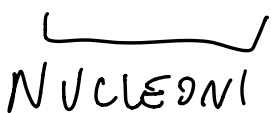


- CINEMATICA / DECADIMENTI
  - SEZ. D'URTO
  - PERDITE DI ENERGIA
- 

## IL NUCLEO

- interazione a corto raggio
  - $Z$  protoni,  $A-Z$  neutroni
- 
  
 NUCLEONI

$$A = 50$$

50! istanze

→ Sono quantità ollettive

- statistiche  $R, E_L, M, \dots$

- dinamiche

decadimenti  
 $\alpha, \beta, \gamma$

NUCLIDE  $(A, Z)$

ISOTOPI

"nuclei di cui lo stesso  $Z$ "

$^{12}\text{C}$ ,  $^{14}\text{C}$

ISOBARI

lo stesso  $A$

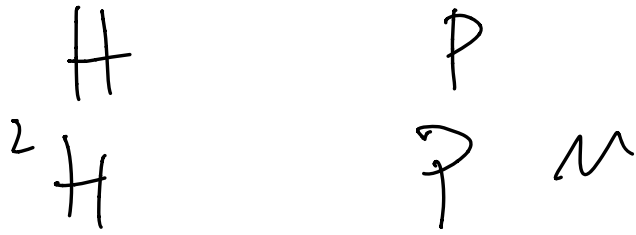
RAGGIO

$$b \propto \frac{1}{\tan \theta/2}$$

$$V = \frac{4}{3} \pi R^3 \propto A$$

$$R = R_0 A^{1/3}$$

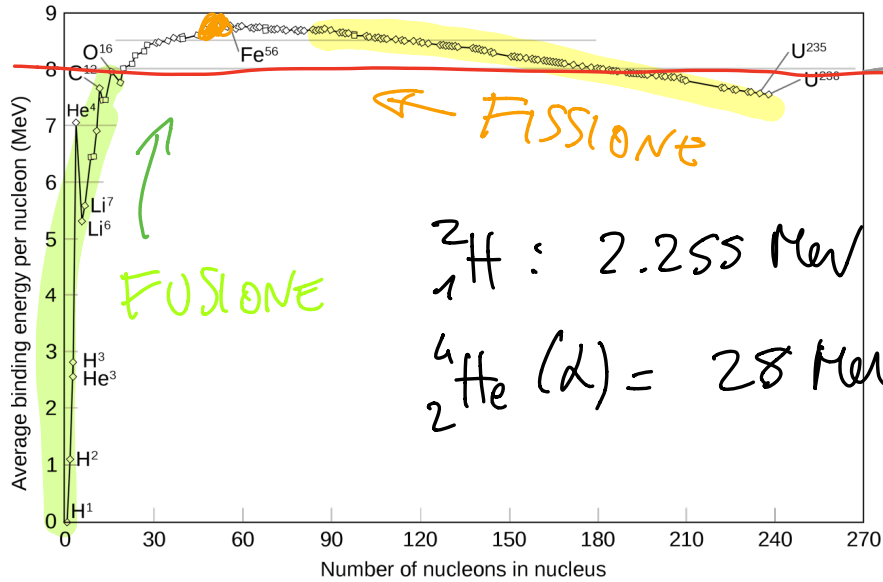
$\underbrace{\quad}_{\sim 1.2 \text{ fm}}$



(c=1)

