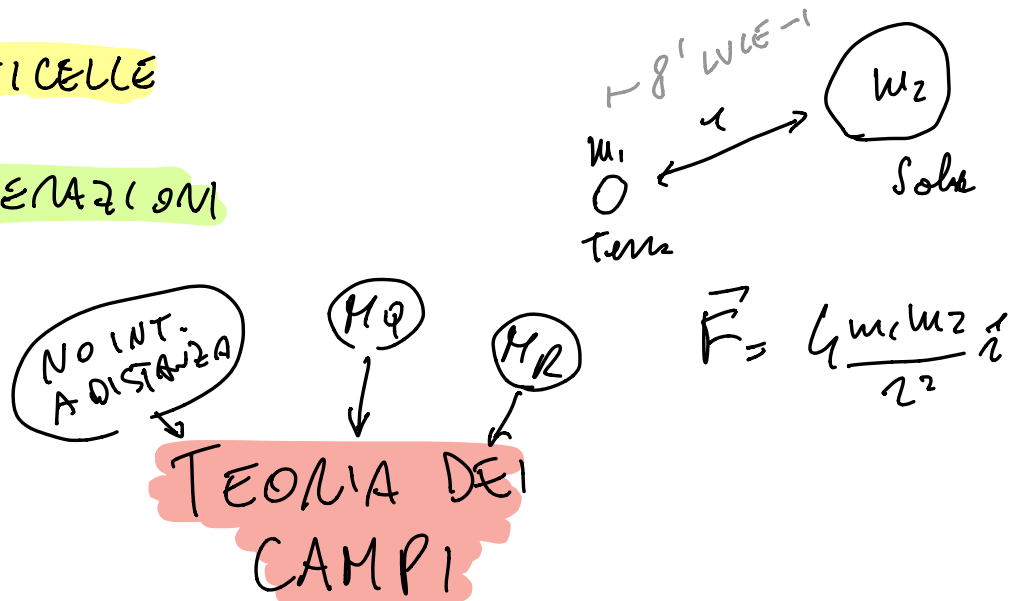


PREVIOUSLY ON FNSN 1

PARTICELLE

INTERAZIONI



- Scoperta elettrone
- Scoperta relatività

PARTICELLE

• sono localizzate nello spazio

$$p, E \dots \quad E^2 = p^2 + m^2 \quad (c=1)$$

• posso creare e s: annichilano

$$A + B \rightarrow C + D + E + F \dots$$

$$\textcircled{es} \quad p + p \rightarrow \pi^+ + \pi^0 + \pi^- + \dots$$

SCATTERING

$$\textcircled{es} \quad e^+ + e^- \rightarrow \mu^+ + \mu^-$$

- alcune di esse sono instabili:

$$A \rightarrow B + C + \dots$$

$$\textcircled{es} \quad n \rightarrow p + e + \bar{\nu}_e$$

CONTROESEMPLO: e^- ELEMENTARE
 p NON ELEMENTARE

$$p \rightarrow n + \text{altro}$$

NON PUÒ ACCADERE

$$\sqrt{s} = m_p = m_n + E$$

$$\text{MA } m_p < m_n$$

- hanno gradi di libertà

→ es: SPIN

- esistono antiparticelle

è la particella che ha
 lo stesso nome, stessa spin,
 ma \textcircled{es} carica elettrica
 opposta

(LS)	e^-	$m_{e^+} = m_{e^-}$
	e^+	spin sempre $1/2$
(LS)	γ	spin 1 massa 0 } E' LA SUA STESSA ANTIPARTIC.

