Suggestion:

Review the first two chapters of Cowan – Statistical Data Analysis



Binomial distribution

Consider *N* independent experiments (Bernoulli trials): outcome of each is 'success' or 'failure', probability of success on any given trial is *p*.

Define discrete r.v. n = number of successes ($0 \le n \le N$).

Probability of a specific outcome (in order), e.g. 'ssfsf' is

$$pp(1-p)p(1-p) = p^n(1-p)^{N-n}$$

But order not important; there are

 $\frac{N!}{n!(N-n)!}$

ways (permutations) to get *n* successes in *N* trials, total probability for *n* is sum of probabilities for each permutation.

Binomial distribution (2)

The binomial distribution is therefore

$$f(n; N, p) = \frac{N!}{n!(N-n)!}p^n(1-p)^{N-n}$$
random parameters
variable

For the expectation value and variance we find:

$$E[n] = \sum_{n=0}^{N} nf(n; N, p) = Np$$
$$V[n] = E[n^{2}] - (E[n])^{2} = Np(1-p)$$

Binomial distribution (3)

Binomial distribution for several values of the parameters:



Example: observe *N* decays of W^{\pm} , the number *n* of which are $W \rightarrow \mu \nu$ is a binomial r.v., *p* = branching ratio.

Multinomial distribution

Like binomial but now *m* outcomes instead of two, probabilities are

$$\vec{p} = (p_1, \dots, p_m)$$
, with $\sum_{i=1}^m p_i = 1$.

For *N* trials we want the probability to obtain:

 n_1 of outcome 1, n_2 of outcome 2, \vdots n_m of outcome *m*.

This is the multinomial distribution for $\vec{n} = (n_1, \ldots, n_m)$

$$f(\vec{n}; N, \vec{p}) = \frac{N!}{n_1! n_2! \cdots n_m!} p_1^{n_1} p_2^{n_2} \cdots p_m^{n_m}$$

Multinomial distribution (2)

Now consider outcome *i* as 'success', all others as 'failure'.

 \rightarrow all n_i individually binomial with parameters N, p_i

$$E[n_i] = Np_i, \quad V[n_i] = Np_i(1-p_i) \quad \text{for all } i$$

One can also find the covariance to be

$$V_{ij} = Np_i(\delta_{ij} - p_j)$$

Example: $\vec{n} = (n_1, \dots, n_m)$ represents a histogram with *m* bins, *N* total entries, all entries independent.

Poisson distribution

Consider binomial n in the limit

$$N \to \infty, \qquad p \to 0, \qquad E[n] = Np \to \nu.$$

 \rightarrow *n* follows the Poisson distribution:

$$f(n;\nu) = \frac{\nu^n}{n!}e^{-\nu} \quad (n \ge 0)$$

$$E[n] = \nu, \quad V[n] = \nu.$$

Example: number of scattering events *n* with cross section σ found for a fixed integrated luminosity, with $\nu = \sigma \int L dt$.



10 n

0

0

5

7

3/13/20

15

20

From Binomial to Poisson to Gaussian $P(k:n,p) = \binom{n}{k} p^{k} (1-p)^{n-k}$

$$P(k:n,p) \xrightarrow{n \to \infty, np = \lambda} Poiss(k;\lambda) = \frac{\lambda^k e^{-k}}{k!}$$
$$\langle k \rangle = \lambda, \ \sigma_k = \sqrt{\lambda}$$

$$k \to \infty \Longrightarrow x = k$$

Using Stirling Formula

prob(x)=G(x,
$$\sigma = \sqrt{\lambda}$$
) = $\frac{1}{\sqrt{2\pi\sigma}}e^{-(x-\lambda)^2/2\sigma^2}$

This is a Gaussian, or Normal distribution with mean and variance of λ

Histograms



Uniform distribution

Consider a continuous r.v. x with $-\infty < x < \infty$. Uniform pdf is:



N.B. For any r.v. x with cumulative distribution F(x), y = F(x) is uniform in [0,1].

