



The Approaching future of Leptonic Mixing

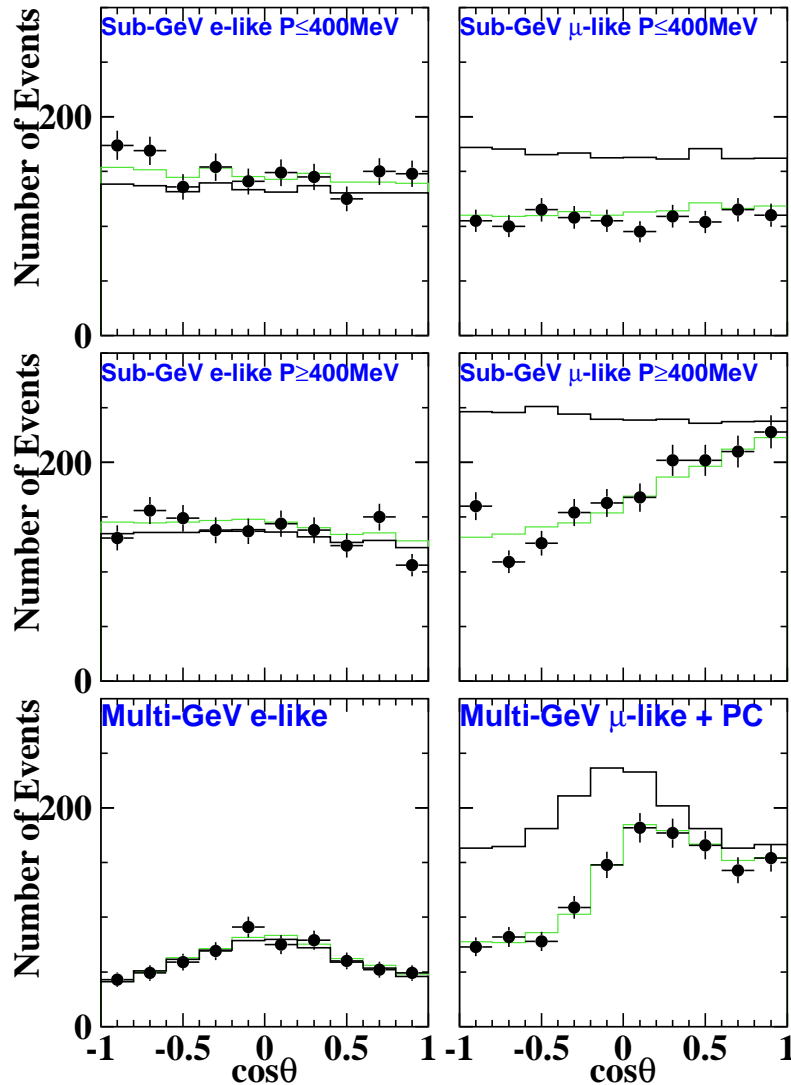
El jardín de los caminos que se bifurcan

A. Donini

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The SuperK atmospheric data



1998:
First strong evidence
of ν_μ oscillations.

Leptonic flavour mixing

The Pontecorvo-Maki-Nakagawa-Sakata matrix:

$$\begin{pmatrix} \text{FlavourBasis} \\ |\nu_\alpha\rangle \end{pmatrix} = \sum_i U_{\text{PMNS}}^{\alpha i} \times \begin{pmatrix} \text{MassBasis} \\ |\nu_i\rangle \end{pmatrix}$$

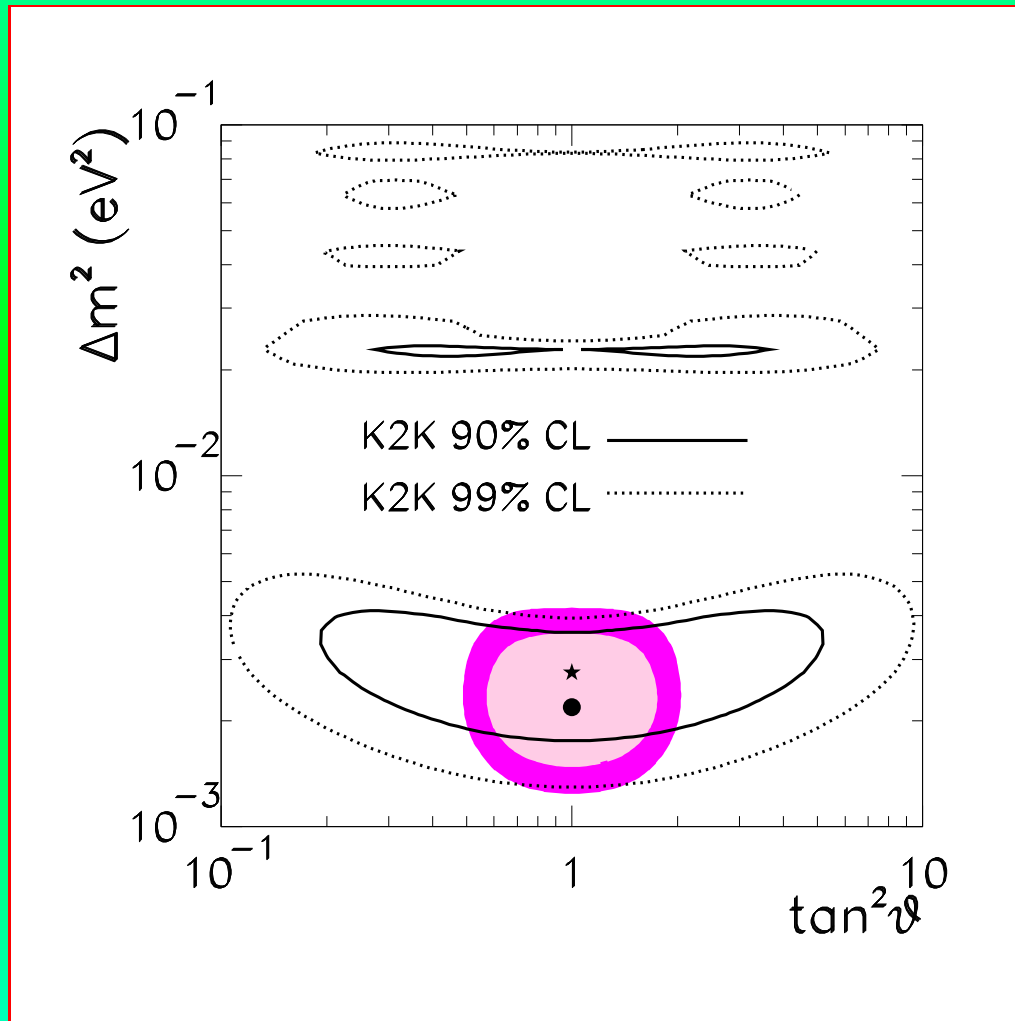
and its space-evolution:

$$i \frac{d}{dx} |\nu_\alpha\rangle = \sum_\beta \left(U \left[\frac{\Delta m^2}{2E_\nu} \right] U^\dagger + [A] \right)_{\alpha\beta} |\nu_\beta\rangle$$

$$\left[\frac{\Delta m^2}{2E_\nu} \right] = \frac{1}{2E_\nu} \text{diag}(0, \Delta m_{21}^2, \Delta m_{31}^2)$$

$$[A] = \text{diag}(\sqrt{2}G_F N_e, 0, 0)$$

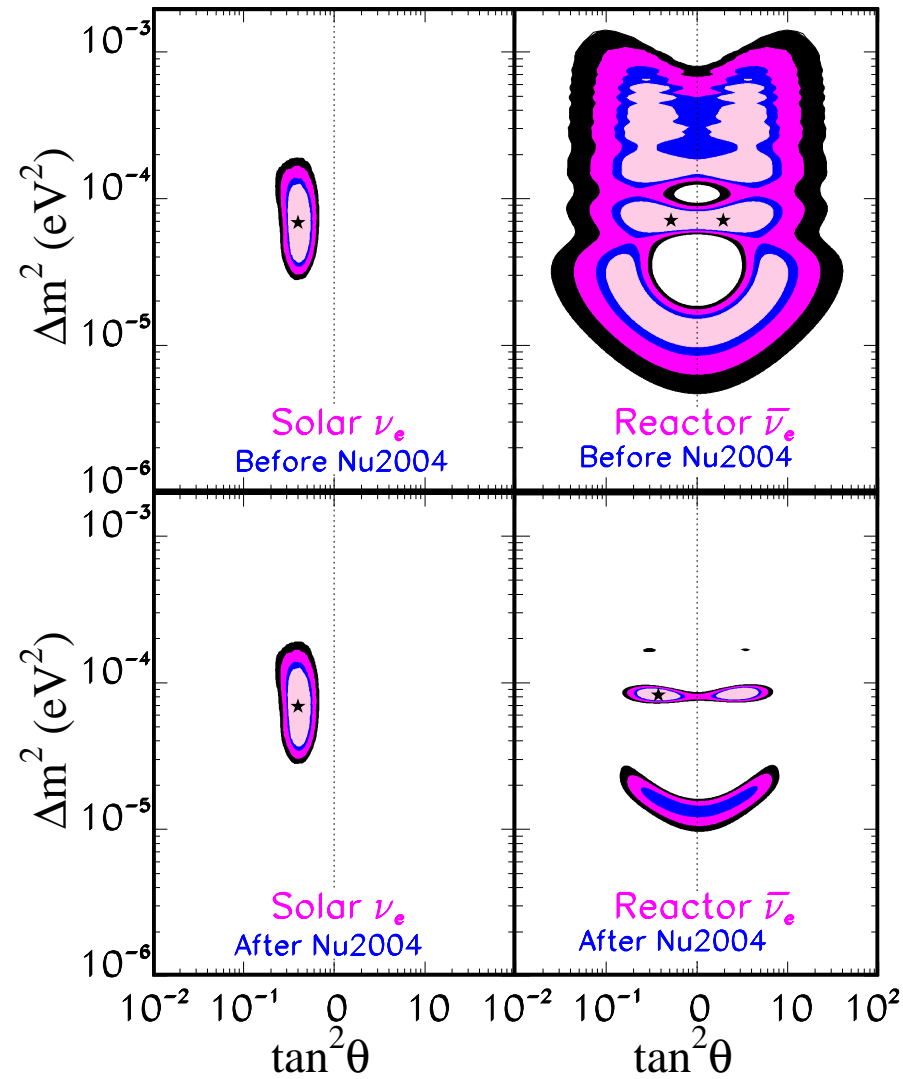
The atmospheric sector



Atmospheric data
plus
K2K

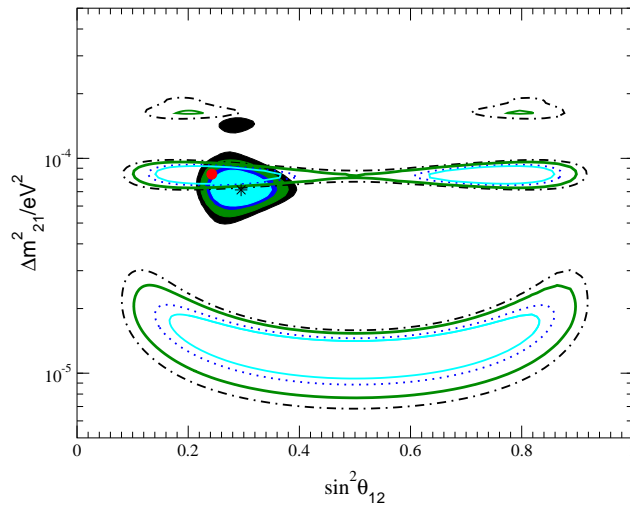
a conventional beam:
 $\pi^\pm \rightarrow \mu^\pm \nu_\mu (\bar{\nu}_\mu)$

