

$$i \hbar \frac{\partial}{\partial t} |\psi\rangle = H |\psi\rangle$$

↳ equazione

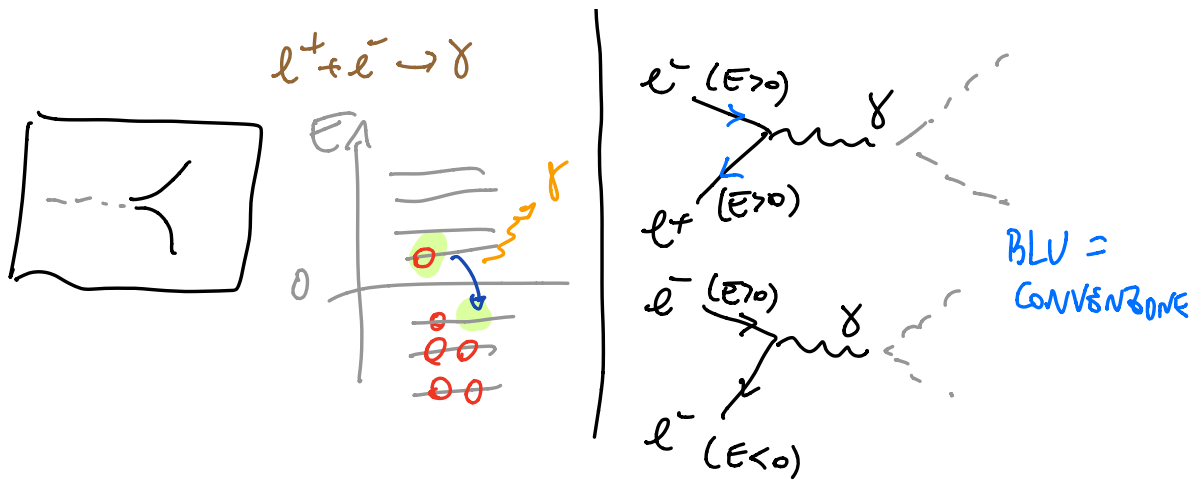
ANTI PARTICELLE

tentativo di conciliare MQ & RELSPEC

- $k < 0$: $E < 0$ vanno tenute
 $p < 0$ 2° ordine in $\partial/\partial x_{\mu}$
- Dirac = $E < 0$
 (spin 1/2) $p > 0$ 1° ordine in $\partial/\partial x_{\mu}$

DIRAC SEA

FERMIONI	FEYNMAN - STÜCKELBERG
<div style="display: flex; align-items: flex-start;"> <div style="margin-right: 20px;"> </div> <div> <p>◦ i livelli con $E < 0$ sono occupati</p> <p>$\gamma \rightarrow e^+ + e^-$ ($\gamma + N \rightarrow N + e^+ + e^-$)</p> </div> </div>	<p>$-iEt$ $-i(-E)(-t)$</p> <p>$e = e$</p> <p>- gli elettroni ad $E < 0$ sono soliti: associati a positroni che viaggiano indietro nel tempo</p>



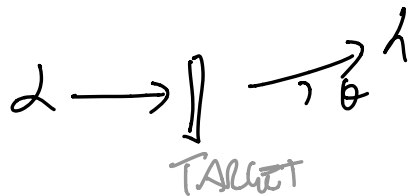
VENSO LE INTERAZIONI FORTI

Coulombiano \rightarrow FORTI \rightarrow DEBOLI

\rightarrow ELETTRODEBOLI

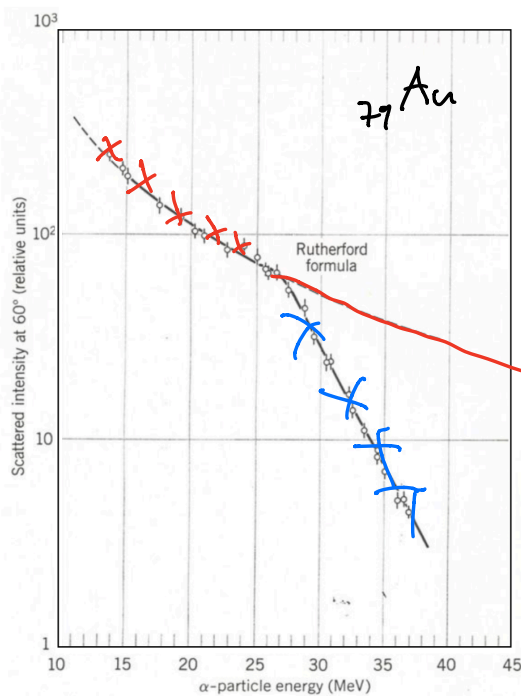
\rightarrow GRAVITAZ.

