

ELEMENTARY PARTICLE PHYSICS
Current Topics in Particle Physics
Laurea Magistrale in Fisica,
curriculum Fisica Nucleare e Subnucleare
Lecture 5

Simonetta Gentile*

* Università Sapienza, Roma, Italia.

October 21, 2018

Simonetta Gentile

terzo piano Dipartimento di Fisica *Guglielmo Marconi*

Tel. 0649914405

e-mail: simonetta.gentile@roma1.infn.it

pagina web: <http://www.roma1.infn.it/people/gentile/simo.html>

Bibliography

♠ Bibliography

- K.A. Olive et al. (Particle Data Group), *The Review of Particle Physics*, Chin. Phys. C, 38, 090001 (2014)(PDG) update 2015, <http://pdg.lbl.gov/>
- F. Halzen and A. Martin, *Quarks and Leptons: An introductory course in Modern Particle Physics*, Wiley and Sons, USA(1984).

♠ Other basic bibliography:

- A.Das and T.Ferbel, *Introduction to Nuclear Particle Physics* World Scientific,Singapore, 2nd Edition(2009)(DF).
- D. Griffiths, *Introduction to Elementary Particles* Wiley-VCH,Weinheim, 2nd Edition(2008),(DG)
- B.Povh et al., *Particles and Nuclei* Springer Verlag, DE, 2nd Edition(2004).(BP)
- D.H. Perkins,*Introduction to High Energy Physics* Cambridge University Press, UK, 2nd Edition(2000).
- Y.Kirsh & Y. Ne'eman, *The Particle Hunters* 

♠ Particle Detectors bibliography:

- William R. Leo *Techniques for Nuclear and Particle Physics Experiments*,
Springer Verlag (1994)(LEO)
- C. Grupen, B. Shawartz *Particle Detectors*,
Cambridge University Press (2008)(CS)
- *The Particle Detector Brief Book*,(BB)
<http://physics.web.cern.ch/Physics/ParticleDetector/Briefbook/>

Specific bibliography is given in each lecture

Lecture Contents - 1 part

1. Introduction. Lep Legacy
2. Proton Structure
3. Hard interactions of quarks and gluons: Introduction to LHC Physics
4. Collider phenomenology
5. The machine LHC
6. Inelastic cross section pp
7. W and Z Physics at LHC
8. Top Physics: Inclusive and Differential cross section $t\bar{t}$ W, $t\bar{t}$ Z
9. Top Physics: quark top mass, single top production
10. Dark matter
 - Indirect searches
 - Direct searches

Specific Bibliography

- ♠ Bibliography of this Lecture
 - ATLAS, ATLAS experiment
 - CMS, CMS experiment

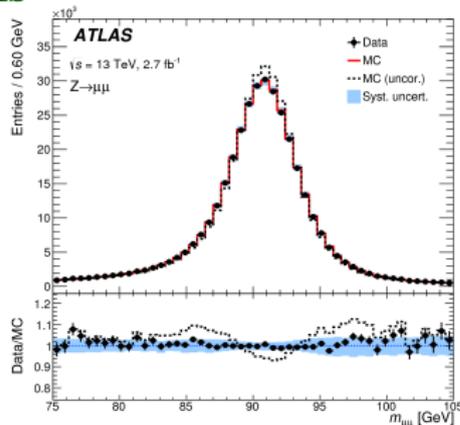
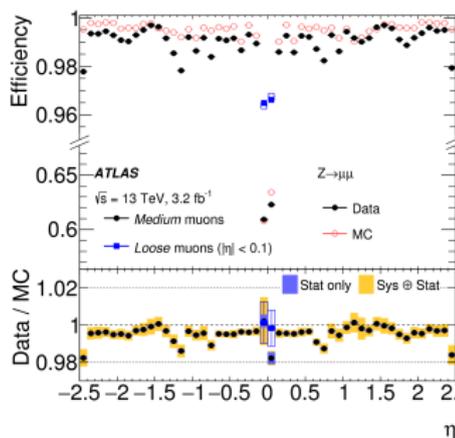
1 Identification of jets and leptons

2 Expected physics

Identification of objects: muons

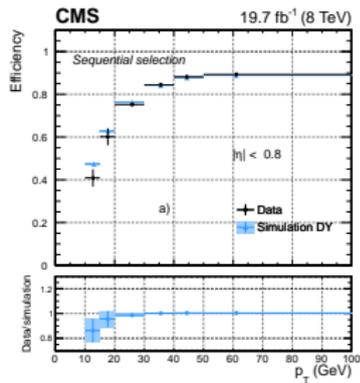
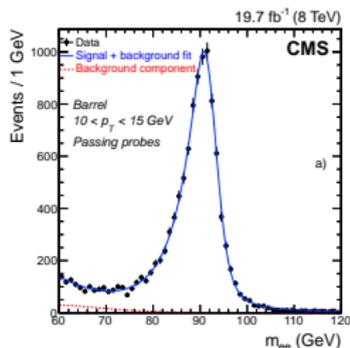
♠ Key points towards the discovery is to prove the capability (and quality) of **identification of objects** (electron, μ , tau, photons, jets, missing transverse energy) and their reconstruction and the capability of **reproduce the expected physics**.