The solar sector



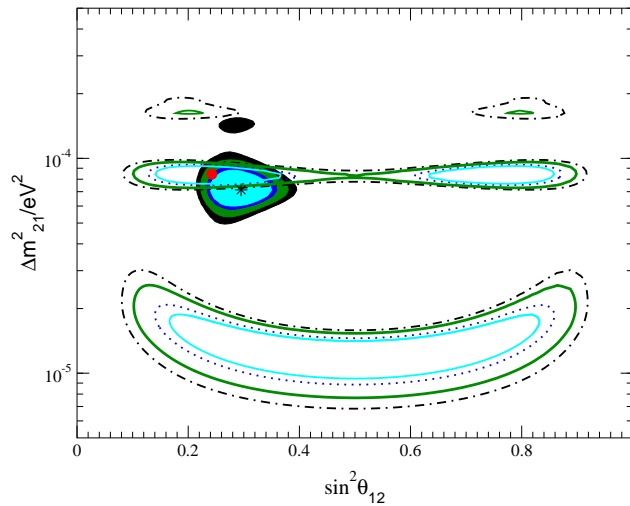
The solar sector, 2

Solar + KamLand (old)

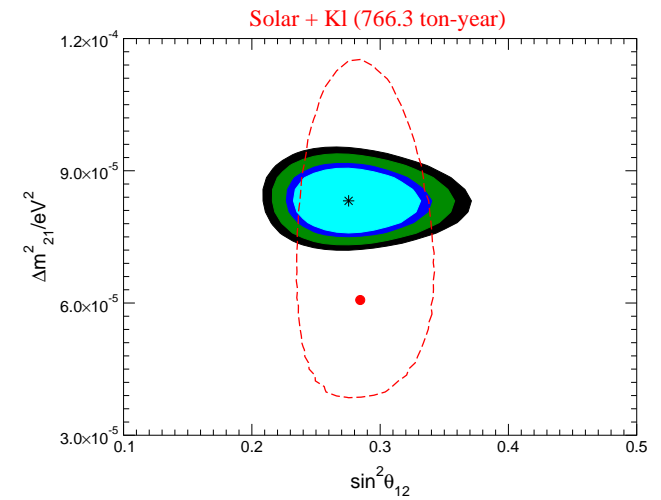


The solar sector, 2

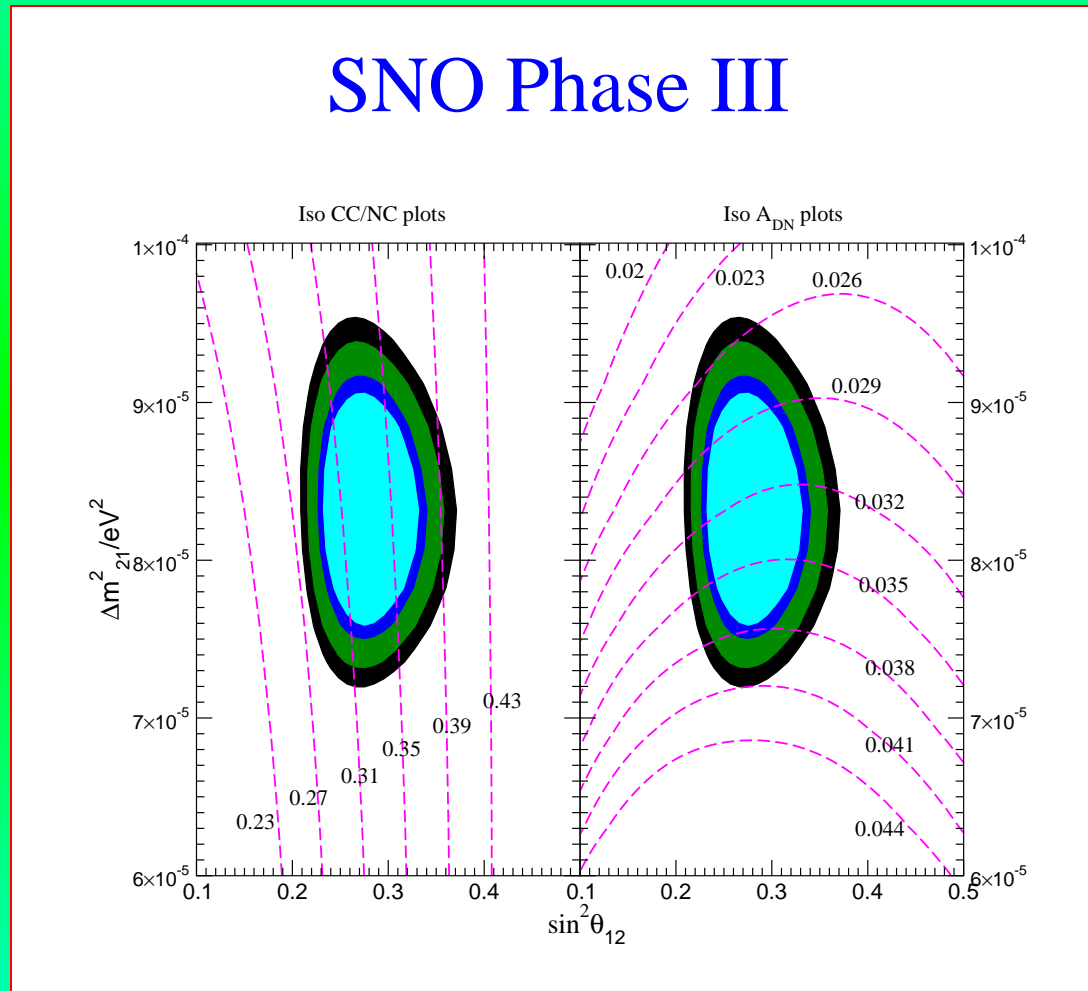
Solar + KamLand (old)



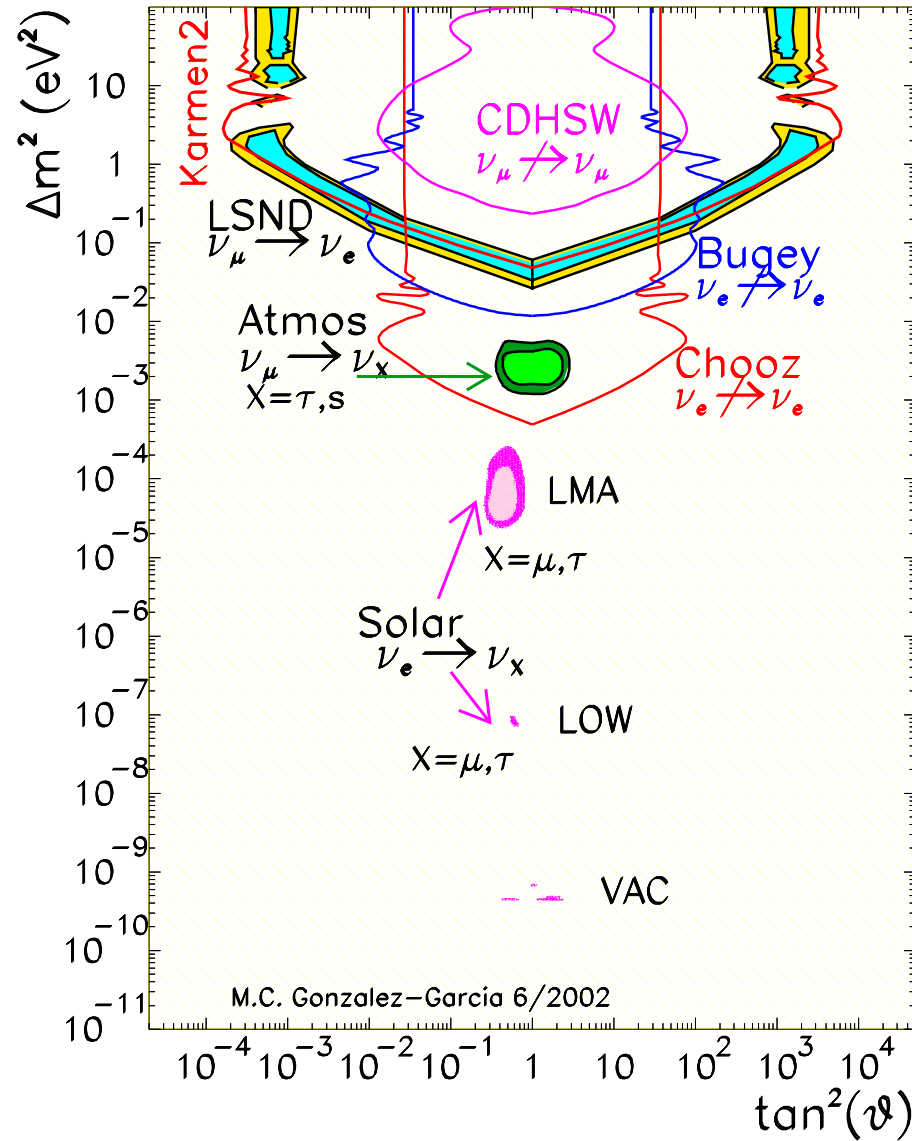
Solar + KamLand (new)