TEMA DELLO SCATTERING DIP. DAL TEMPO

$$i\hbar \frac{\partial}{\partial t} |\psi\rangle = H |\psi\rangle$$

particelle libere

$$p = -i\hbar \nabla$$

$$H_0 = \frac{p^2}{2m} = -\frac{\hbar^2}{2m} \nabla^2$$

$$\psi_k = \psi_k(\vec{x}, t) = \underbrace{\psi_k(\vec{x})}_{\text{SPAZIO}} e^{-\frac{i}{\hbar} E_k t}_{\text{TEMPO}}$$

$$H_0 \psi_k(\vec{x}, t) = E_k \psi_k(\vec{x}, t)$$

$$\int_{\mathcal{V}} d^3x \psi_k^*(\vec{x}, t) \psi_l(\vec{x}, t) = \delta_{kl}$$

1 PARTICELLA PER VOLUME \mathcal{V}

caso generale

$$H = H_0 + V(\vec{x}, t)$$

$$H\psi = E\psi$$

$$\psi(\vec{x}, t) = \sum_k \underbrace{a_k(t)}_{\text{soluz. di } H_0} \psi_k(\vec{x}, t)$$

$$i\hbar \frac{\partial \psi}{\partial t} = i\hbar \sum_k \left[\frac{\partial a_k(t)}{\partial t} \psi_k(\vec{x}, t) + a_k(t) \frac{\partial \psi_k(\vec{x}, t)}{\partial t} \right]$$

$$= i\hbar \sum_k \left[\frac{\partial a_k(t)}{\partial t} \psi_k(\vec{x}) \cdot e^{-iE_k t} + a_k(t) \cdot \psi_k(\vec{x}) (-iE_k) e^{-iE_k t} \right]$$

$$= H\psi = (H_0 + V(\vec{x}, t))\psi$$

$$= \sum_k a_k(t) \cdot (E_k) \cdot \psi_k(\vec{r}) e^{-i E_k t}$$

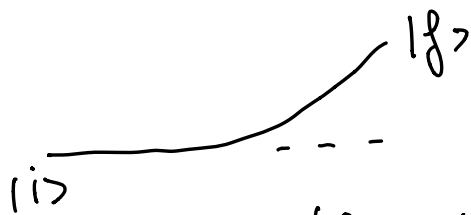
$$H_0 \psi = E_k \psi + \sum_k a_k(t) \cdot V(\vec{r}, t) \cdot \psi_k(x)$$

È SEMPRE $H\psi = i \frac{\partial \psi}{\partial t} e^{-i E_k t}$

$$i \sum_k \frac{\partial a_k(t)}{\partial t} \psi_k(\vec{r}) e^{-i E_k t}$$

$$= \sum_k a_k(t) \cdot V(\vec{r}, t) \psi_k(x) e^{-i E_k t}$$

tipicamente



$(|i\rangle, |f\rangle)$ LIBERA

$$|f\rangle = \psi_f(\vec{r}) e^{-i E_f t}$$

assumo $V(\vec{r}, t) = V(\vec{r})$

proiettiamo su $\langle f|$

per avere $\frac{\partial a_f}{\partial t} = \text{bla bla bla}$

$$i \sum_k \psi_f^*(\vec{r}) e^{i E_f t} \frac{\partial a_k(t)}{\partial t} \psi_k(\vec{r}) e^{-i E_k t}$$

$$= \sum_k \psi_f^*(\vec{r}) \cdot e^{i E_f t} \cdot a_k(t) \cdot V(\vec{r}, t) \psi_k(\vec{r}) e^{-i E_k t}$$

$$\int d^3x \dots = i \cdot \frac{\partial a_f(t)}{\partial t} = \sum_k \int d^3x \psi_f^*(\vec{r}) \cdot V(\vec{r}, t) \psi_k(\vec{r}) \times a_k(t) \cdot e^{i(E_f - E_k)t}$$

δ_{kl} impone $E_f = E_k$

$$= \sum_k a_k(t) \cdot \langle f | V | k \rangle e^{-i(E_k - E_f)t}$$

"T_{fk}"

IPOTESI $|i\rangle$ $|f\rangle$ iniz. stati

$$a_i(t) \equiv 1 \quad \forall t$$

$$a_{k \neq i}(t) \equiv 0$$

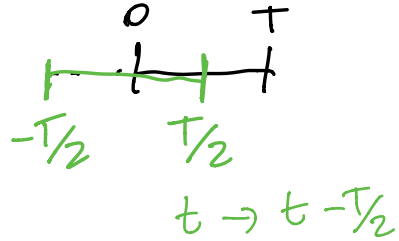
$$i \frac{\partial a_f(t)}{\partial t} = \langle f | V | i \rangle e^{-i(E_i - E_f)t}$$

$$\frac{\partial a_f(t)}{\partial t} = -i T_{fi} e^{-i(E_i - E_f)t}$$

$$a_f(t) = -i T_{fi} \int dt e^{-i(E_i - E_f)t}$$

PROB. DI TRANS.