muons



- Reconstruction **efficiency** $\approx 99\%$ ($|\eta| < 2.5$ and $5 < p_T < 100 \text{ GeV}$)
- **Momentum resolution** 1.7% (2.3%) for muons from $J/\psi \rightarrow \mu\mu (Z \rightarrow \mu\mu)$

electrons



- The **momentum scale** is calibrated with an uncertainty smaller than **0.3%**.
- The **momentum resolution** for electrons produced in Z boson decays ranges from **1.7 to 4.5%**.

Identification of objects: photons

photons

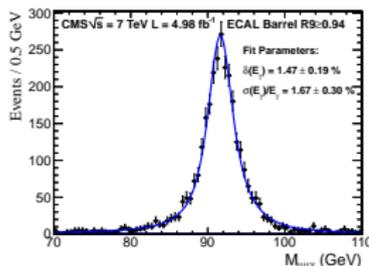
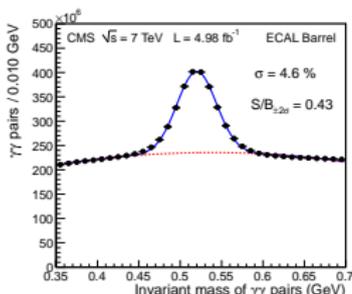


Figure : $Z \rightarrow \mu\mu\gamma$ events, deviation of the reconstructed photon energy from that expected from the decay kinematics δ , mean energy resolution.

Figure : an example of the η -meson peak reconstructed from the invariant mass of photon pairs in barrel,

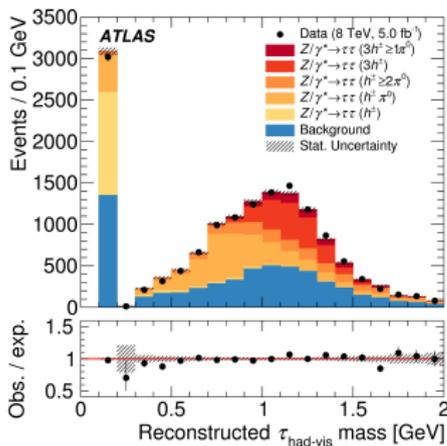
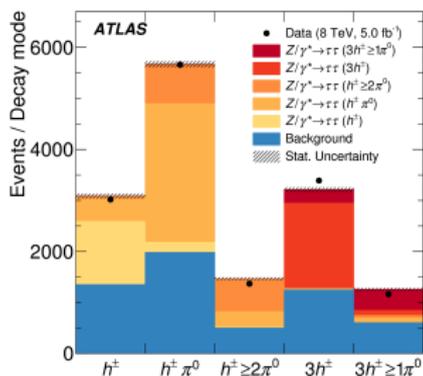
$$(m_\eta = 547.86 \text{ MeV}, \Gamma_\eta = 1.31 \text{ keV}, Br(\gamma\gamma) \rightarrow 72.12\%)$$

- The derived **energy resolution** for photons from 125 GeV Higgs boson decays varies across the barrel from **1.1%** to 2.6% and from 2.2% to 5% in the endcaps(CMS).

Identification of objects: taus

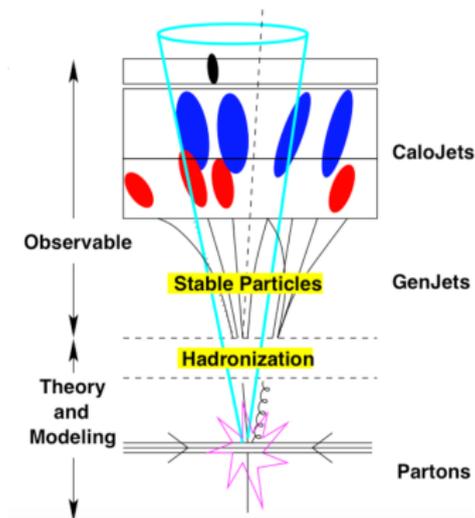
Lepton τ : $m_\tau = 1776$ MeV lifetime $\tau \simeq 290 \cdot 10^{-15}$ s \implies
 $c\tau = 87.03\mu\text{m}$

main decays: $\tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau (\sim 17.4\%), \tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau (\sim 17.4\%), \tau^- \rightarrow \pi^- \nu_\tau (\sim 10.8\%), \tau^- \rightarrow K^- \nu_\tau (\sim 7\%)$ and
 $\tau^- \rightarrow h^- \nu_\tau (\simeq 11\%), \tau^- \rightarrow h^- h^+ h^- \nu_\tau (\simeq 15\%),$



Identification of objects: jets

- What are jets?