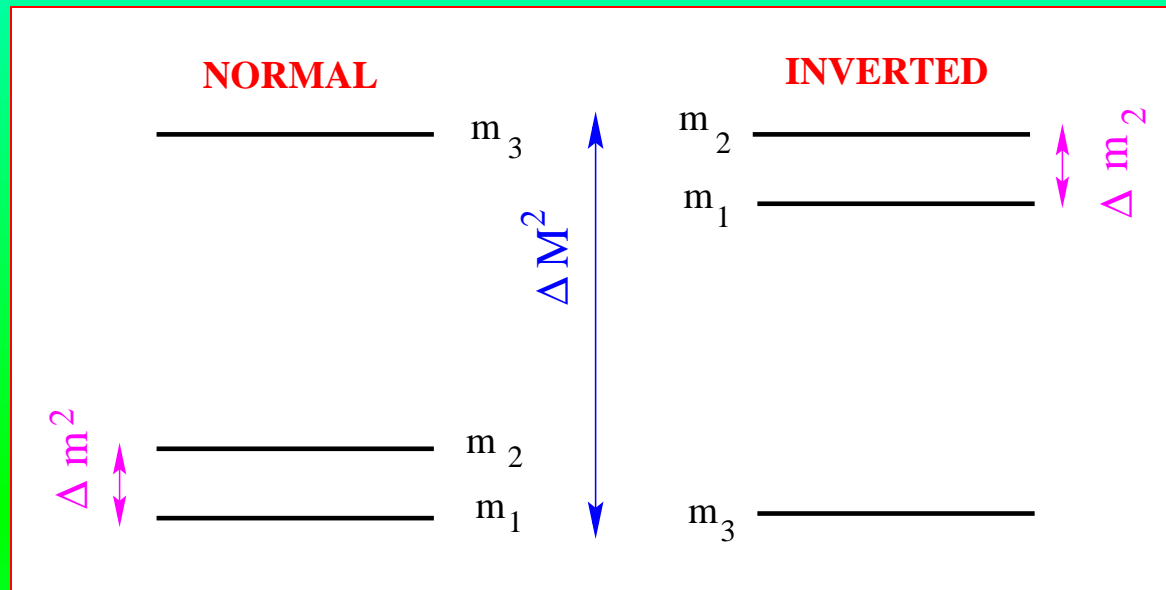
The solar sector, 3



The LSND puzzle



The mass hierarchy



Present data do not distinguish between
normal and **inverted** hierarchy

The PMNS matrix

- $|\Delta m_{atm}^2| = (2.6 \pm 0.4) \cdot 10^{-3} \text{ eV}^2$
- $\Delta m_{sol}^2 = (8.3 \pm 0.4) \cdot 10^{-5} \text{ eV}^2$
- $\sin^2 \theta_{23} = (0.33 - 0.68) \text{ at } 3\sigma$
- $\sin^2 \theta_{12} = (0.22 - 0.38) \text{ at } 3\sigma$
- $\sin^2 \theta_{13} \leq 0.041 \text{ at } 3\sigma$
- The sign of Δm_{atm}^2
- The θ_{23} -octant
- Is θ_{13} different from zero?
- Is δ different from zero?

Two more missing pieces

- The absolute neutrino mass:
tritium decay experiments;
cosmological bounds (WMAP, PLANCK, ...)
- Majorana or Dirac neutrinos?
neutrinoless double- β decay experiments;
plenty of additional phases,
needed for **LEPTOGENESIS**.

I will not cover these two items in this talk.

Neutrino sources

- Natural sources:
 - ★ The Sun $\longrightarrow \nu_e$
 - ★ Cosmic rays $\longrightarrow \nu_e, \nu_\mu$
 - ★ Supernovae and relic SNs $\longrightarrow \nu_e, \nu_\mu, \nu_\tau$
(wait for gadolinium SK!)
 - ★ Geoneutrinos !!! $\longrightarrow \nu_e$

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(wait for gadolinium SK!)

- ★ Geoneutrinos !!! $\longrightarrow \nu_e$

- Man-made sources:

- △ Reactors: $n \rightarrow pe \bar{\nu}_e$

- △ Conventional beams: $\pi^\pm \rightarrow \mu^\pm \nu_\mu (\bar{\nu}_\mu)$

- △ Neutrino Factory: $\mu^\mp \rightarrow e^\mp \nu_\mu, \bar{\nu}_e (\bar{\nu}_\mu, \nu_e)$

- △ Beta beam: ${}^6\text{He} \rightarrow \bar{\nu}_e, {}^{18}\text{Ne} \rightarrow \nu_e$

Foreseen bounds on θ_{13}

EXP	θ_{13}	$\sin^2(2\theta_{13})$	$\sin^2 \theta_{13}$
Global Fit	11.5°	0.157	0.041
BEAMS			
K2K	?	?	?
MINOS	6°	0.04	0.01
	→ 8°	→ 0.08	→ 0.02
CNGS	5°	0.03	0.008
	→ 7°	→ 0.06	→ 0.015

$$P_{\mu\mu} \simeq$$

$$1 - \sin^2(2\theta_{23}) \sin^2 \left[\frac{\Delta_{atm} L}{2} \right] + \mathcal{O} \left[\left(\frac{\Delta_{sol}}{\Delta_{atm}} \right) \sin \theta_{13} \cos \delta \right]$$

Sensitivity loss due to $(\theta_{13} - \delta)$ -correlations

Foreseeable bounds on θ_{13} (1)

EXP	θ_{13}	$\sin^2(2\theta_{13})$	$\sin^2 \theta_{13}$
Global Fit	11.5°	0.157	0.041
SBEAMS			
JHF-I	2.2°	0.006	0.0015
(T2K)	→ 3.3°	→ 0.013	→ 0.0030
NUMI-OA	2°	0.005	0.0010
(NO ν A)	→ 3.5°	→ 0.015	→ 0.0040
SPL	1.2°	0.003	0.0007
	→ 2.9°	→ 0.010	→ 0.0025

Sensitivity loss due to $(\theta_{13} - \delta)$ -correlations

Foreseeable bounds on θ_{13} (2)

EXP	θ_{13}	$\sin^2(2\theta_{13})$	$\sin^2 \theta_{13}$
Global Fit	11.5°	0.157	0.041
REACT.			
Japan	4.5°	0.025	0.006
USA	3.5°	0.015	0.004
EU (D-CHOOZ)	5°	0.030	0.008

$$P_{ee} \simeq 1 - \sin^2(2\theta_{13}) \sin^2 \left[\frac{\Delta_{atm} L}{2} \right] + \mathcal{O} \left[\left(\frac{\Delta_{sol}}{\Delta_{atm}} \right)^2 \right]$$

no sensitivity loss due to $(\theta_{13} - \delta)$ -correlations

Around 2012...

After the wave of conventional beams and first generation superbeams, and of high-power reactors experiments, we will know something more on the PMNS matrix:

- ▷ mass differences $\Delta m_{atm}^2, \Delta m_{sol}^2$ at some %;

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- ▷ the value of θ_{13} , if large;
- ▷ the sign of Δm_{23}^2 , if lucky.

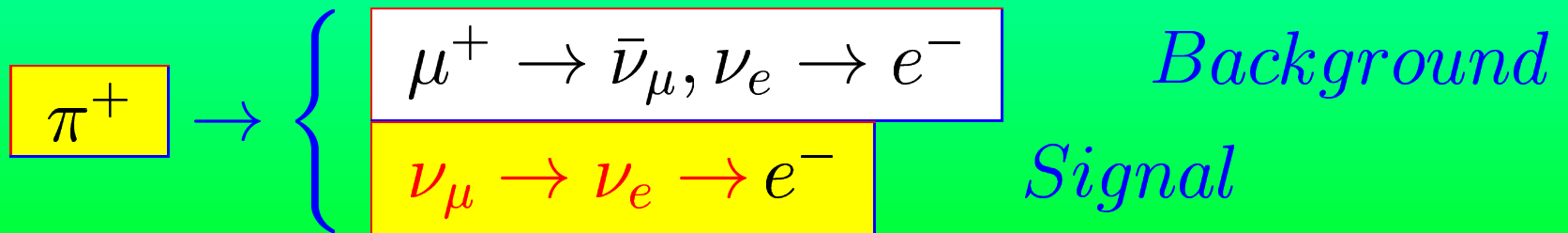
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- ▷ the sign of Δm_{23}^2 , if lucky.

Precision measurements of the **LEPTONIC MIXING** will start with the next-to-next generation man-made sources, the Neutrino Factory and/or the Beta Beam.

Appearance Signal at a SB



The oscillation probability is

$$P_{\mu e}^\pm \simeq X_\pm \sin^2(2\theta_{13})$$

$$+ Y_\pm \cos\left(\delta \pm \frac{\Delta_{atm} L}{2}\right) \cos\theta_{13} \sin(2\theta_{13})$$

$$+ Z + \dots$$

The Coefficients

For the golden channel we have:

$$\begin{cases} X_{\pm} &= \Delta_{atm}^2 \times f_X^{\pm}(\theta_{23}, A, L, E_{\nu}) \\ Y_{\pm} &= \Delta_{sol} \times \Delta_{atm} \times f_Y^{\pm}(\theta_{12}, \theta_{23}, A, L, E_{\nu}) \\ Z &= \Delta_{sol}^2 \times f_Z(\theta_{12}, \theta_{23}, A, L, E_{\nu}) \end{cases}$$

(+ neutrinos, - antineutrinos)

The (θ_{13}, δ) correlation

The number of signal electrons is:

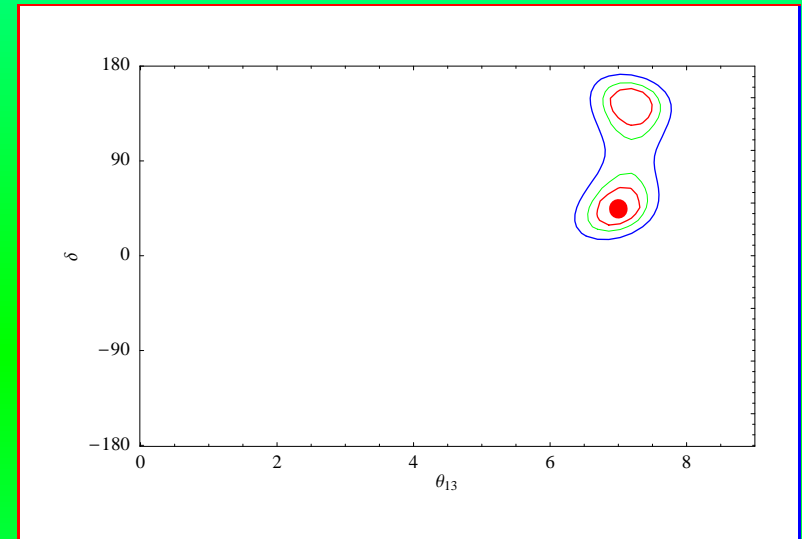
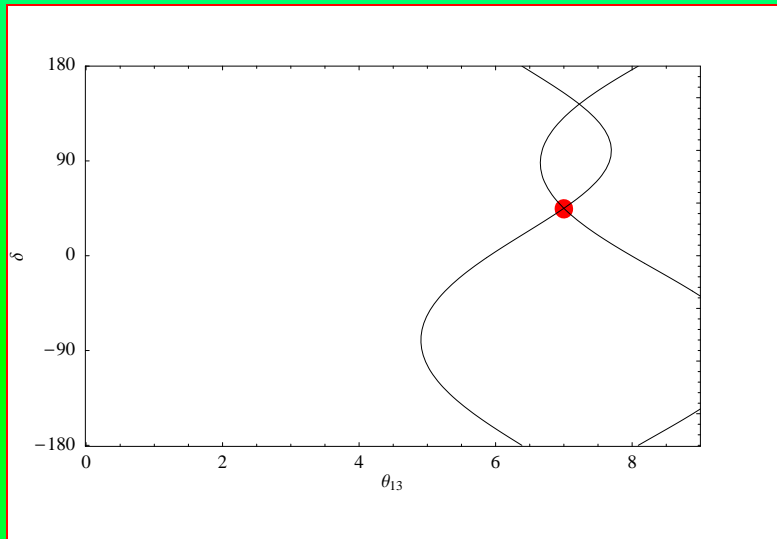
$$N_{e^-}(\bar{\theta}_{13}, \bar{\delta}) = \left\{ \epsilon_e \otimes \sigma_{\nu_e} \otimes P_{\mu e}^+(\bar{\theta}_{13}, \bar{\delta}) \otimes \Phi_{\nu_\mu} \right\}_E^{E+\Delta E}$$

$$N_{\pm}^i(\bar{\theta}_{13}, \bar{\delta}) = N_{\pm}^i(\theta_{13}, \delta)$$

By changing (θ_{13}, δ) accordingly,
curves are drawn in the (θ_{13}, δ) plane.