$$P_{fi} = P_{fi}(T) = |a_f(T)|^2$$



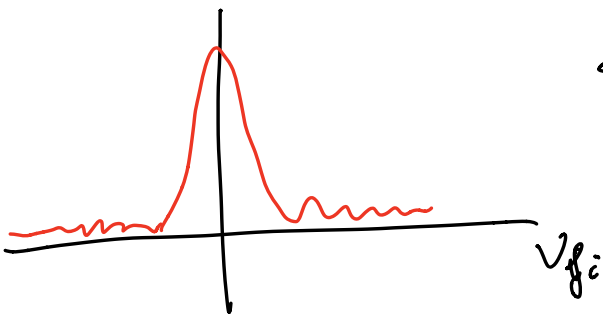
$$= |T_{fi}|^2 \int_{-T/2}^{T/2} dt \int_{-T/2}^{T/2} dt' e^{-i(E_i - E_f)t} e^{i(E_i - E_f)t'}$$

VOLGLO IL RATE DI TRANSIZIONE

$$d\Gamma_{fi} = \frac{P_{fi}}{T} = \frac{1}{T} |T_{fi}|^2 \int_{-T/2}^{T/2} \frac{2 \sin(\nu_{fi} t/2)}{\nu_{fi}} \left(\lim_{T \rightarrow \infty} \int_{-T/2}^{T/2} e^{i(E_i - E_f)t'} dt' \right)$$

$\nu_{fi} = \frac{E_f - E_i}{\hbar}$

$$= \frac{1}{T} |T_{fi}|^2 \cdot \frac{4 \sin^2(\nu_{fi} T/2)}{\nu_{fi}^2}$$



se c'è un continuo di stati con energia $[E_f, E_f + dE_f]$

$$\lim_{t \rightarrow \infty} \frac{1}{T} \cdot |T_{fi}|^2 \cdot \int_{-T/2}^{T/2} dt \int_{-T/2}^{T/2} dt' e^{i(E_f - E_i)t} \times e^{-i(E_f - E_i)t'}$$

$\times d\mu$

QUANTI STATI ACCESSIBILI
 CISONO FRA E_f ed $E_f + dE_f$

APPLICO

$$\lim_{T \rightarrow \infty} \int_{-T/2}^{T/2} e^{ik(x-x_0)} dk = 2\pi \delta(x-x_0)$$

E OTTENGO

$$= dT_{fi} = |T_{fi}|^2 \cdot \lim_{T \rightarrow \infty} \frac{1}{T} \int_{-T/2}^{T/2} dt' \frac{d\mu}{dE_f} \cdot dE_f \cdot 2\pi$$

$$= |T_{fi}|^2 \cdot 2\pi \cdot \left. \frac{d\mu}{dE_f} \right|_{E_f = E_i}$$

$$T_{fi}^2 = 2\pi \cdot |T_{fi}|^2 \cdot \left. \frac{d\mu}{dE_f} \right|_{E_f = E_i}$$

$$= \frac{2\pi}{\hbar} \cdot |T_{fi}|^2 \rho(E_i)$$

REGOLA DI
 ORA DI
 FERMI

ORDINE SUCCESSIVO

assumo $a_i(t) \cong 1$

$a_{k \neq i}(t) \cong 0$ (per $\epsilon_k \cong 0$)

$$i \frac{\partial a_j(t)}{\partial t} = \sum_k a_k(t) \langle j | V | k \rangle e^{-i(\epsilon_k - \epsilon_j)t}$$

$$= \langle j | V | i \rangle e^{-i(\epsilon_i - \epsilon_j)t}$$

$$+ \sum_{k \neq i} a_k(t) \langle j | V | k \rangle e^{-i(\epsilon_k - \epsilon_j)t}$$