$$M(A, Z) = Z m_p + (A - Z) m_n - E_b$$

$\frac{E_b}{A} \uparrow$



$${}^2_1\text{H} : 2.255 \text{ MeV}$$

$${}^4_2\text{He} (\alpha) = 28 \text{ MeV}$$

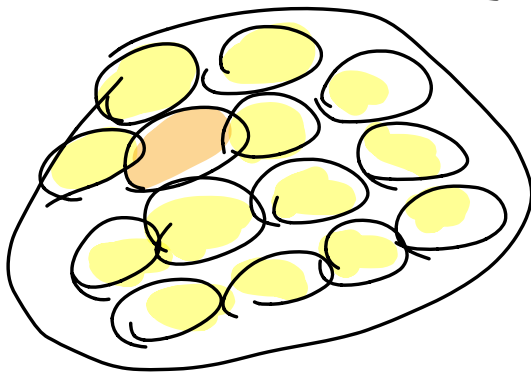
8 MeV

$\longrightarrow A$

$$E_L \propto A$$

MODELLO NON MD : a goccia

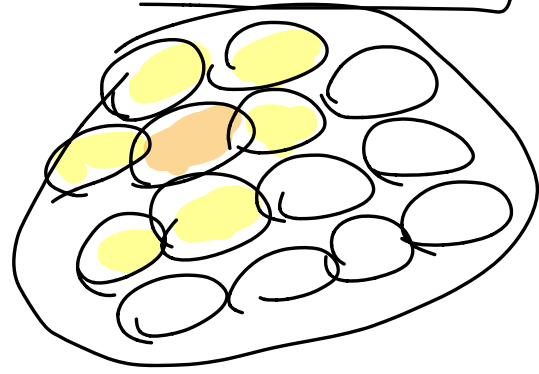
LUNGO NAGGIO



es. EM

(VOLUME)  $E_L \propto A(A-1) \propto A^2$

COLTO NAGGIO



es. FORTE

$$E_L \propto A \equiv a, A$$

(SUPERFICIE)  $E_L \propto R^2 = -A^{2/3} \equiv -a_2 A^{2/3}$

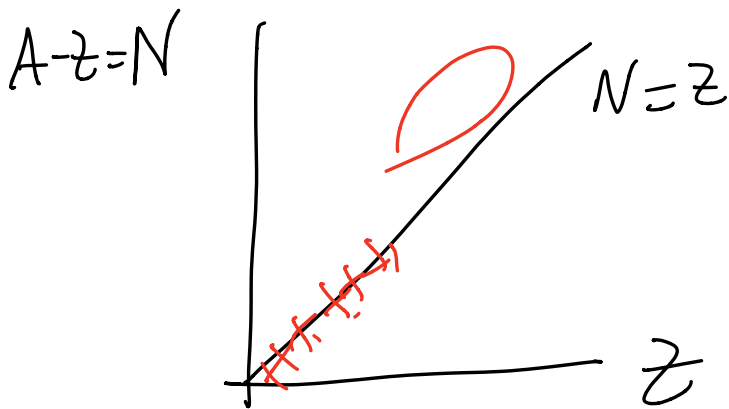
• i poteri si desprendono

$$E_L \propto -\frac{3}{5} \frac{Q^2}{R \cdot 4\pi\epsilon_0}$$

$$\begin{aligned}
&= -\frac{3}{5} \cdot \frac{Z^2 e^2}{R_0 A^{1/3} \cdot 4\pi\epsilon_0} \\
&= -\frac{3}{5} \cdot (\lambda_{tc}) \frac{Z^2}{R_0 A^{1/3}} \\
&\approx -0.7 \text{ MeV} \cdot \frac{Z^2}{A^{1/3}} \\
&\equiv -a_3 \frac{Z^2}{A^{1/3}}
\end{aligned}$$

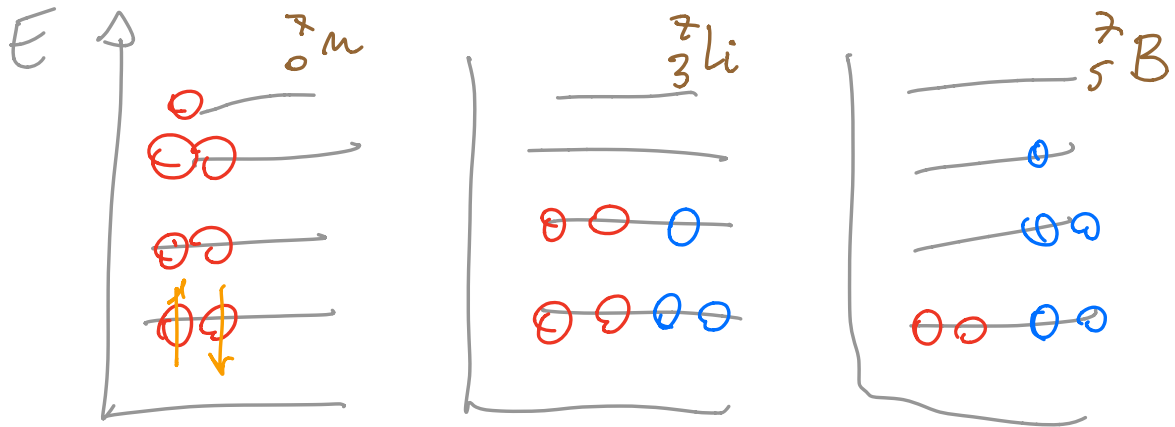
Modello A Gocce

$$E_L = a_1 A - a_2 A^{2/3} - a_3 \frac{Z^2}{A^{1/3}}$$



tenice di Pauli

$$\begin{aligned}
&N-Z = (A-Z) - Z \\
&E_L = -a_4 \frac{(A-2Z)^2}{A}
\end{aligned}$$