Degeneracy in (θ_{13}, δ) at the SPL

2 years for π^+ and 8 years for π^-



$$L = 130 \text{ Km}, \bar{E}_{\nu_{\mu}} = 0.27 \text{ GeV}, \bar{E}_{\bar{\nu}_{\mu}} = 0.25 \text{ GeV}$$

Input parameters: $\bar{\theta}_{13} = 7^{\circ}, \bar{\delta} = 45^{\circ}$

The (θ_{13}, δ) correlation (2)

The number of signal electrons is:

$$N_{e^-}(\bar{\theta}_{13}, \bar{\delta}) = \left\{ \epsilon_e \otimes \sigma_{\nu_e} \otimes P_{\mu e}^+(\bar{\theta}_{13}, \bar{\delta}) \otimes \Phi_{\nu_\mu} \right\}_E^{E+\Delta E}$$

$$N_{\pm}^i(\bar{\theta}_{13}, \bar{\delta}, \bar{s}_{atm}, \bar{s}_{oct}) = N_{\pm}^i(\theta_{13}, \delta, s_{atm}, s_{oct})$$

where

$$\begin{cases} s_{atm} = \text{sign}(\Delta m_{atm}^2) = \pm 1 \\ s_{oct} = \text{sign}(\tan 2\theta_{23}) = \pm 1 \end{cases}$$

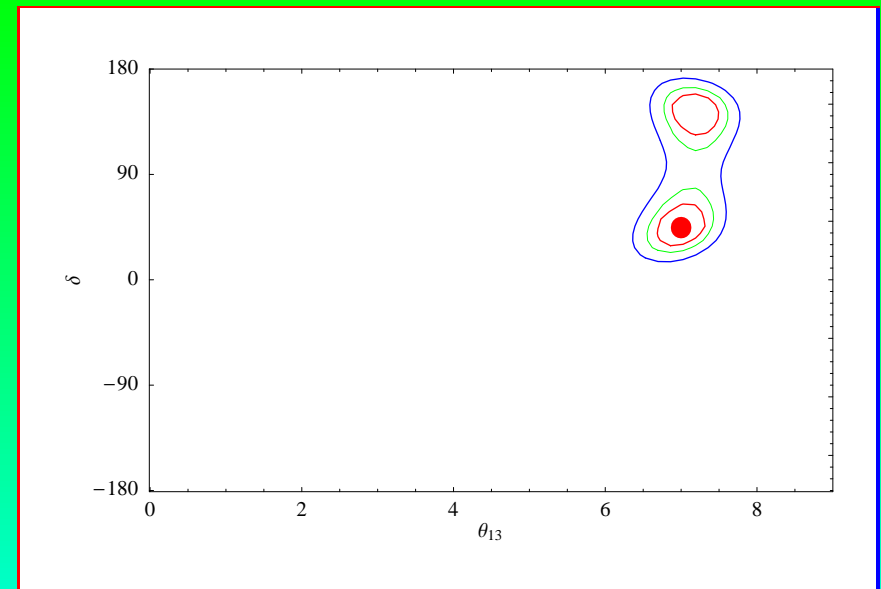
Of other clones

As a first step:

- ▷ $\theta_{23} = 45^\circ$
- ▷ Sign of Δ_{atm} fixed

J. Burguet-Castell *et al.*, hep-ph/0103258

The intrinsic clone



Of other clones

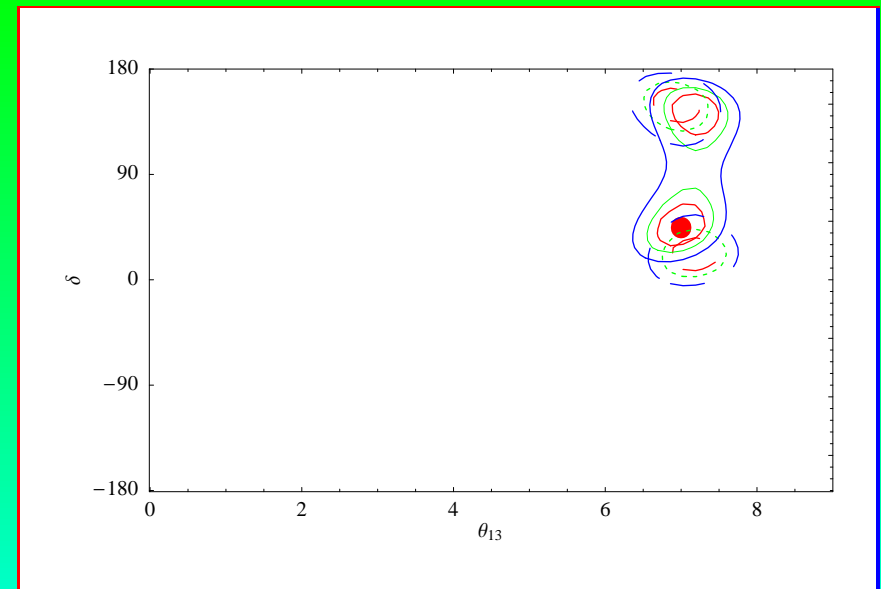
As a third step:

- ▷ $\theta_{23} = 45^\circ$
- ▷ Sign of Δ_{atm} variable

H. Minakata, H. Nunokawa, hep-ph/0108085

One more ambiguity:

- ▷ the **sign** clone



Of other clones

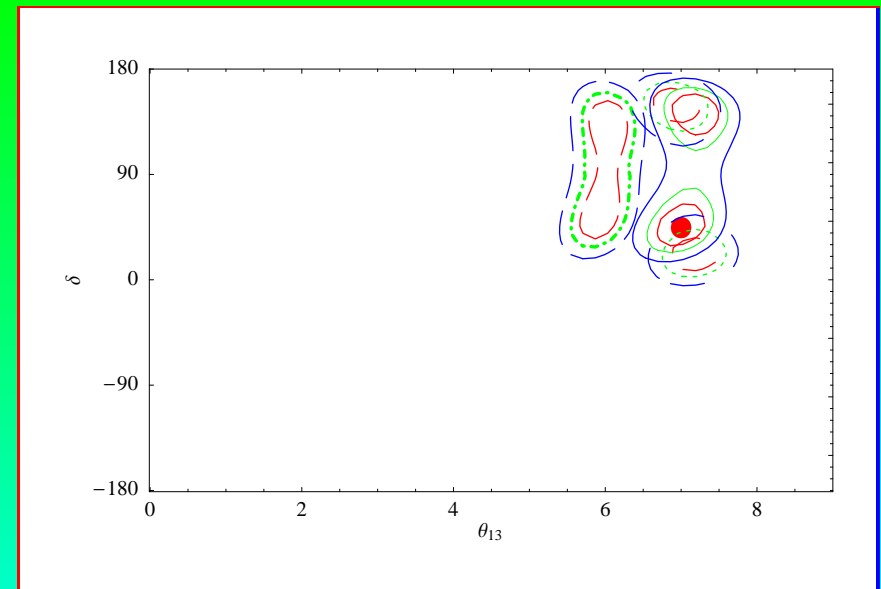
As a second step:

- ▷ $\theta_{23} \neq 45^\circ$
- ▷ Sign of Δ_{atm} fixed

G.L. Fogli, E. Lisi, hep-ph/9604415

Two more ambiguities:

- ▷ the **octant** clone
- ▷ the **sign** clone



Of other clones

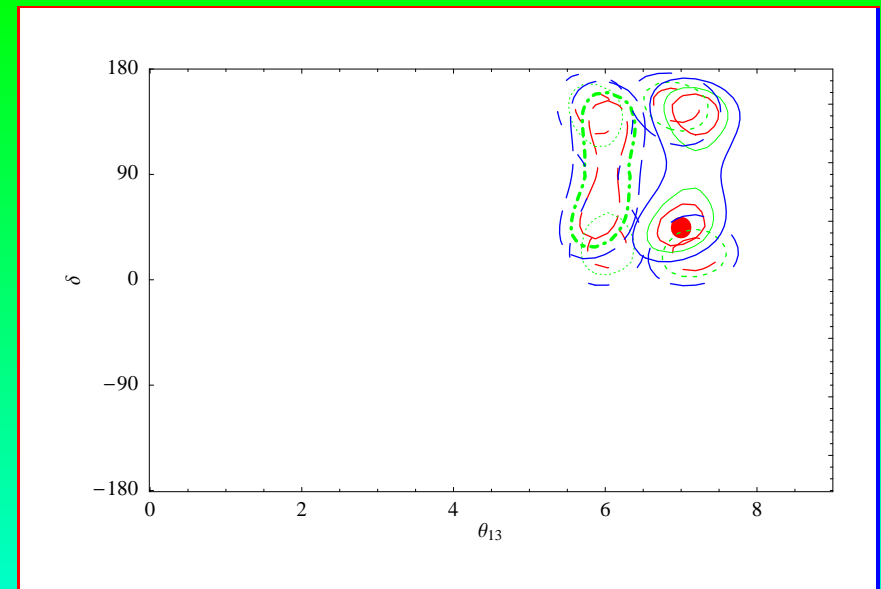
As a fourth step:

- ▷ $\theta_{23} \neq 45^\circ$
- ▷ Sign of Δ_{atm} variable

V. Barger *et al.*, hep-ph/0112119

Three more ambiguities:

- ▷ the octant clone
- ▷ the sign clone
- ▷ the mixed clone



Solving degeneracies

It is possible to compute analytically the location of each clone in the (θ_{13}, δ) plane as a function of

- the flux and the beam characteristics
- the input parameters $(\bar{\theta}_{13}, \bar{\delta})$

A.Donini, D.Meloni and S.Rigolin, hep-ph/0312072

E.g., for the intrinsic clones in vacuum:

$$\sin^2 2\theta_{13} = \sin^2 2\bar{\theta}_{13} + \frac{1}{D^2} [1 + 2D \cos \bar{\delta} \sin 2\bar{\theta}_{13}]$$

These results can be used, in principle, to **design experiments with complementary characteristics** to solve the parameter degeneracies.

The understanding of the clones location permits to study a SETUP to solve most of the degeneracies.

