{-}_{1}, \tilde{\chi}^{+}_{1} \rightarrow W\tilde{V}^{\dagger}_{1}, \tilde{\chi}^{0}_{1} \rightarrow \tau r\tilde{v}_{e}, er\tilde{v}_{\tau} \\ & \tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}^{0}_{1}, \tilde{\chi}^{0}_{1} \rightarrow \tau r\tilde{v}_{e}, er\tilde{v}_{\tau} \\ & \tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}^{0}_{1}, \tilde{\chi}^{0}_{1} \rightarrow qqq \\ & \tilde{g}\tilde{g}, \tilde{g} \rightarrow t_{1}, \tilde{t}_{1} \rightarrow bs \\ & \tilde{t}_{1}\tilde{t}_{1}, \tilde{t}_{1} \rightarrow b\ell \end{array} $	0 0 2 <i>e</i> , µ (SS)	- 0-3 b - - 6-7 jets 6-7 jets 0-3 b 2 jets + 2 b 2 b	- Yes Yes - - Yes b -	20.3 20.3 20.3 20.3 20.3 20.3 20.3 20.3	\$\vec{r}_r\$ 760 GeV \$\vec{r}_1\$ 760 GeV \$\vec{r}_1\$ 450 GeV \$\vec{r}_2\$ 917 GeV \$\vec{r}_2\$ 980 GeV \$\vec{r}_2\$ 880 GeV \$\vec{r}_1\$ 320 GeV \$\vec{r}_1\$ 0.4-1.0 TeV	15 TeV $m(\tilde{q})$ =n $m(\tilde{\chi}_1^0)$ > $m(\tilde{\chi}_1^0)$ > $m(\tilde{\chi}_1^0)$ > $m(\tilde{\chi}_1^0)$ > BR(t)= $m(\tilde{\chi}_1^0)$ =	11, $\lambda_{132/133/233}=0.07$ $n(\tilde{g}), c\tau_{LSP} < 1 \text{ mm}$ $0.2 \times m(\tilde{K}_1^+), \lambda_{121} \neq 0$ $0.2 \times m(\tilde{K}_1^+), \lambda_{133} \neq 0$ BR(b)=BR(c)=0% 600 GeV $+be/\mu)>20\%$	1503.04430 1404.2500 1405.5086 1502.05686 1502.05686 1502.05686 1404.2500 1601.07453 ATLAS-CONF-2015-015
			Yes	20.3	510 GeV		200 GeV	1501.01325

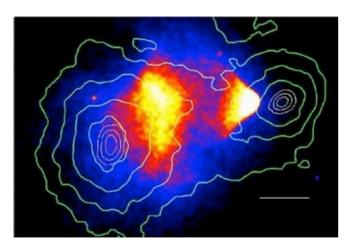
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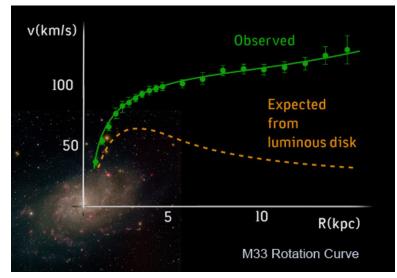
Evidence for Dark Matter



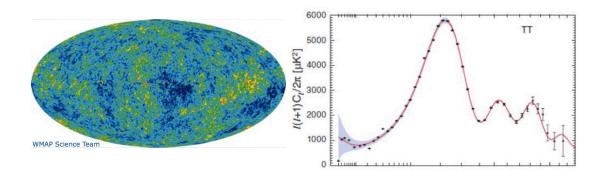
Comprises **majority** of **mass** in Galaxies Missing mass on Galaxy Cluster scale Zwicky (1937)



Almost collisionless Bullet Cluster Clowe+(2006)

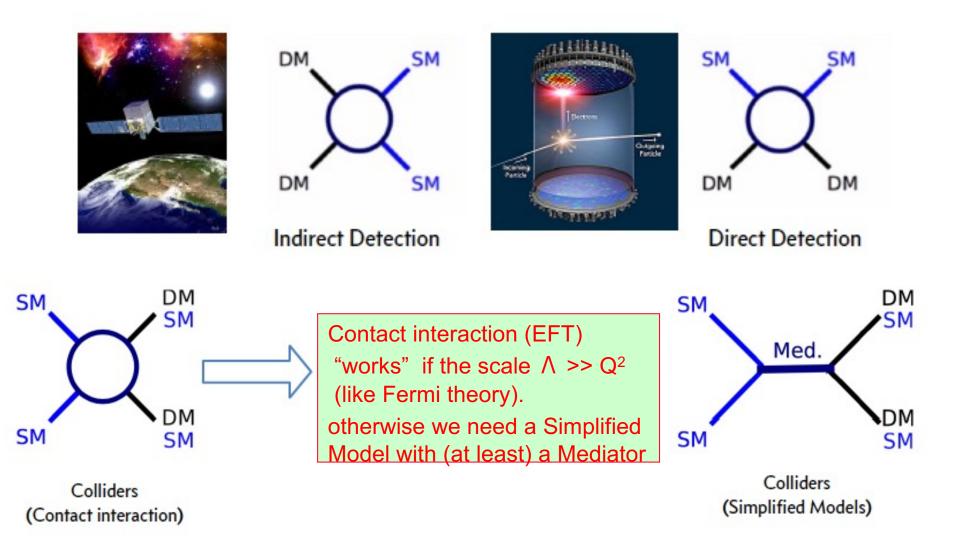


Large **halos** around Galaxies Rotation Curves Rubin+(1980)