⊙ = n

⊙ = p

$$A = 7$$

A	N	Z	NUCLEI STABILI
PARI	PARI	PARI	156
PARI	DISPARI	DISPARI	5
DISPARI	PARI	D	48
DISPARI	D	PARI	50

→ ODD - ODD è RARO

$$E_L = \pm a_s A^{-3/4}$$

0 per A dispari

-as  
+as

ODD - ODD  
EVEN - EVEN

## BETHE-WEIZSÄCKELN

$$E_L = a_1 A - a_2 \cdot A^{2/3} - a_3 \frac{Z^2}{A^{1/3}} - a_4 \frac{(A-2Z)^2}{A} \pm a_5 A^{-3/4}$$

$$a_1 = 16 \text{ MeV}$$

$$a_2 = 17 \text{ MeV}$$

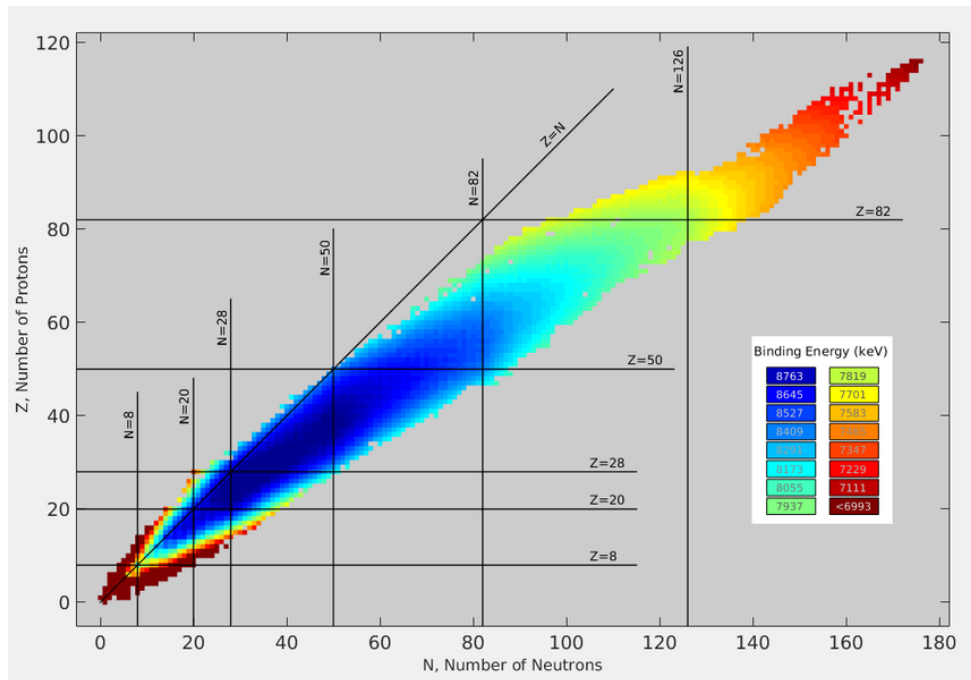
$$a_3 = 0.7 \text{ MeV}$$

$$a_4 = 24 \text{ MeV}$$

$$a_5 = 34 \text{ MeV}$$

$$M(A, Z) = Z m_p + (A-Z) m_n - E_L(A, Z)$$

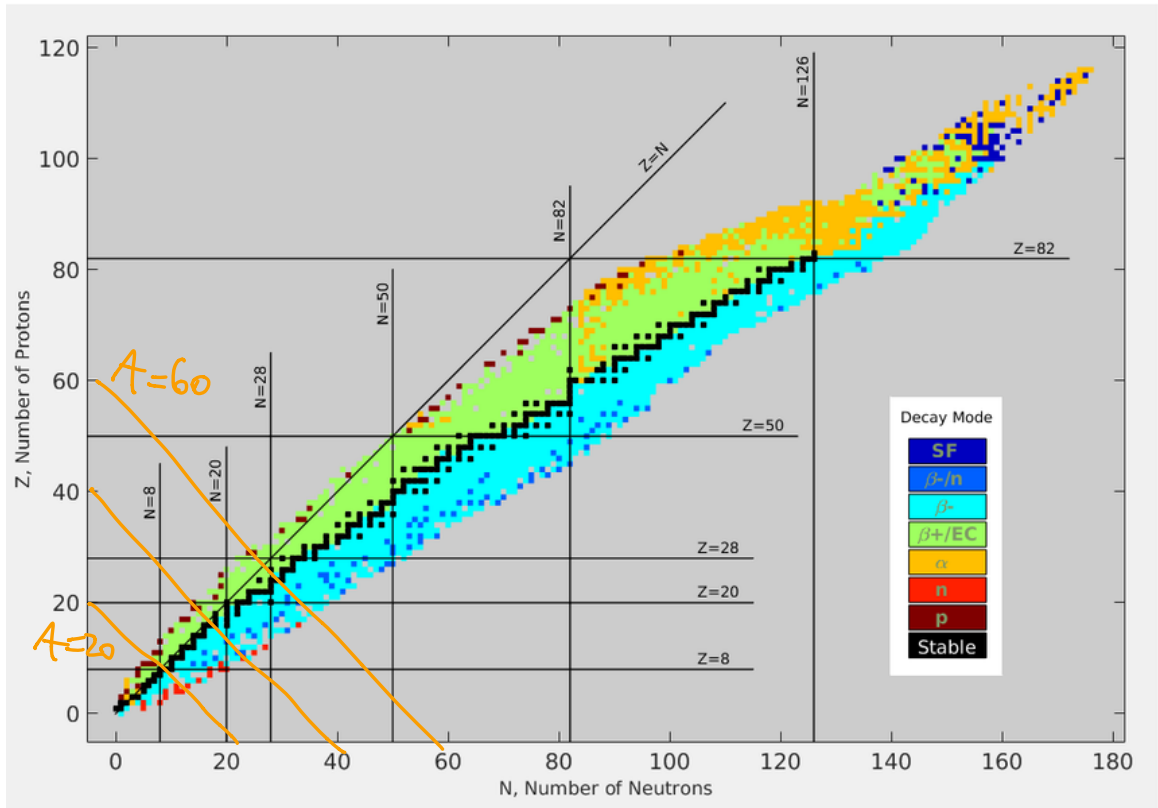
CARTA  
DI  
SEGRE'



asse z:  $Z$   
asse x:  $N$   
asse y:  $Z$

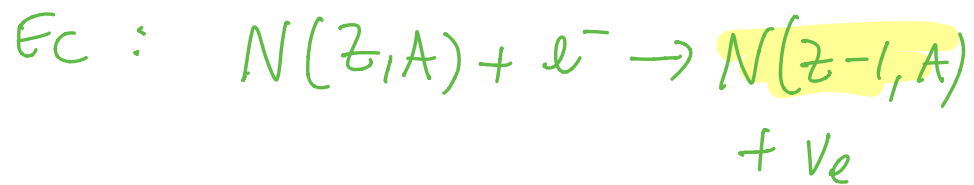
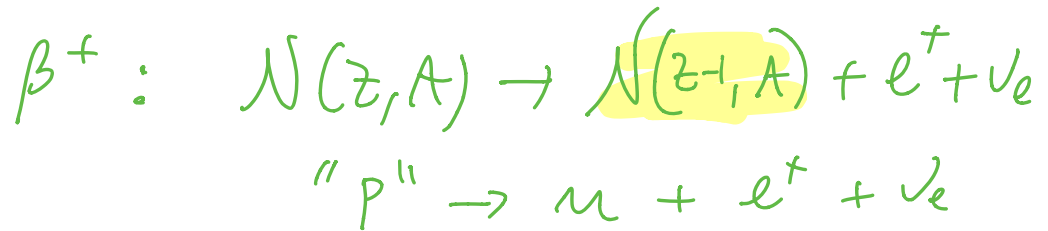