Two possible setups:

- ▷ The Neutrino Factory
- ▷ The Beta Beam

Both facilities are multi-sources plants, including:

- one SuperBeam facility
- two μ or heavy ion decay tunnels

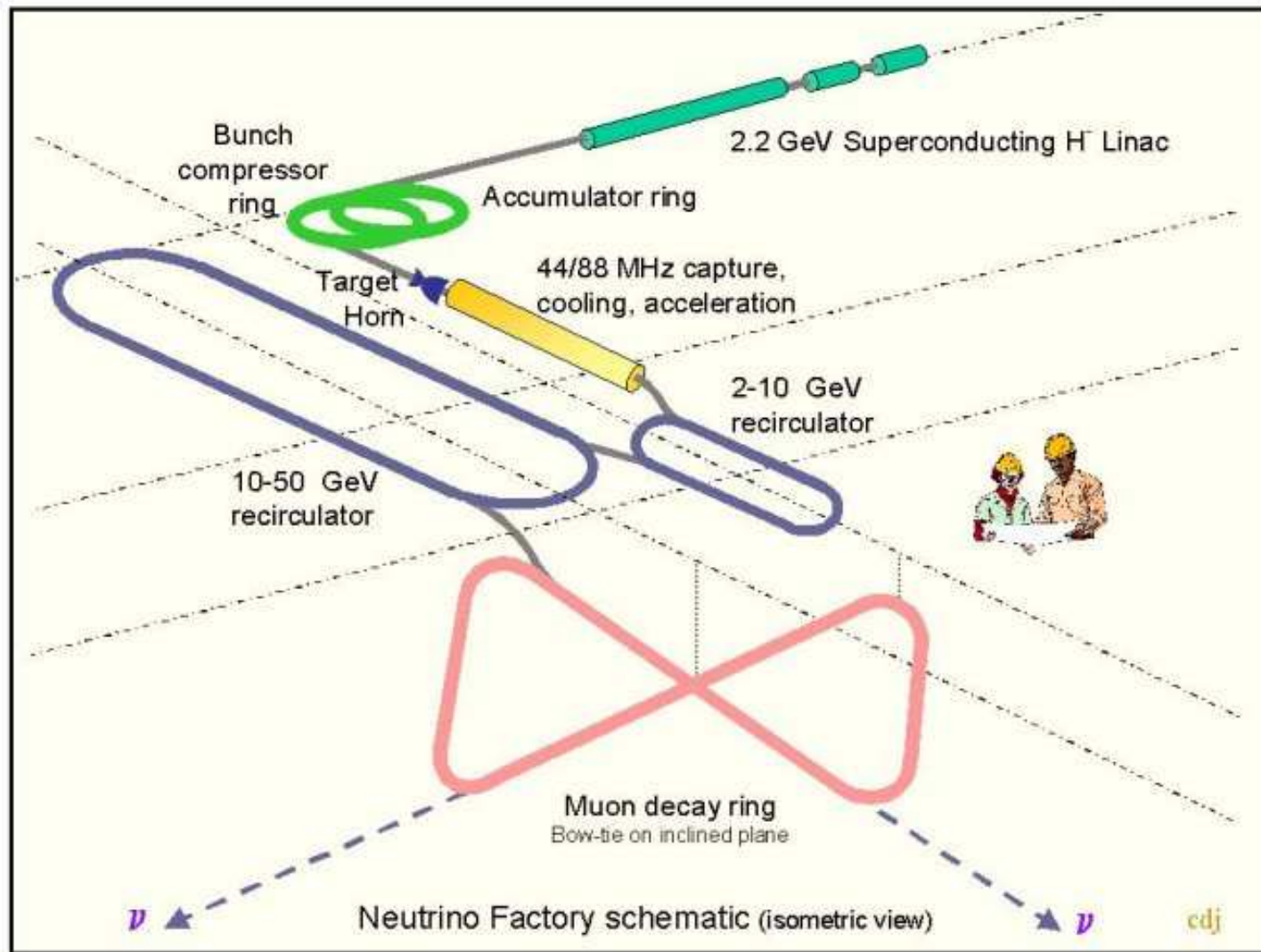
Neutrino Factory setup:

A. Donini, hep-ph/0310014; NuFact03, New York

BetaBeam setup:

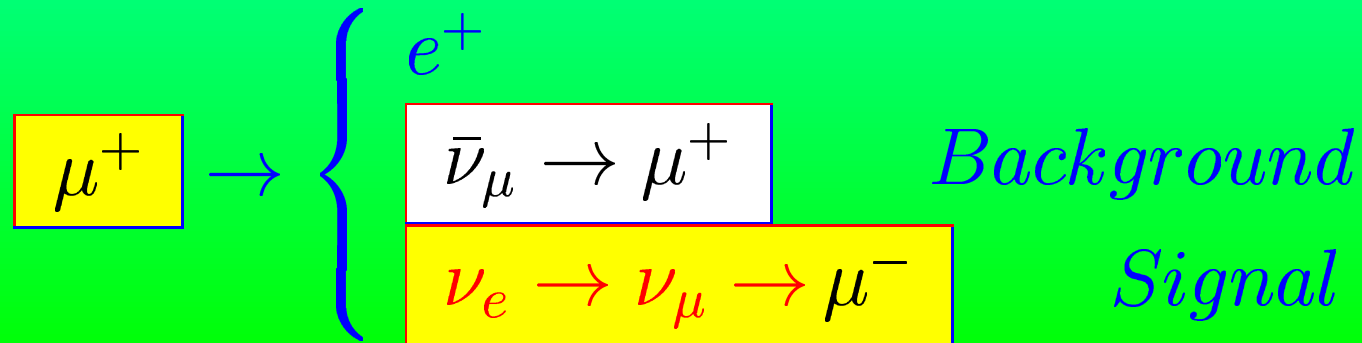
J. Bouchez *et al.*, hep-ph/0310059

The Neutrino Factory at CERN



The Golden channel: ν -factory

A. Cervera *et al.*, hep-ph/0002108

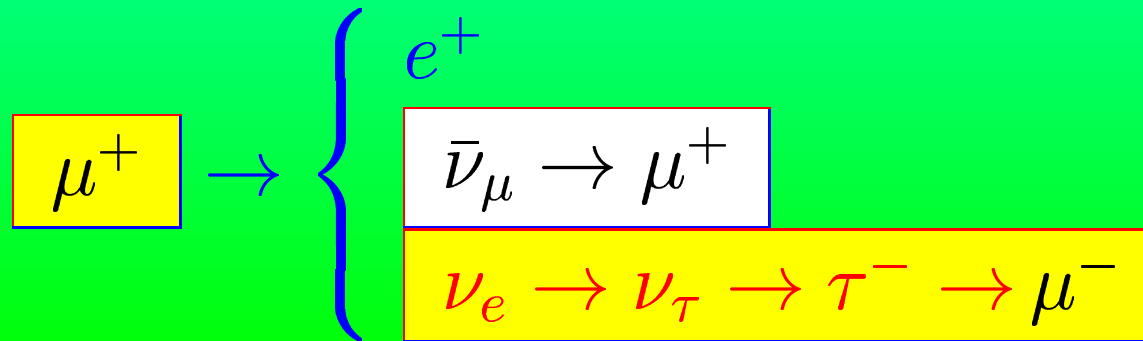


The oscillation probability is

$$P_{e\mu}^\pm = X_\pm \sin^2(2\theta_{13}) + Y_\pm \cos\left(\delta \mp \frac{\Delta_{atm}L}{2}\right) \cos\theta_{13} \sin(2\theta_{13}) + Z + \dots$$

The Silver channel: ν -factory

A. Donini, D. Meloni and P. Migliozzi, hep-ph/0206034



The oscillation probability is

$$P_{e\tau}^\pm = X_\pm^\tau \sin^2(2\theta_{13})$$

$$-Y_\pm^\tau \cos\left(\delta \mp \frac{\Delta_{atm}L}{2}\right) \cos\theta_{13} \sin(2\theta_{13}) \\ +Z^\tau + \dots$$

The Coefficients

For the golden channel we have:

$$\begin{cases} X_{\pm} &= \Delta_{atm}^2 \times f_X^{\pm}(\theta_{23}, A, L, E_{\nu}) \\ Y_{\pm} &= \Delta_{sol} \times \Delta_{atm} \times f_Y^{\pm}(\theta_{12}, \theta_{23}, A, L, E_{\nu}) \\ Z &= \Delta_{sol}^2 \times f_Z(\theta_{12}, \theta_{23}, A, L, E_{\nu}) \end{cases}$$

For the silver channel we have:

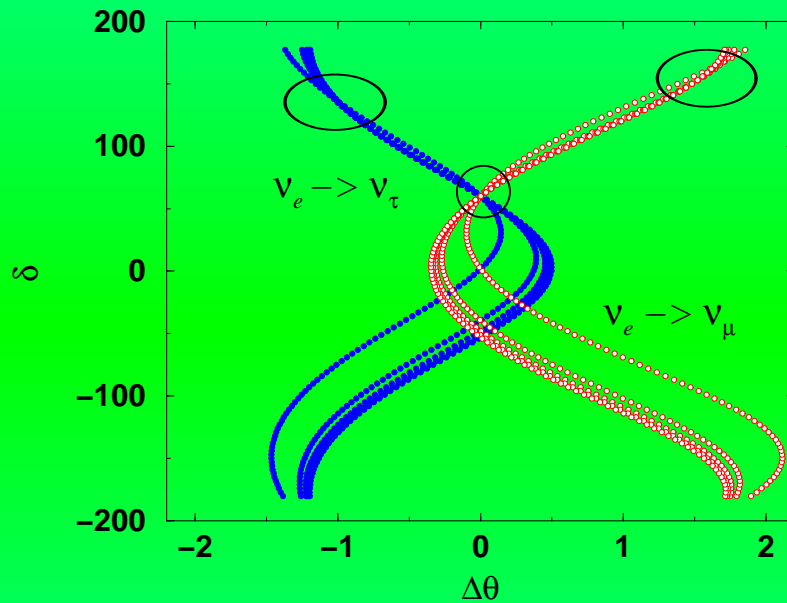
$$\begin{cases} X_{\pm}^{\tau} &= (c_{23}^2/s_{23}^2)X_{\pm} \\ Y_{\pm}^{\tau} &= Y_{\pm} \\ Z^{\tau} &= (s_{23}^2/c_{23}^2)Z \end{cases}$$

(+ neutrinos, - antineutrinos)

Notice: X, Z interchange $\theta_{23} \rightarrow \pi/2 - \theta_{23}$.