- Rutherford (1911)



$$A = \frac{e^2}{4\pi\epsilon_0} \cdot \frac{1}{r}$$

$$b = A \frac{1}{\tan^2(\theta/2)}$$



↑
28 MeV

$\theta = 60^\circ$
 $T = 28 \text{ MeV}$

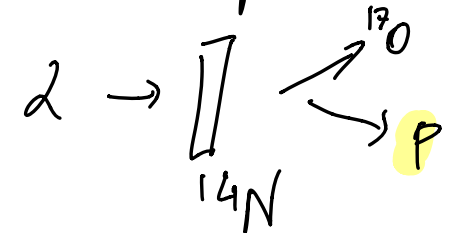
$$b = \frac{zZe^2}{4\pi\epsilon_0 T} \frac{1}{\tan \frac{\theta}{2}}$$

RUTHERFORD

$b \approx 14 \text{ fm}$

DAL BLU, "SENTIAMO"
UNA WICIAZ. A
CORTO RAGGIO

- 1917 : scoperta il positron



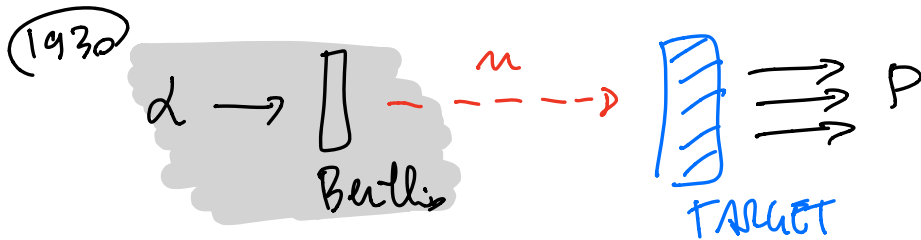
PROBLEMA:

$m_{nuclei} > Z m_p$

→ c'è altro

se fossero elettrici, $T = E - m_e \approx P$

se $p \approx \Delta p \geq \frac{\hbar}{\Delta x} = \frac{\hbar}{1 \text{ fm}} \approx 200 \text{ MeV}$



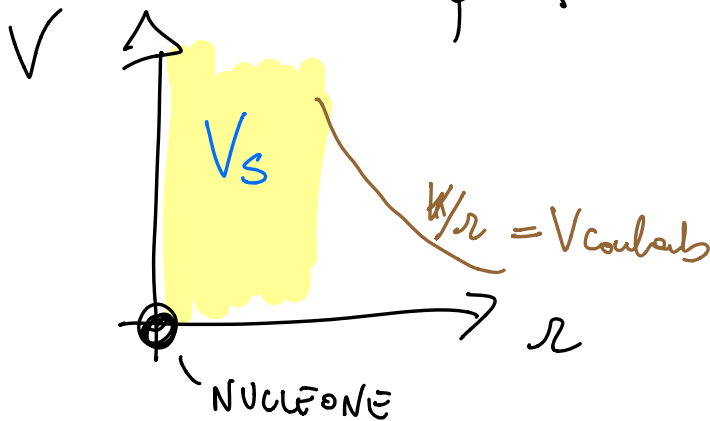
- H (proton) $T_p \approx 5 \text{ MeV}$
- ${}^{14}\text{N}$ $T_p \approx 1 \text{ MeV}$

$T_{\text{MAX}}^{(p)} = ?$



→ give $M_n \approx M_p$

NUCLEONI : p & n



$V_s(\text{strong})$
 here energy
 a cuts
 range

CAMPO EM

$$-\nabla^2 \phi = \rho/\epsilon_0 = \frac{e}{\epsilon_0} \delta^{(3)}(\vec{r})$$

$$\phi = \frac{e}{r} \frac{1}{4\pi\epsilon_0}$$

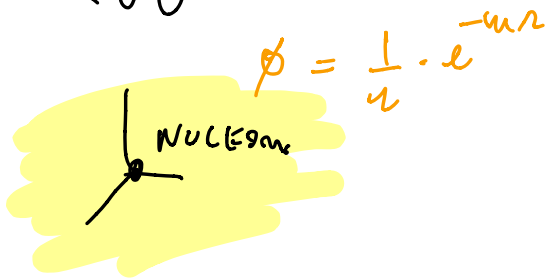
$$\partial_\mu \partial^\mu \phi = 0$$

$$\hookrightarrow \partial/\partial t \phi = 0$$

CAMPO
SCALARE

INT. FONTE

$$\left(\frac{\partial^2}{\partial t^2} - \nabla^2 + m^2 \right) \phi = 0 \quad (K-6)$$