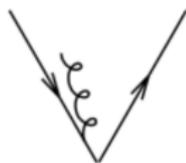
- Collimated bunches of stable hadrons, originating from **partons** (quarks & gluons) after **fragmentation** and **hadronization**
- Jet Finding is the **approximate** attempt to reverse-engineer the quantum mechanical processes of fragmentation and hadronization. **Not a unique procedure** \Rightarrow several different approaches.
- Jets are the observable objects to relate experimental **observations** to theory **predictions** formulated in terms of quarks and gluons

Identification of objects: jets

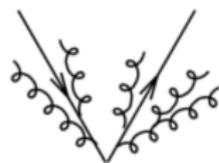
A **hadronic jet is a collimated cone of particles associated with a final state parton**, produced through fragmentation. The jet has a four momentum: $E = \sum_i E_i, \vec{p} = \sum_i \vec{p}_i$. Where the constituents (i) are hadrons detected as **charged tracks** and **neutral energy** deposits.



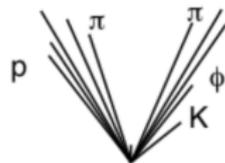
LO partons



NLO partons



parton shower



hadron level

How to define experimentally a *jet*?

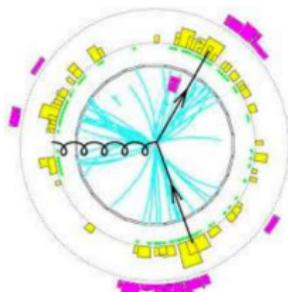
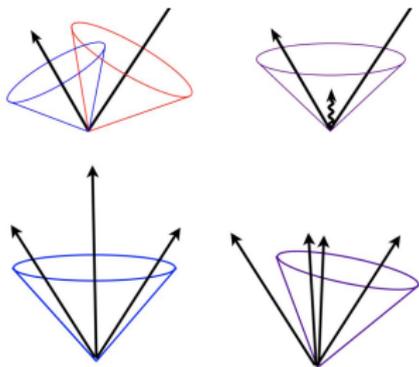
A few different approach.

- **Cone algorithm**: include all particles inside a cone of given radius
- **Particle Flow Algorithms** (used at LHC, next generation of collider): for each individual particle in a jet, use the detector part with the best energy resolution.

Identification of objects: jets

typical jet:

- $\sim 60\%$ charged particles
- $\sim 30\%$ photons
- $\sim 10\%$ neutral hadrons
- $\sim 1\%$ neutrinos



Identification of objects: jet algorithm

K_T /anti K_T algorithm.

Based on the following distance measures:

- Distance d_{ij} between two particles i, j :

$$d_{ij} \equiv \min(p_{ti}^{2p}, p_{tj}^{2p}) \frac{\Delta R_{ij}^2}{R^2}$$

$$\Delta R_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2$$

- Distance between any particle i and beam (B) d_{iB}

$$d_{iB} \equiv p_{T,i}^{2p}$$

- From the list of all final state particles, determine all the distances d_{ij} and d_{iB}
- Find the minimum distance
- If smallest is a d_{ij} recombine particles i and j (sum four momenta) and go back to the first step: update distance and find again smallest
- If smallest is a d_{iB} , declare particle i to be a jet and remove from the list of particles.
- Stop the algorithm when no particles remain and all

Identification of objects: jet algorithm

K_T /anti K_T algorithm.

Based on the following distance measures:

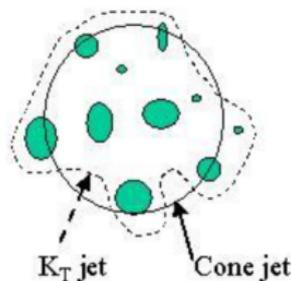
- Distance d_{ij} between two particles i, j :