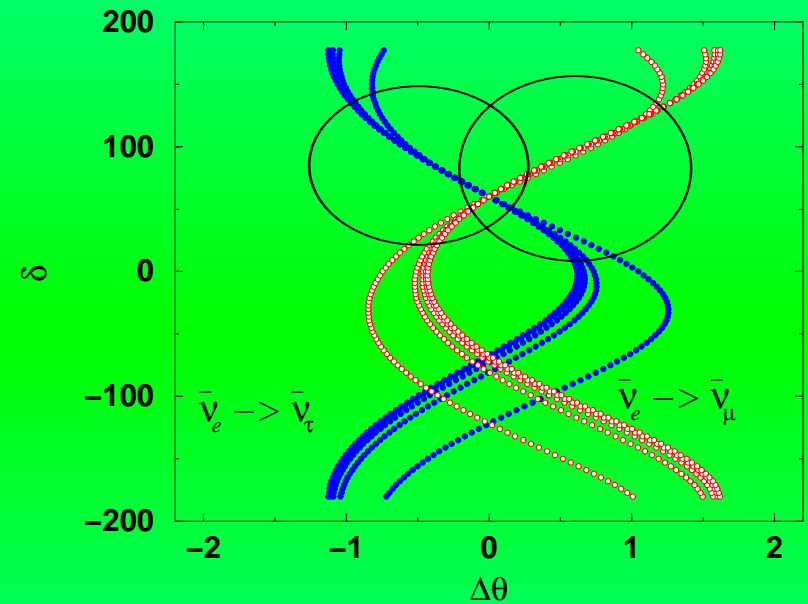
The intrinsic clones

$L = 732 \text{ Km}$



Neutrinos

$L = 732 \text{ Km}$



Antineutrinos

$$\begin{cases} \bar{\theta}_{13} &= 5^\circ \\ \bar{\delta} &= 90^\circ \end{cases}$$

$$\Delta\theta = \theta_{13} - \bar{\theta}_{13}$$

The ν -factory/detectors setup

CERN design for a 2.2 GeV superbeam
and a 50 GeV Neutrino Factory

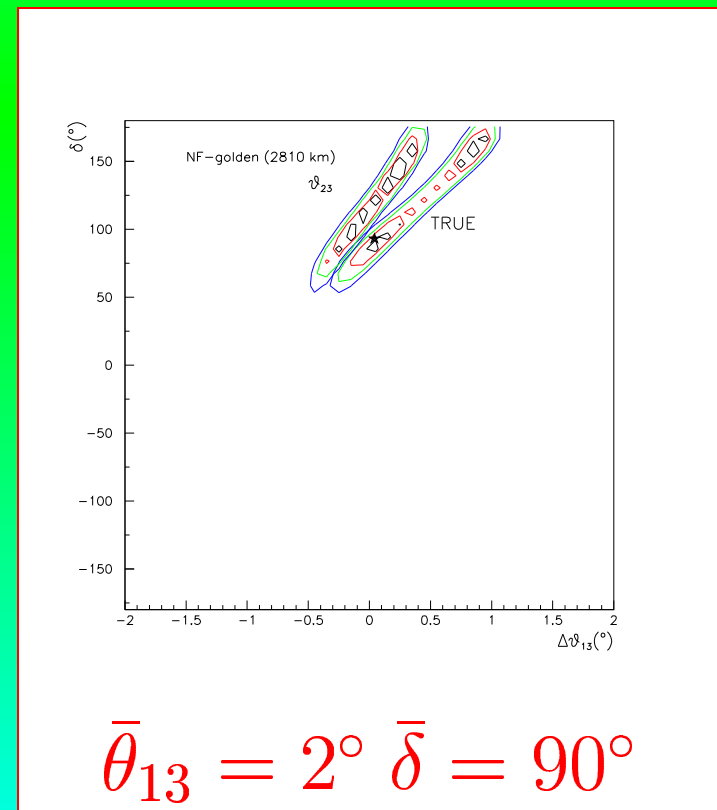
- 40 Kton Magnetized iron detector (**MID**)
 $L = 2810$ Km (Canary Islands)
A. Cervera *et al.*,
Nucl. Instr. Meth. A 451 (2000) 123; NuFact99, Lyon
- 400 Kton Water Cherenkov (**WC**)
 $L = 130$ Km (Frejus)
A. Blondel *et al.*,
Nucl. Instr. Meth. A 503 (2001) 173; NuFact01, Tsukuba
- 4 Kton Emulsion Cloud Chamber (**ECC**)
 $L = 732$ Km (Gran Sasso) or $L = 2810$ Km
D. Autiero *et al.*, hep-ph/0305185; NuFact03, New York

One detector

Consider the **NuFact golden channel**:
best option for one detector, with baseline $L = 2810$
(no sign degeneracies for $\theta_{13} \geq 1^\circ$).

A. Cervera *et al.*, hep-ph/0002108

- 40 Kton MID



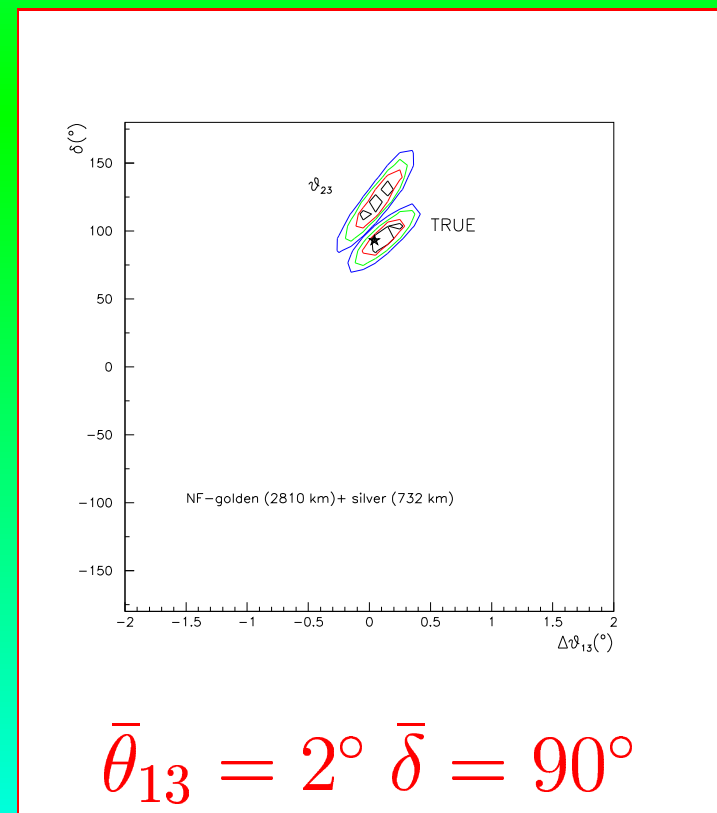
Two detectors

You can now add a second detector.

We can take advantage of the **NuFact silver channel...**

A. Donini *et al.*, hep-ph/0206034

- 40 Kton MID
- 4 Kton ECC

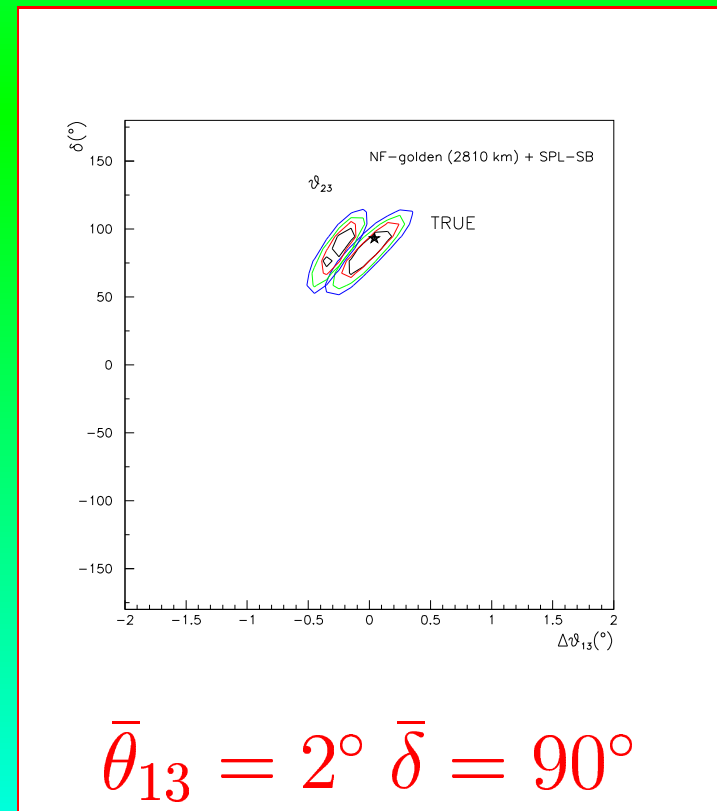


Two detectors

... or of the Superbeam-driven water Cherenkov.