ϕ potenziale di Yukawa

M massa del mediatore

$$(\nabla^2 - m^2) \phi = -g \delta^{(3)}(\vec{r})$$

PIONE: π^0
 π^+
 π^-

$$\frac{\partial \phi}{\partial \theta} = \frac{\partial \phi}{\partial \varphi} = 0$$

$$\phi = \phi(r)$$

$$\nabla^2 = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) + \underline{\underline{\text{celtre}}}$$

devo risolvere

$$(\nabla^2 - m^2) \phi = 0$$

$$\phi \equiv \frac{C}{4\pi} \frac{e^{-mr}}{r}$$

POTENZIALE
DI
YUKAWA

ADDAONE
(particelle
di
interazione
forte)

(M)

ADDAONE

$|i\rangle, |f\rangle$ ondeggiare

$$E_i = E_f$$

$$\vec{P}_i - \vec{P}_f \equiv \vec{P}$$

$$i \vec{P}_i \cdot \vec{x} - i E_i t$$

$$\psi_i = N e^{i \vec{P}_i \cdot \vec{x} - i E_i t}$$

$$a = \frac{\Gamma_{if}}{N} \quad (V=1)$$

$$\Gamma_{if} = 2\pi |T_{if}|^2 \rho(E_i)$$

$$T_{if} = \langle \psi_f | \phi | \psi_i \rangle$$

$$= \frac{C}{4\pi} |N|^2 \int_0^{2\pi} d\varphi \int_0^\pi \sin\theta d\theta \int_0^\infty dr r^2$$

$$\times \frac{1}{r} e^{-mr} \cdot e^{i\vec{p}\vec{r}}$$

$$\vec{p}\vec{r} = pr \cos\theta$$

$$y \equiv \cos\theta$$

$$dy = -\sin\theta d\theta$$

(|N|=1)

$$= \frac{C}{4\pi} \cdot 2\pi \int_{-1}^1 dy \int_0^\infty dr r e^{-mr} e^{ipy}$$

$$= \frac{C}{2} \int_0^\infty dr \cdot r e^{-mr} \left[\frac{e^{ipy} - e^{-ipy}}{ipy} \right]$$

$$= \frac{C}{2ip} \left[-\frac{1}{ip-m} - \frac{1}{ip+m} \right]$$

$$= \frac{C}{2ip} \left[\frac{-ip-m - ip+m}{-p^2 - m^2} \right]$$

$$= \frac{c}{p^2 + m^2}$$

$$|T_{fi}|^2 = \frac{c^2}{(p^2 + m^2)^2}$$

$$p(E_i) = \left. \frac{d\omega}{dE_f} \right|_{E_f = E_i} = \left. \frac{d\omega}{dp_f} \left| \frac{dp_f}{dE_f} \right| \right|_{E_f = E_i}$$

$$= \left. \frac{p^2 dp_f d\Omega}{(2\pi)^3 dp_f} \right| \left. \frac{dp_f}{dE_f} \right|_{E_f = E_i}$$

$$\frac{dp_f}{dE_f} = \frac{\sqrt{2M}}{\sqrt{4E_f}} = \sqrt{\frac{M}{2E_f}}$$

$$E_f = p_f^2 / 2M$$

$$p_f = \sqrt{2ME_f}$$

$$p(E_i) = \sqrt{\frac{M}{2E_i}} \cdot \frac{d\Omega}{(2\pi)^3} p_f^2$$

$$= M \sqrt{2ME_j} \frac{d\Omega}{(2\pi)^3}$$

$$\Gamma_{fi} = 2\pi |T_{fi}|^2 \rho(E_i)$$

$$\phi = \nu = \frac{p_f}{M} = \frac{\sqrt{2ME_j}}{M}$$

$$d\sigma = 2\pi \cdot |T_{fi}|^2 \cdot M \sqrt{2ME_j} \frac{d\Omega}{(2\pi)^3}$$

$$\frac{M}{\sqrt{2ME_j}}$$

$$\sigma = \int d\Omega "d\sigma" =$$

$$= \frac{c^2}{\pi} \left(\frac{M}{(p^2 + m^2)^2} \right)^2$$

se $p \ll m$ basro

$$\sigma \approx \frac{c^2}{\pi} \left(\frac{M}{m^2} \right)^2$$

$$\equiv 4\pi a^2$$

- vogliamo stime
m del perimetro di
Yukawa (PION)