Example: for $\pi^0 \to \gamma \gamma$, E_{γ} is uniform in $[E_{\min}, E_{\max}]$, with $E_{\min} = \frac{1}{2} E_{\pi} (1 - \beta)$, $E_{\max} = \frac{1}{2} E_{\pi} (1 + \beta)$

Exponential distribution

The exponential pdf for the continuous r.v. x is defined by:



Example: proper decay time *t* of an unstable particle

 $f(t; \tau) = \frac{1}{\tau} e^{-t/\tau}$ (τ = mean lifetime)

Lack of memory (unique to exponential): $f(t - t_0 | t \ge t_0) = f(t)$

Gaussian distribution

The Gaussian (normal) pdf for a continuous r.v. *x* is defined by:

$$f(x;\mu,\sigma) = \frac{1}{\sqrt{2\pi}\sigma} e^{-(x-\mu)^2/2\sigma^2}$$

 $E[x] = \mu$ (N.B. often μ , σ^2 denote
mean, variance of any $V[x] = \sigma^2$ r.v., not only Gaussian.)



Special case: $\mu = 0$, $\sigma^2 = 1$ ('standard Gaussian'):

$$\varphi(x) = \frac{1}{\sqrt{2\pi}} e^{-x^2/2}, \quad \Phi(x) = \int_{-\infty}^x \varphi(x') dx'$$

If $y \sim \text{Gaussian with } \mu, \sigma^2$, then $x = (y - \mu) / \sigma$ follows $\varphi(x)$.

Gaussian pdf and the Central Limit Theorem

The Gaussian pdf is so useful because almost any random variable that is a sum of a large number of small contributions follows it. This follows from the Central Limit Theorem:

For *n* independent r.v.s x_i with finite variances σ_i^2 , otherwise arbitrary pdfs, consider the sum

$$y = \sum_{i=1}^{n} x_i$$

In the limit $n \to \infty$, y is a Gaussian r.v. with

$$E[y] = \sum_{i=1}^{n} \mu_i \qquad V[y] = \sum_{i=1}^{n} \sigma_i^2$$

Measurement errors are often the sum of many contributions, so frequently measured values can be treated as Gaussian r.v.s.

Meaning of parameter estimate



- We are interested in some physical unknown parameters
- Experiments provide samplings of some PDF which has among its parameters the physical unknowns we are interested in
- Experiment's results are statistically "related" to the unknown PDF
 - PDF parameters can be determined from the sample within some approximation or uncertainty
- Knowing a parameter within some error may mean different things:
 - **Frequentist**: a large fraction (68% or 95%, usually) of the experiments will contain, in the limit of large number of experiments, the (fixed) unknown true value within the quoted confidence interval, usually $[\mu \sigma, \mu + \sigma]$ ('coverage')
 - Bayesian: we determine a degree of belief that the unknown parameter is contained in a specified interval can be quantified as 68% or 95%
- We will see that there is still some more degree of arbitrariness in the definition of confidence intervals...

Statistical inference









Hypothesis tests





Parameter estimators



- An estimator is a function of a given sample whose statistical properties are known and related to some PDF parameters
 - "Best fit"
- Simplest example:
 - Assume we have a Gaussian PDF with a known $\,\sigma$ and an unknown μ
 - A single experiment will provide a measurement x
 - We estimate μ as $\mu^{est} = x$
 - The distribution of μ^{est} (repeating the experiment many times) is the original Gaussian
 - 68.27%, *on average*, of the experiments will provide an estimate within: $\mu \sigma < \mu^{est} < \mu + \sigma$
- We can determine: $\mu = \mu^{est} \pm \sigma$