ipotizzo

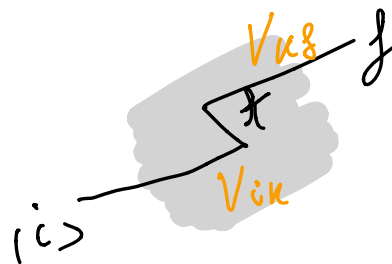
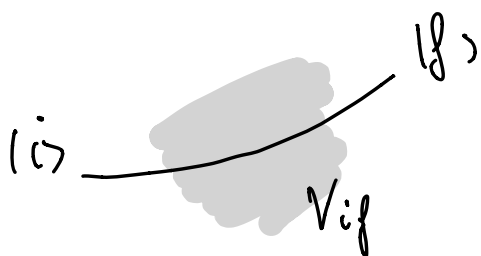
$$a_k(T) \equiv -i T \kappa_i \int_0^T e^{i(\epsilon_k - \epsilon_i)t} dt$$

$$= -i \langle k | V | i \rangle \frac{e^{i(\epsilon_k - \epsilon_i)T} - 1}{i(\epsilon_k - \epsilon_i)}$$

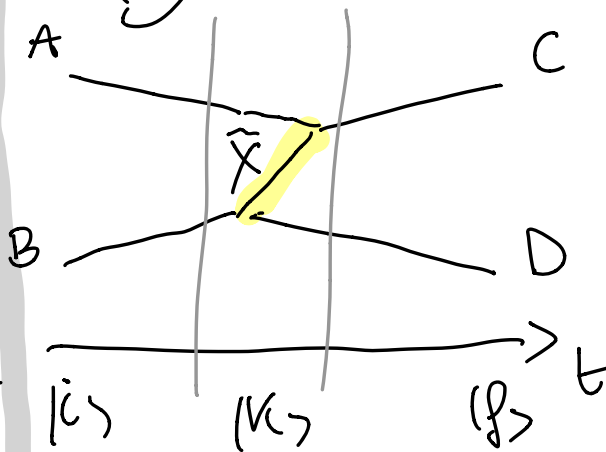
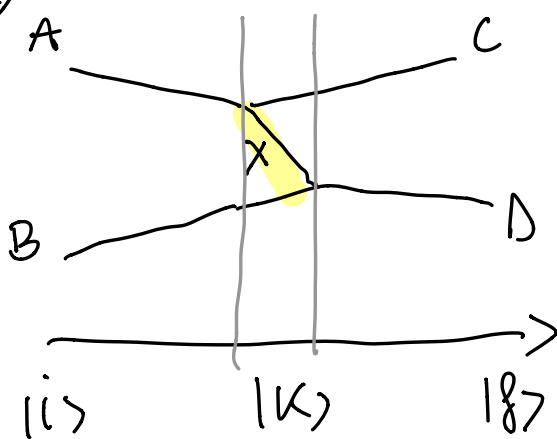
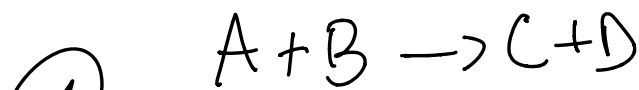
$$\frac{\partial a_j(t)}{\partial t} = -i \left[\langle j | V | i \rangle + \sum_{k \neq i} \frac{\langle j | V | k \rangle \langle k | V | i \rangle}{\epsilon_k - \epsilon_i} \right]$$

$$\times e^{-i(\epsilon_i - \epsilon_j)t}$$

$$= -i \int_{f_i}^{(al 2^o ordina)} e^{-i(E_i - E_f)t}$$



CAPIAMO!



$$|i\rangle : A + B$$

$$|k\rangle : (X + C) + B$$

$$|f\rangle : C + D$$

$$|k\rangle = (X + D) + A$$

X e \tilde{X} sono le puppe
autro.

$$\text{con } w_x = w_x^z$$

altra assunzione: scambio di
capitale

$$\textcircled{1} \vec{P}_A = \vec{P}_X + \vec{P}_C$$