- vogliamo stime
c

$$4\pi a^2 = \frac{c^2}{\hbar} \left(\frac{M}{m^2} \right)^2$$

$$\rightarrow a = \frac{1}{2\pi} c \left(\frac{M}{m^2} \right)$$

Stime m

$$e^{-mr} \equiv e^{-\frac{r}{\lambda}}$$

$$\lambda \approx 1 \text{ fm} \rightarrow M \approx 200 \text{ MeV}$$

$$\frac{1}{1 \text{ fm}} \Rightarrow \frac{197 \text{ MeV} \cdot \text{fm}}{1 \text{ fm}}$$

$$\bullet a \approx \lambda \sim 1 \text{ fm}$$

$$M \approx 1 \text{ GeV} \quad (\sim M_p, M_n)$$

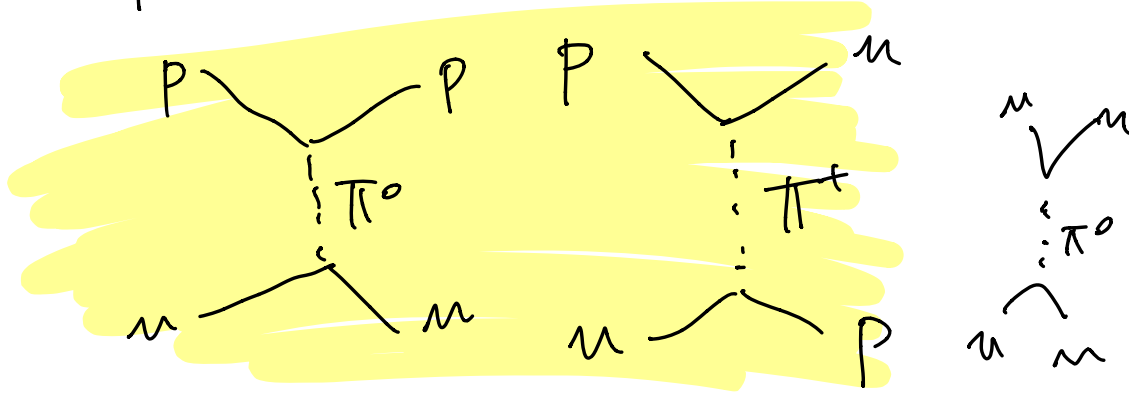
$$C \approx O(1)$$

per caso EM:

$$V = \frac{e^2}{4\pi\epsilon_0 r}$$

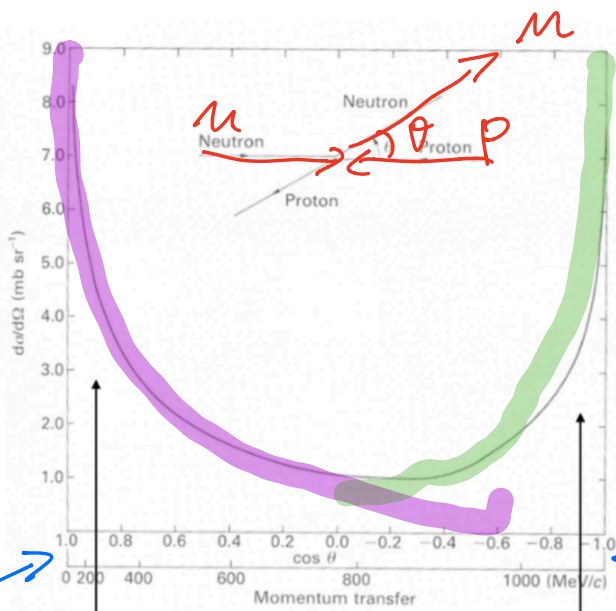
$$C_{EM} = \frac{e^2}{\epsilon_0} = 4\pi\alpha \approx 0.1$$

$$C \gg \underbrace{C_{EM}}_{(\text{è Forte})}$$



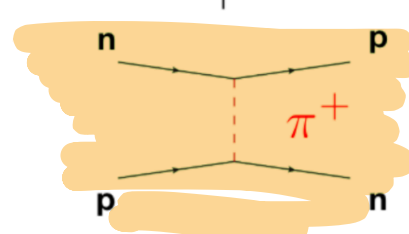
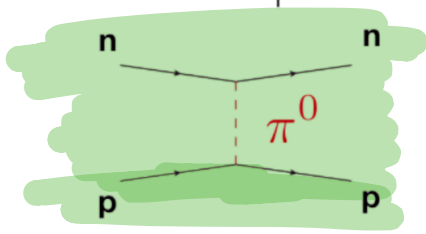
$\rightarrow t$