Likelihood function

N



• Given a sample of *N* events each with variables $(x_1, ..., x_n)$, the likelihood function expresses the probability density of the sample, as a function of the unknown parameters:

$$L = \prod_{i=1}^{i} f(x_1^i, \cdots, x_n^i; \theta_1, \cdots, \theta_m)$$

 Sometimes the used notation for parameters is the same as for conditional probability:

$$f(x_1,\cdots,x_n|\theta_1,\cdots,\theta_m)$$

• If the size *N* of the sample is also a random variable, the extended likelihood function is also used:

$$L = p(N; \theta_1, \cdots, \theta_m) \prod_{i=1}^N f(x_1^i, \cdots, x_n^i; \theta_1, \cdots, \theta_m)$$

- Where *p* is most of the times a Poisson distribution whose average is a function of the unknown parameters
- In many cases it is convenient to use $-\ln L$ or $-2\ln L$: $\prod_{i} \rightarrow \sum_{i}$



Maximum likelihood estimates



- ML is the widest used parameter estimator
- The "best fit" parameters are the set that maximizes the likelihood function
 - "Very good" statistical properties
- The maximization can be performed analytically, for the simplest cases, and numerically for most of the cases
- Minuit is historically the most used minimization engine in High Energy Physics
 - F. James, 1970's; rewritten in C++ recently



CL&CImeasurement
$$\hat{\mu} = 1.1 \pm 0.3$$
 $L(\mu) = G(\mu; \hat{\mu}, \sigma_{\hat{\mu}})$ $\Rightarrow CI of \ \mu = [0.8, 1.4] at 68\% CL$

- A confidence interval (CI) is a particular kind of interval estimate of a population parameter.
- Instead of estimating the parameter by a single value, an interval likely to include the parameter is given.
- How likely the interval is to contain the parameter is determined by the confidence level
- Increasing the desired confidence level will widen the confidence interval.

Confidence Interval & Coverage

-Say you have a measurement μ_{meas} of μ with μ_{true} being the unknown true value of μ

-Assume you know the probability distribution function $\rho(\mu_{meas}|\mu)$

•based on your statistical method you deduce that there is a 95% Confidence interval $[\mu_1, \mu_2]$. (it is 95% likely that the μ_{true} is in the quoted interval)

The correct statement:

In an ensemble of experiments 95% of the obtained confidence intervals will contain the true value of μ .



Confidence intervals in practice

The recipe to find the interval [a, b] boils down to solving

$$\alpha = \int_{u_{\alpha}(\theta)}^{\infty} g(\hat{\theta}; \theta) \, d\hat{\theta} = \int_{\hat{\theta}_{obs}}^{\infty} g(\hat{\theta}; a) \, d\hat{\theta} \,,$$

$$\beta = \int_{-\infty}^{v_{\beta}(\theta)} g(\hat{\theta}; \theta) \, d\hat{\theta} = \int_{-\infty}^{\hat{\theta}_{obs}} g(\hat{\theta}; b) \, d\hat{\theta} \,.$$



 $\rightarrow a$ is hypothetical value of θ such that $P(\hat{\theta} > \hat{\theta}_{obs}) = \alpha$. $\rightarrow b$ is hypothetical value of θ such that $P(\hat{\theta} < \hat{\theta}_{obs}) = \beta$. Methods in Experimental Particle Physics

Meaning of a confidence interval

N.B. the interval is random, the true θ is an unknown constant. Often report interval [a, b] as $\hat{\theta}_{-c}^{+d}$, i.e. $c = \hat{\theta} - a, d = b - \hat{\theta}$. So what does $\hat{\theta} = 80.25^{+0.31}_{-0.25}$ mean? It does not mean: $P(80.00 < \theta < 80.56) = 1 - \alpha - \beta$, but rather: repeat the experiment many times with same sample size, construct interval according to same prescription each time, in $1 - \alpha - \beta$ of experiments, interval will cover θ .



Confidence Interval & Coverage

- •You claim, $Cl_{\mu}=[\mu_{1},\mu_{2}]$ at the 95% CL i.e. In an ensemble of experiments CL (95%) of the obtained confidence intervals will contain the true value of μ .
 - •If your statement is accurate, you have full coverage
 - olf the true CL is>95%, your interval has an over coverage
 - •If the true CL is <95%, your interval has an undercoverage



Neyman, J. (1937) <u>"Outline of a Theory of Statistical Estimation Based on the Classical Theory of</u> <u>Probability</u> Philosophical Transactions of the Royal Society of London A, 236, 333-380.