$Z \uparrow$



$\rightarrow N$

$A$  is constant  $\rightarrow Z + N = A$



$$\beta^- : N(z, A) \rightarrow N(z+1, A)$$

$$+ e^- + \bar{v}_e$$

Stabilità: fissato  $A$ ,

$$\left. \frac{\partial E_L}{\partial z} \right|_{A=\text{cost}} = 0 = 2a_3 \frac{z}{A^{1/3}} - 4a_4 \frac{(A-2z)}{A}$$

$$z = \frac{4a_4}{\frac{2a_3}{A^{1/3}} + \frac{8a_4}{A}} = \frac{2a_4 A}{4a_4 + a_3 A^{2/3}}$$

per  $A$  piccolo

$$z \simeq \frac{A}{2}$$

per  $A$  grande

$$z \simeq \frac{2a_4}{a_3} A^{1/3}$$

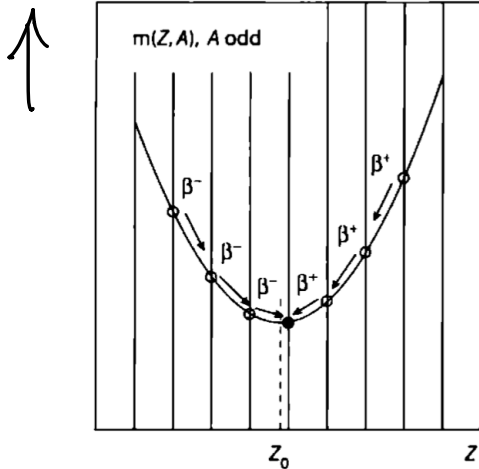
$E_L$  è una parabola in  $z$   
(fissato  $A$ )

$$E_L(A, z) = a_1 A - a_2 A^{2/3} - a_3 \frac{z^2}{A^{1/3}} - a_4 \frac{(A-zz)^2}{A} \pm a_5 A^{-3/6}$$

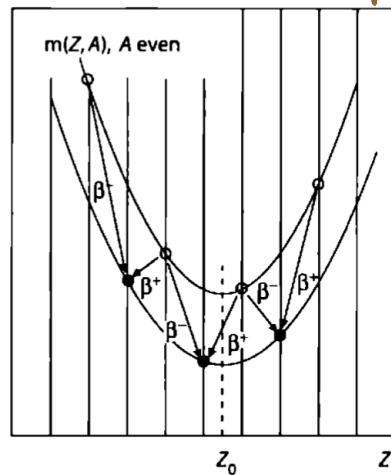
se plots

$$M(A, z) = z \omega_p + (A-z) \omega_m - E_L(A, z)$$

*A dispari*



*A pari*

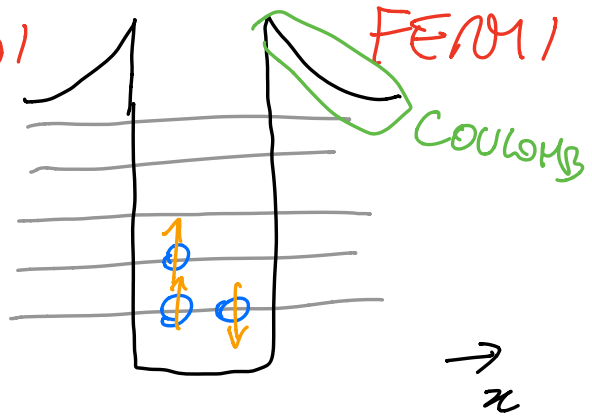
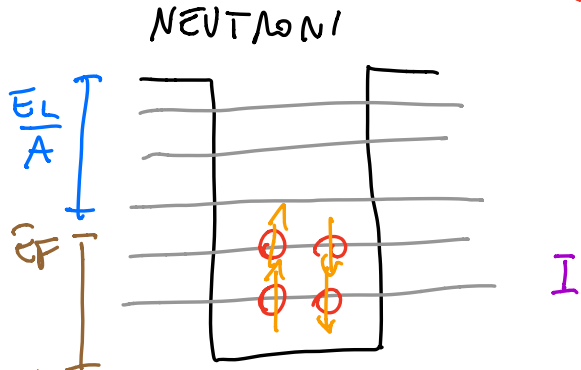