J. Burguet-Castell *et al.*, hep-ph/0207080

- 40 Kton MID
- 400 Kton WC

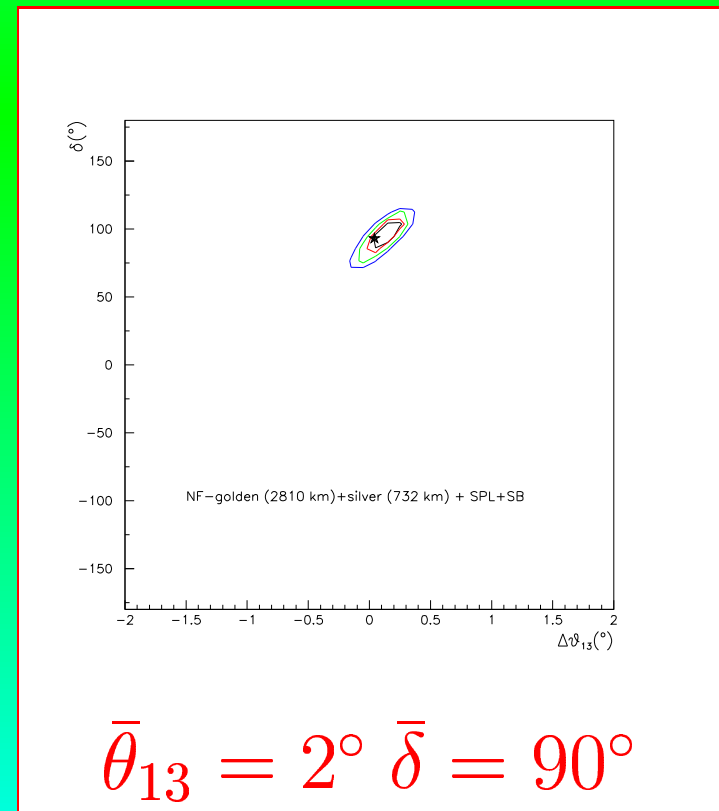


The Three Detectors

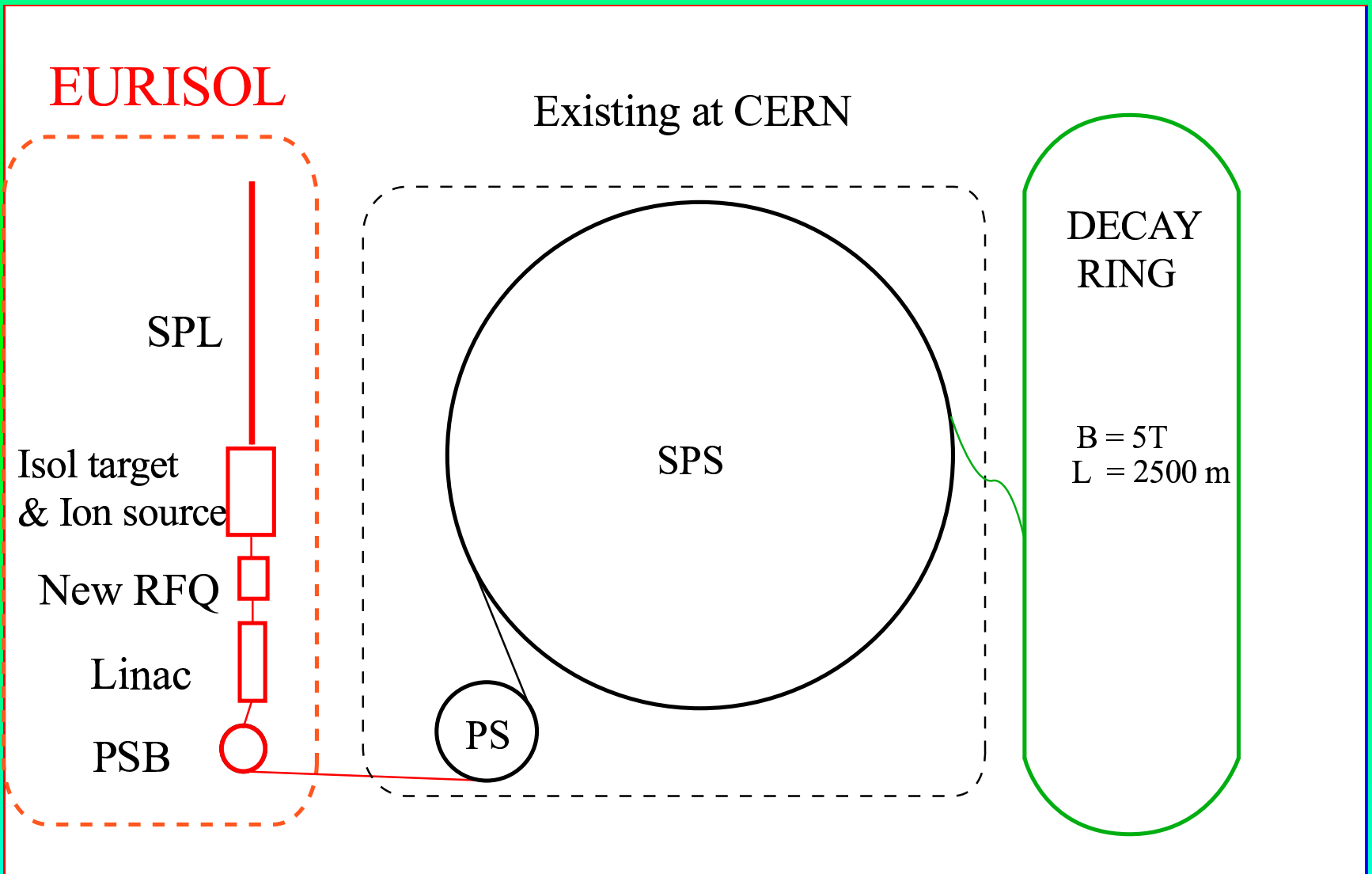
However, the very best possibility is to combine the three detectors in their **FULL GLORY**.

A. Donini, hep-ph/0310014

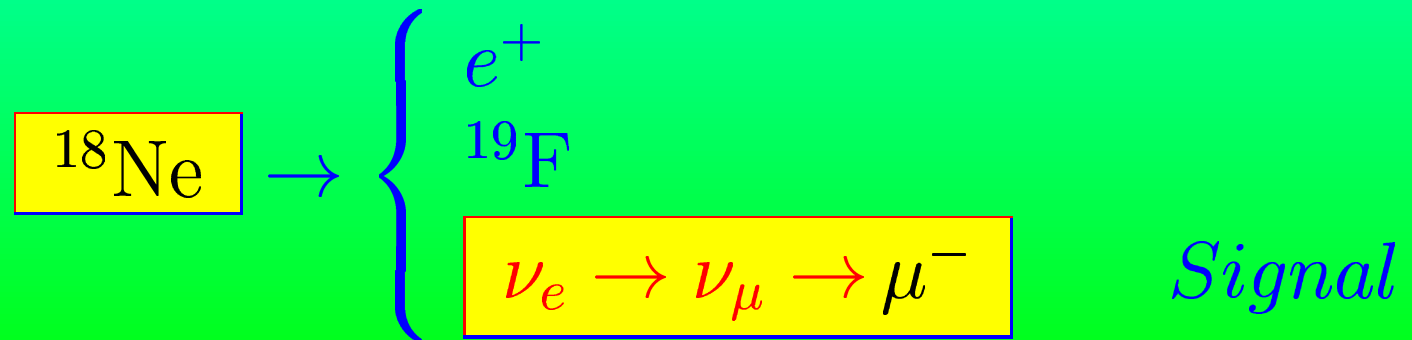
- 40 Kton MID
- 4 Kton ECC
- 400 Kton WC



The Beta-Beam at CERN



The Golden channel: β -beam



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$$+ Y_\pm \cos\left(\delta \mp \frac{\Delta_{atm}L}{2}\right) \cos\theta_{13} \sin(2\theta_{13})$$

$$+ Z + \dots$$

The β -beam/detector setup

CERN design for 2.2 GeV superbeam
and a low- γ β -beam:

$\gamma = 60$ for ${}^6\text{He}$; $\gamma = 100$ for ${}^{18}\text{Ne}$

- 440 Kton Water Cherenkov (WC)
 $L = 130$ Km (Frejus)

UNO Collaboration, hep-ex/0005046;

D. Casper, Nucl. Phys. Proc. Suppl. 112 (2002) 161.

The high- γ option must be considered seriously:

J. Burguet-Castell *et al.*, hep-ph/0312068

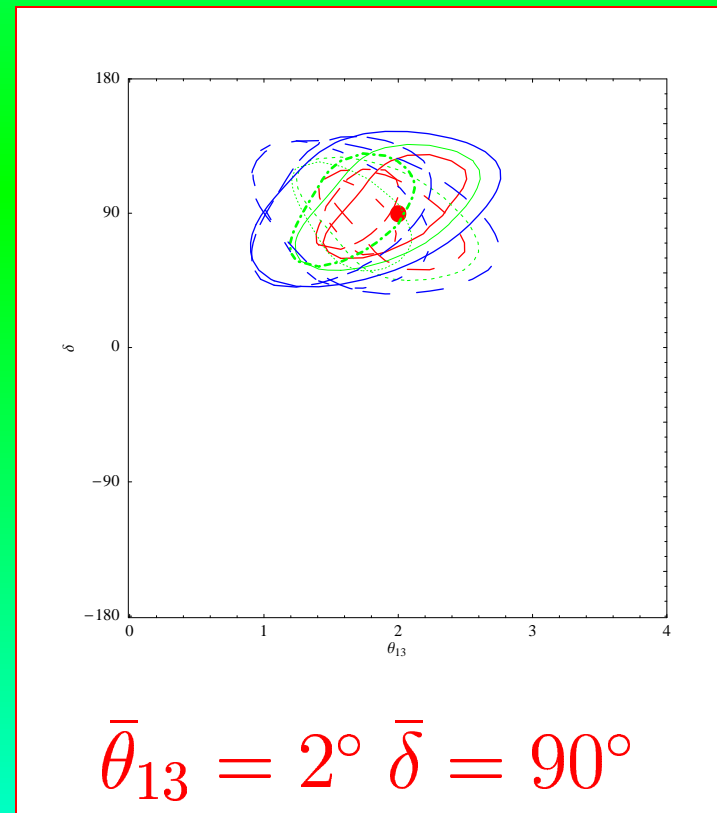
The SPL and low- γ Beta Beam

Consider the $\nu_e \rightarrow \nu_\mu$ at the BB:

one massive Water Cherenkov, with baseline $L = 130$

J. Bouchez *et al.*, hep-ph/0310059

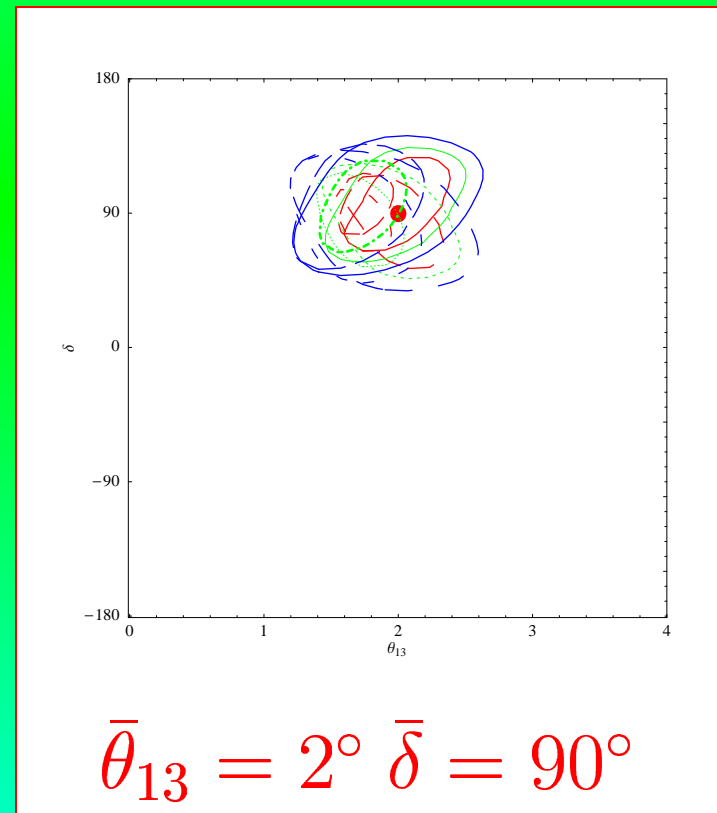
- 440 Kton WC-BB



The SPL and low- γ Beta Beam

You can now add $\nu_\mu \rightarrow \nu_e$ at the SPL:
same detector, same baseline