The Frequentist Game a 'la Neyman

Or

How to ensure a Coverage with Neyman construction



Fig. 7.1 Graphical illustration of Neyman belt construction (*left*) and inversion (*right*)

$$1 - \alpha = \int_{x^{\mathrm{lo}}(\theta_0)}^{x^{\mathrm{up}}(\theta_0)} f(x \mid \theta_0) \,\mathrm{d}x$$






















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By construction the probability to measure $x_0' < x_0$ if the true value $\mu = \mu_1(x_0)$ is $\alpha/2$ The determined C.I. is $[\mu_2(x_0), \mu_1(x_0)]$. $x_0' > x_0$ if the true value $\mu = \mu_2(x_0)$ is $\alpha/2$

Check the correct coverage: suppose μ^* is the true value. I repeat N times the measurement and determine each time the C.I. By construction in a fraction 1- α of the cases x_0 is within $x_1(\mu^*)$ and $x_2(\mu^*)$ and the corresponding C.I. provides coverage of μ^* .

$$P(x_1(\mu^*) < x_0 < x_2(\mu^*)) = 1 - \alpha$$



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Fig. 7.1 Graphical illustration of Neyman belt construction (*left*) and inversion (*right*)



Fig. 7.3 Neyman belt for the parameter μ of a Gaussian with $\sigma = 1$ at the 68.27% confidence level

Suppose Poisson variable and n=0 is measured (no background) Upper limit (lower limit =0) => 0 ± 0 (freq) or 1 ± 1 (Bayes)?

By construction the probability to measure $x_0' < x_0$ if the true value $\mu = \mu_1(x_0)$ is $(1-\alpha)$ (only one limit) or the probability to measure $x_0' > x_0$ if the true value $\mu = \mu_1(x_0)$ is α and λ

$$P(n > 0/\lambda) = \sum_{n=1}^{\infty} \frac{\lambda^{n} e^{-\lambda}}{n!} = 1 - e^{-\lambda} = \alpha \qquad \lambda \text{ frequentist} \qquad \lambda$$

Note:
in this example
 α has complementary
meaning than in
previous slides
 $\alpha => 1 - \alpha$

$$g(\lambda/n = 0) = \frac{p(n = 0/\lambda)f_0(\lambda)}{\int_{0}^{\infty} p(n \neq 0/\lambda)f_0(\lambda)d\lambda^{\lambda}} = \frac{e^{-\lambda}}{\int_{0}^{\infty} e^{-\lambda}d\lambda} = e^{-\lambda}$$
Bayesian
(uniform prior)
 λ

$$p(\lambda < \overline{\lambda}) = \int_{0}^{\overline{\lambda}} e^{-\lambda}d\lambda = 1 - e^{-\overline{\lambda}} = \alpha$$

$$\lambda$$

Methods in Experimental Particle Physics

(V==1/V)g

51

Λ

λ

α=0.95

1.5

2.5

3/13/20

frequentist limits

By construction the probability to measure $n_0' < n_0$ if the true value $s = s_{up}(n_0)$ is $(1-\beta)$ (only one limit) or the probability to measure $n_0' > n_0$ if the true value $s = s_{up}(n_0)$ is β



FIGURE 19. Neyman construction for the case of an upper limit. In this case a segment between $n_1(\theta)$ and ∞ is drawn for each value of the parameter θ . The segments define the confidence region. Once a value of n, n_0 is obtained, the upper limit s_{up} is found. (For simplicity the discrete variable n is considered as a real number here).



Fig. 9.9 Upper limits ν_s^{up} at a confidence level of $1 - \beta = 0.95$ for different numbers of events observed n_{obs} and as a function of the expected number of background events ν_b . (a) The classical limit. (b) The Bayesian limit based on a uniform prior density for ν_s .