M

$z$

$z$

# Modello A CAS DI



energia di Fermi  
 $E_C/A \sim 8 \text{ MeV}$

quantità  $\psi$   
 dentro il volume del  
 nucleo

$$du = \frac{d^3p}{(2\pi\hbar)^3}$$

$$= \frac{p^2 dp d\Omega}{(2\pi\hbar)^3} \cdot V = \frac{p^2 dp \cdot 4\pi}{(2\pi\hbar)^3} V$$

$$V = \frac{4}{3}\pi R^3 = \frac{4}{3}\pi R_0^3 A$$

assumiamo  $A = Z + N = 2Z = 2N$

SPIN SU  
 SPIN GIU'  $\rightarrow Z = A/2 = N$

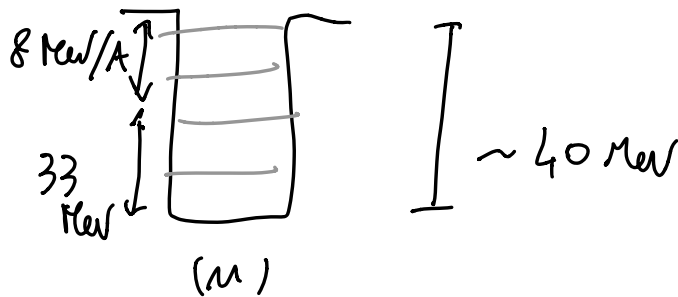
$$Z = N = \frac{A}{2} = 2 \int_0^{p_F} du = 2 \int_0^{p_F} \frac{p^2 dp \cdot 4\pi}{(2\pi\hbar)^3} \cdot \frac{4}{3}\pi R_0^3 A$$

$$= \frac{3/2 \pi^4}{8} \frac{1}{(\pi \hbar)^3} \cdot \frac{1}{3} p_F^3 R_0^3 A \pi$$

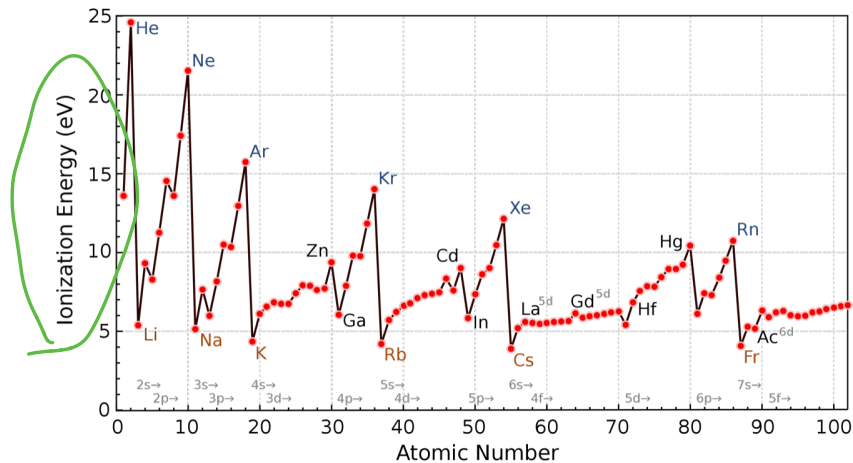
$$P_F = \left( \frac{\hbar}{R_0} \right) \left[ \frac{9\pi}{8} \right]^{1/3} = 250 \text{ MeV/c}$$

$$\underbrace{\quad}_{1.2 \text{ fm}} \quad \hbar c = 197 \text{ MeV}\cdot\text{fm}$$

$$E_F = \frac{P_F^2}{2m_p} = 33 \text{ MeV}$$



## MODELLO A SHELL



NUMERI MAGICI : valori di  $Z$  (o  $N$ )  
 per cui  $E_c \approx 0$   
 particolarmente  
 alte

OMEGMA: il più stabile il  
 nucle

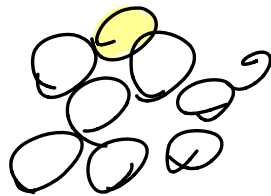
$$N = 2, 8, 20, 28, 50, 82, 126$$

livelli energetici:  $n$   $l$   $m$   
 $l \in \{0, n-1\}$   
 $m \in \{-l, l\}$