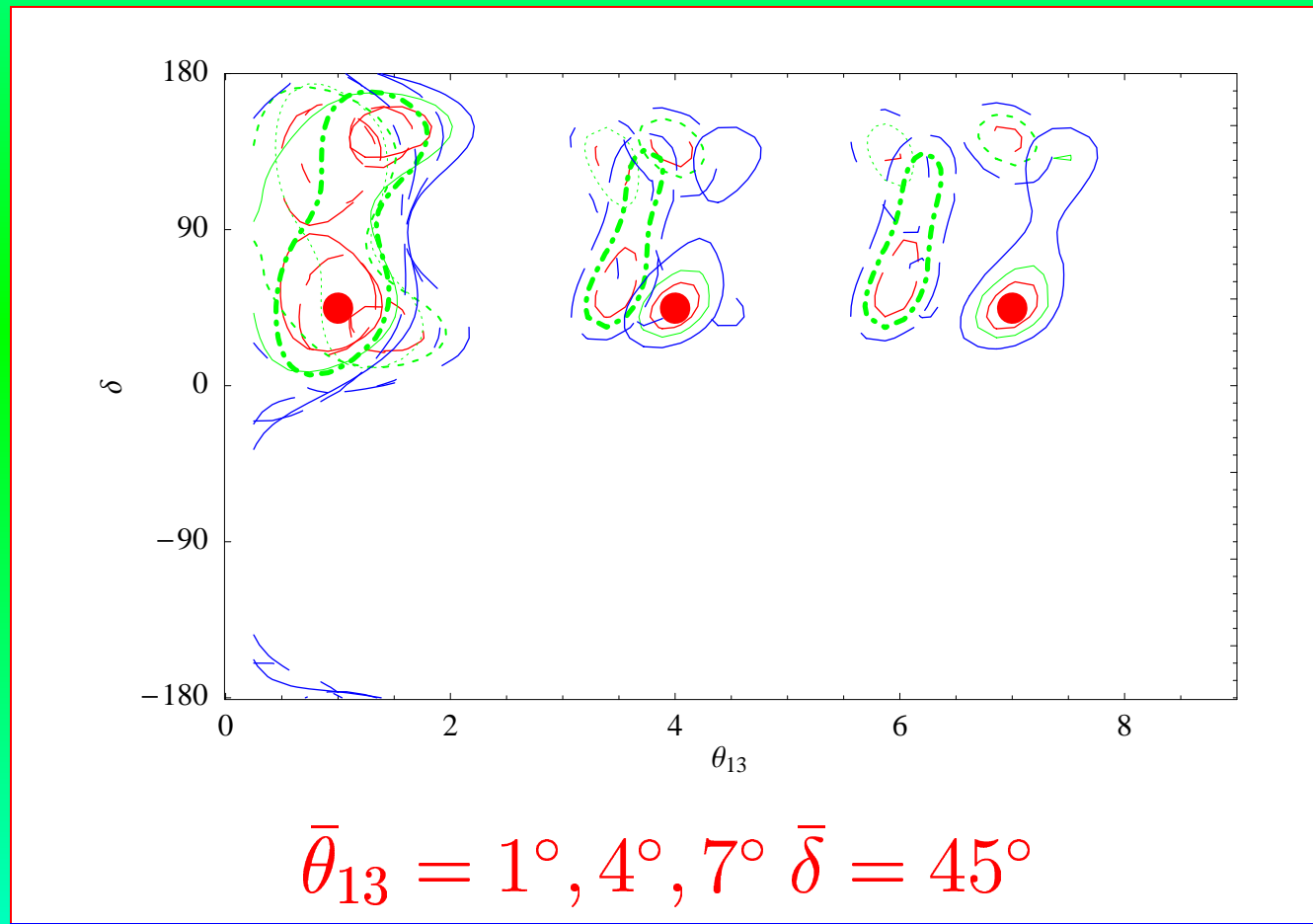
- 440 Kton WC-BB
- 440 Kton WC-SPL



The SPL and low- γ Beta Beam

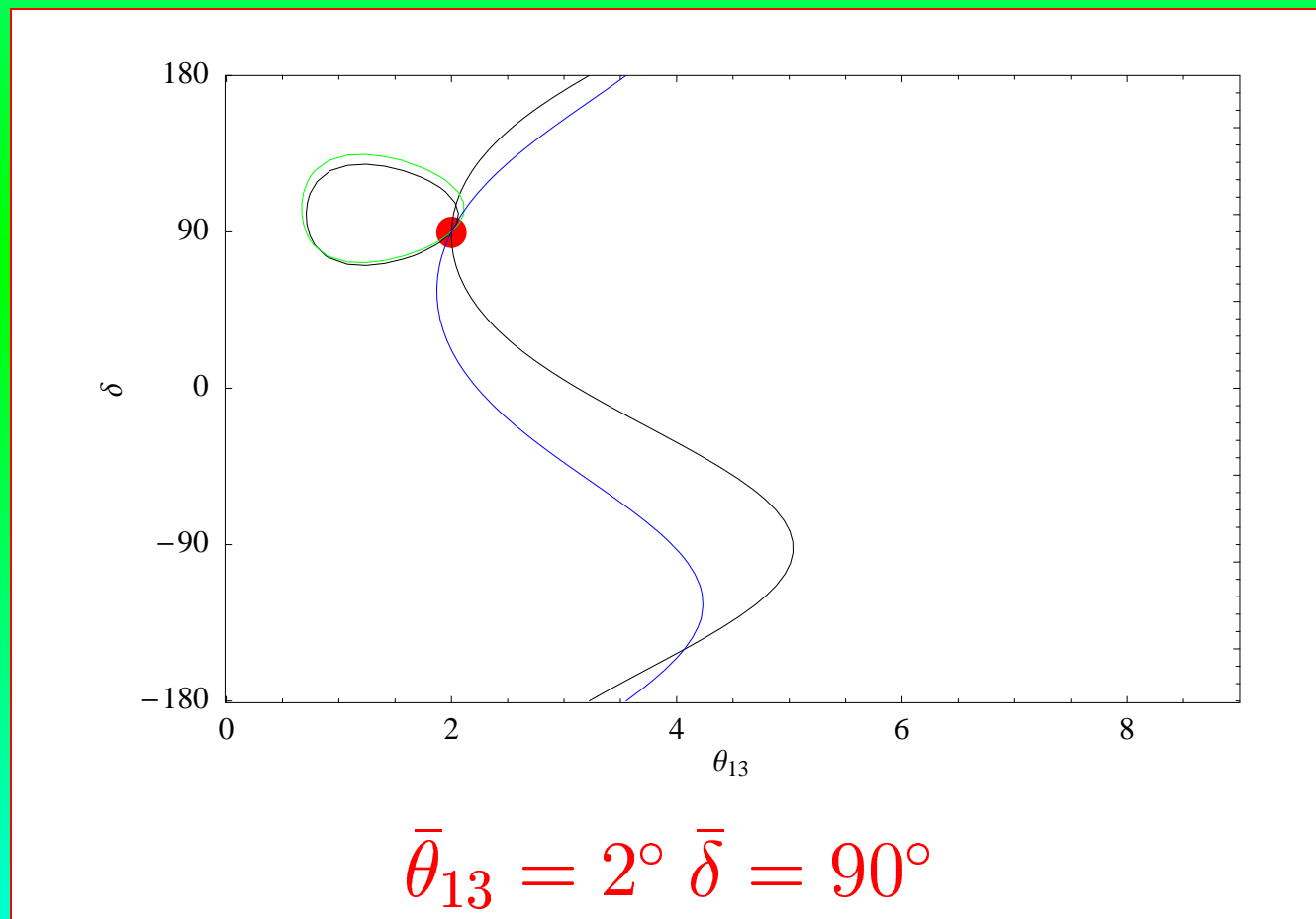
This is the general situation for $\delta \neq 90^\circ$.

A. Donini *et al.*, hep-ph/0406132



The SPL and low- γ Beta Beam

Unfortunately, there is NO SYNERGY:
same detector, same baseline, SAME ENERGY!



Conclusions

The next generation of neutrino experiments (beams, first generation superbeams and reactors) will improve our knowledge of the leptonic mixing matrix:

- ▷ mass differences $\Delta m_{23}^2, \Delta m_{12}^2$ at some %;
- ▷ mixing angles θ_{12}, θ_{23} at 10 %;
- ▷ the sign of Δm_{23}^2 , if lucky;
- ▷ the value of θ_{13} , if $\theta_{13} > 3^\circ$ ($\sin^2 2\theta_{13} > 0.01$).

Some (unexpected ?) presents could originate in natural source experiments, in particular from cosmology and supernovae.

Conclusions,2

The precision era in the leptonic mixing will start only with the next-to-next generation: the Neutrino Factory and/or the BetaBeam concepts.

In this phase, we will start
the quest for leptonic CP violation.

A non-vanishing phase δ is a key ingredient in leptogenesis (the most popular theory for baryogenesis, nowadays).

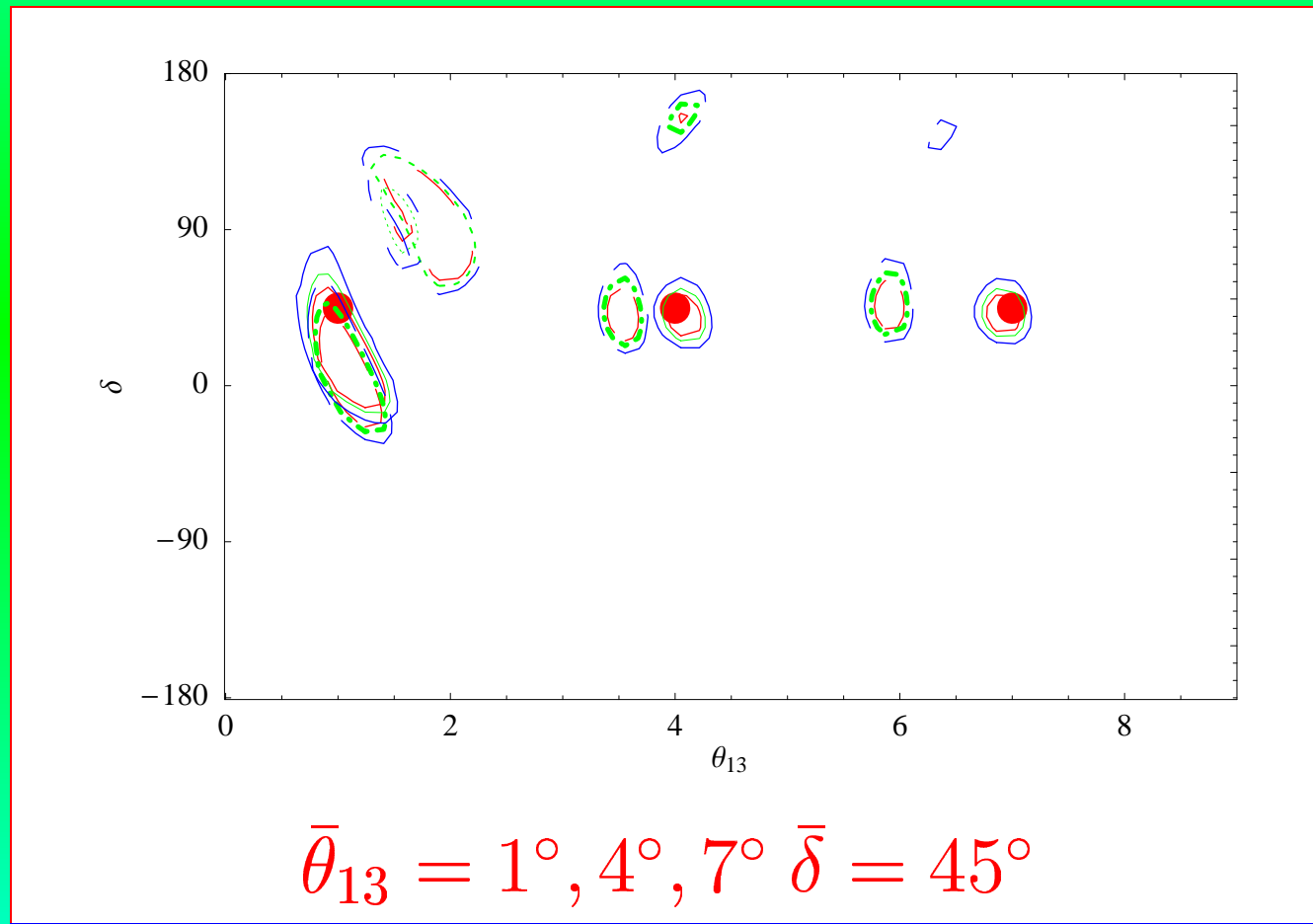
It is crucial to combine experiments and neutrino sources to solve the severe parameter degeneracy that obstacles a clean measurement of δ .

So what?

Outlook

First possibility: the High-gamma Beta Beam.

Half the statistics at $\gamma_{^{18}\text{Ne}} = 2500$, $\gamma_{^6\text{He}} = 1500$:



$L_{\text{golden}} = 2810 \text{ Km}$, $L_{\text{silver}} = 732 \text{ Km}$