Non-BaryonicBig-Bang Nucleosynthesis,CMBAcousticOscillationsWMAP(2010),Planck(2015)VMAP(2015)

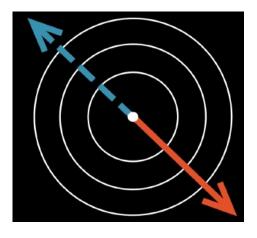
Assumption: non-gravitational interaction with ordinary matter



Detecting Dark Matter at LHC

Non-interacting DM particles → Missing transverse energy (MET)

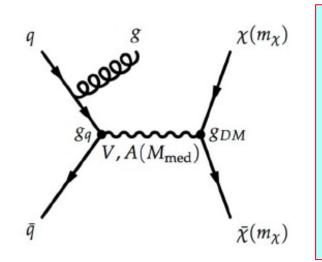
1/3



X (jet, photon, etc..)

General analysis strategy

- Require MET
- Select for X
- Veto other objects
- Additional cuts to suppress background
- Data-driven techniques to estimate background → invert vetoes



Results are interpreted in the Simplified Model framework to allow comparison with Direct Detection

- Mediator particle connects the SM quarks to DM particles:
 - Axial Vector, Pseudoscalar, etc...
- Model depends on four parameters:
 - DM mass, Mediator mass, SM-mediator coupling, DM-mediator coupling

DM at ATLAS; one example: monojet

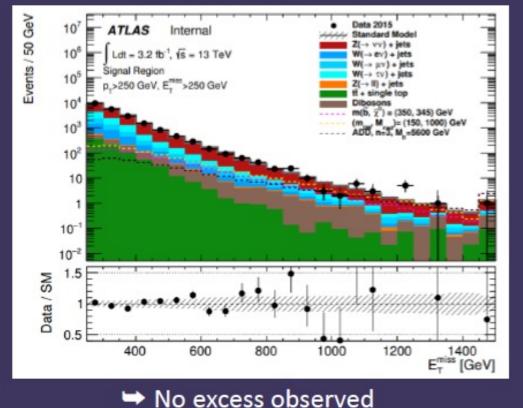
Backgrounds

- Main backgrounds are EW processes with intrinsic E_T^{miss}, accompanied by jets:
 - Z(vv)+jets: irreducible background
 - W(lv)+jets: with unrecostructed or misidentified lepton
- Both estimated from data using leptonic Z or W control regions

Other backgrounds:

- Non-collision background (data)
- Multijet background (data)
- ➤ Z→ee, top, diboson (MC)

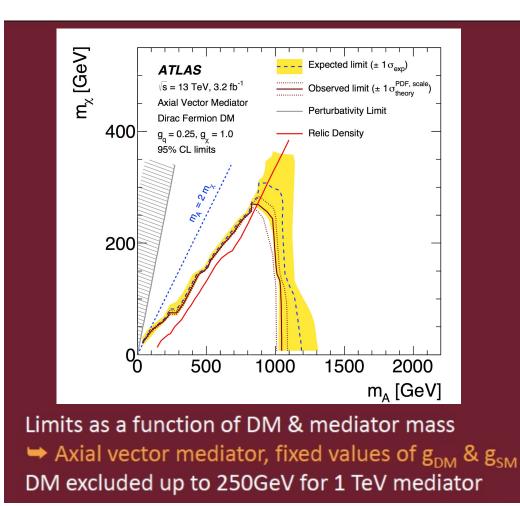
Signal Region MET

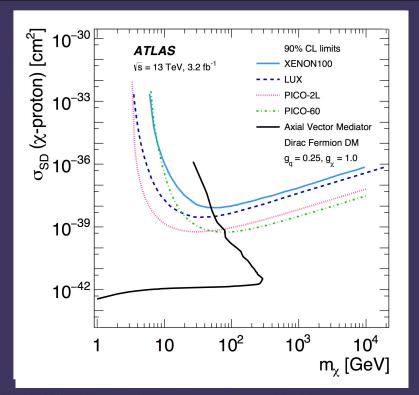


Dominant uncertainties: Statistical (3-10%), top (~3%), boson+jet modeling (2-4%)

DM at ATLAS; one example: monojet

Results





LHC limits reinterpreted as limit on DM-proton scattering cross-section LHC complementary at low m_{DM}

Parameter values & limit interpretation as recommended by the LHC Dark Matter Working Group [ArXiv:1603.04156]



End of chapter 13

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