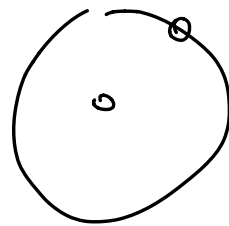
$$2 \times \sum_{l=0}^{n-1} (2l+1)$$

$$= 2n^2$$

MA



$\neq$



difficile MA possibile

→ solve  $V$

$V$	results
-----	---------

buckingham

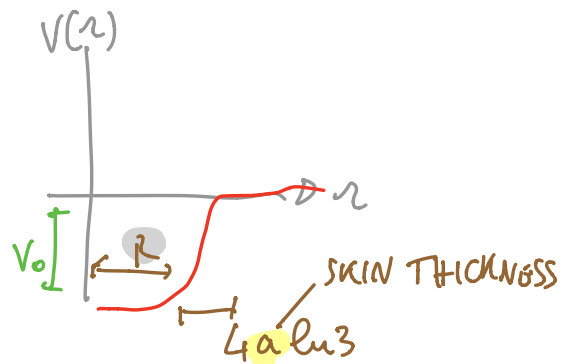
$\ddot{n}$

Armas

2, 8, 20, 40, 70, 112

Woods-Saxon

$$V(r) = \frac{-V_0}{1 + e^{\frac{r-R}{a}}}$$



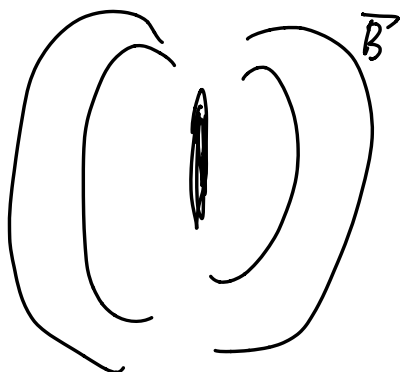
2, 8, 20, 40, 58, 92,

112

NO! ::

→ solve spin-orbit

$$\Delta \vec{E} \propto \vec{\mu} \cdot \vec{B} \propto \vec{L} \cdot \vec{S}$$



$$\vec{B} = \frac{\mu_0}{4\pi} \left[ \frac{3\hat{r}(\vec{\mu} \cdot \hat{r})}{r^3} - \frac{\vec{\mu}}{r^3} \right]$$

$$U = -\vec{\mu} \cdot \vec{B}$$

$$= \sim \text{keV}$$

$$\text{radius } r^3 \sim R^3$$

$$\sim 1.2 \mu\text{m}^3$$

in fssic atoms:

$$\Delta E \sim \text{meV}$$

	harmonic potential	Woods-Saxon potential	spin-orbit coupling		
li 2g 3d 4s	—168—				
		li	1i <sub>13/2</sub>	14	<b>126</b>
			3p <sub>1/2</sub>	2	112
1h 2f 3p	—112—	3p	3p <sub>3/2</sub>	4	110
		2f	2f <sub>5/2</sub>	6	106
			2f <sub>7/2</sub>	8	100
		1h	1h <sub>9/2</sub>	10	92
			1h <sub>11/2</sub>	12	<b>82</b>
1g 2d 3s	—70—	3s	3s <sub>1/2</sub>	2	70
		2d	2d <sub>3/2</sub>	4	68
			2d <sub>5/2</sub>	6	64
		1g	1g <sub>7/2</sub>	8	58
			1g <sub>9/2</sub>	10	<b>50</b>
1f 2p	—40—	2p	2p <sub>1/2</sub>	2	40
		1f	1f <sub>5/2</sub>	6	38
			2p <sub>3/2</sub>	4	32
			1f <sub>7/2</sub>	8	<b>28</b>
1d 2s	—20—	2s	1d <sub>3/2</sub>	4	<b>20</b>
		1d	2s <sub>1/2</sub>	2	16
			1d <sub>5/2</sub>	6	14
1p	—8—	1p	1p <sub>1/2</sub>	2	<b>8</b>
			1p <sub>3/2</sub>	4	6
1s	—2—	1s	1s <sub>1/2</sub>	2	<b>2</b>
	$\Sigma 2(2l+1)$		$2j+1$		$\Sigma 2j+1$

low splitting  $\bar{e}$   
(dispense)

$$J = L + S$$

$$L \quad S = 1/2$$

$$\Delta E = l + 1/2$$