$\sigma \uparrow$



$\theta = 0$

$\theta = \pi$



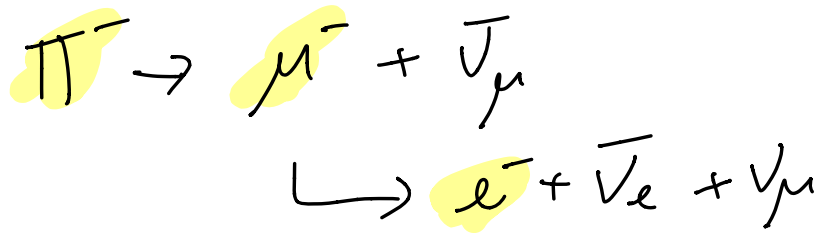
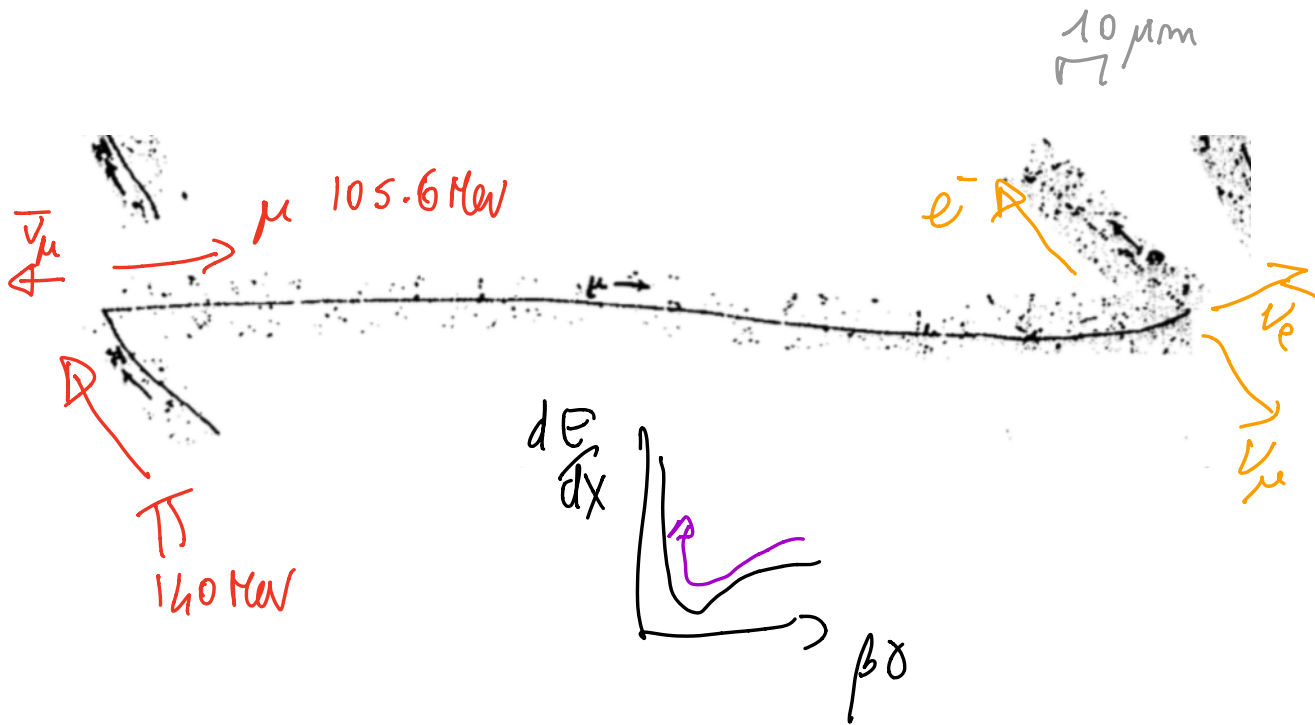
Size reports of PION

0 mgl, π^+ π^-
 π^0

$$M_{\pi^\pm} \approx 140 \text{ MeV}$$

$$M_{\pi^0} \approx 135 \text{ MeV}$$





\hookrightarrow π^- (π^+ , π^0) particelle di Yukawa

MESONI

$$\mu^- = \dots$$

- staccante: MESONE $m \in [m_e, m_p]$
- MESONE stato legato q \bar{q} (FERMIONI)

o BACON 9 9 9

o 4, \$ quark