$$d_{ij} \equiv \min(p_{ti}^{2p}, p_{tj}^{2p}) \frac{\Delta R_{ij}^2}{R^2}$$

$$\Delta R_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2$$

- Distance between any particle i and beam (B) d_{iB}

$$d_{iB} \equiv p_{T,i}^{2p}$$

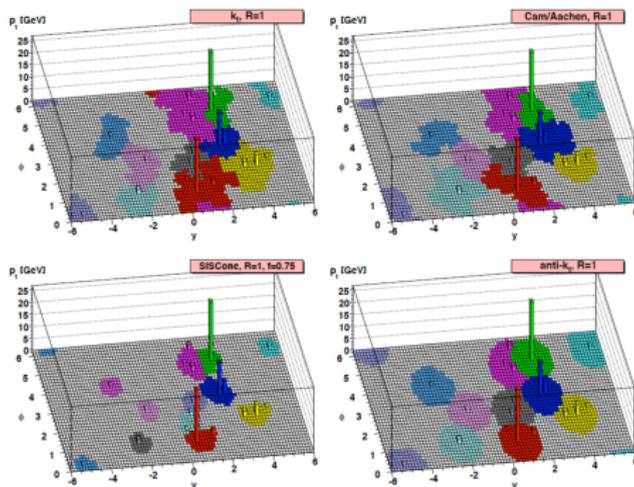


- Parameter R : Scales the d_{ij} w.r.t. the d_{iB} such that any pair of final jets a and b are at least separated by $\Delta_{ab}^2 = R^2$
- Parameter p : governs the relative power of energy vs geometrical scales to distinguish the three algorithms K_T , anti- K_T ...

Comparison jet algorithm

$p=+1$: kT algorithm. $p=-1$: anti-kT algorithm (favoring recombination of high-pt particles)

- It is important to emphasize that sequential recombination algorithms **can be applied** exactly in the same way to **QCD partons, hadrons or at the detector level**. Allows straightforward comparison between data and theory. **anti K_T** algorithm used at LHC gives **more regular jet**.



Comparison jet algorithm

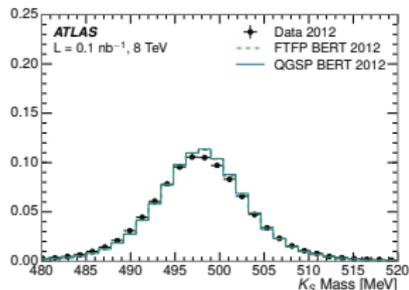


Figure : $m_{K^0} \simeq 497.6$ MeV.
Decays $\pi^0\pi^0$ (39%) $\pi^+\pi^-$ (69%)

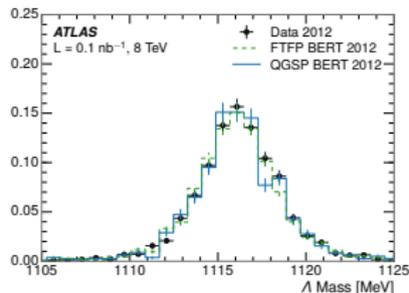


Figure : $m_{\Lambda} \simeq 1115$ MeV. Decays
 $p\pi^-$ (63.9%) $n\pi^0$ (35.8%)

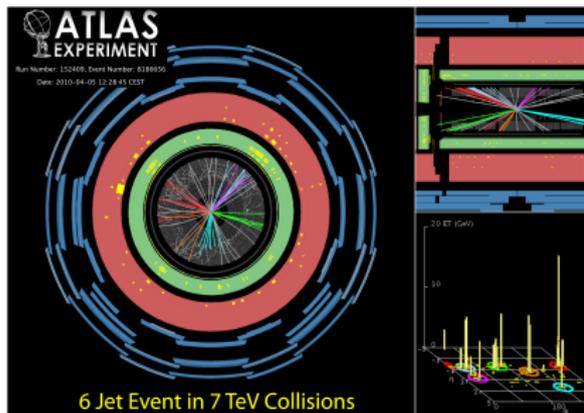
Issues for Jet Observables

- Jet triggers and Jet selection
 - lower p_T thresholds, matching sample with different triggers
- Choice of Jet algorithm and jet size
 - use different algorithm (collinear safe algorithms)
 - Standard in CMS: anti-KT, $R= 0.5, 0.7$; standard in ATLAS: anti-KT, $R= 0.4, 0.7$
- Jet Energy Scale
 - Absolute and relative (as function of rapidity)
 - jet cross section falls like power law, power $\approx -5-6$
 - great progress reached $< 1\%$ ¹
- Jet energy resolution
 - smearing of distribution
 - CMS jet energy resolutions at η central 15-20% at 30 GeV, about 10% at 100 GeV, and 5% at 1 TeV

Comparison with theory at "hadron (or particle) level"

¹CMS Coll Jet energy scale and resolution in the CMS experiment in pp collisions at 8 TeV, subm J of Instr [arXiv:1607.03663](https://arxiv.org/abs/1607.03663)

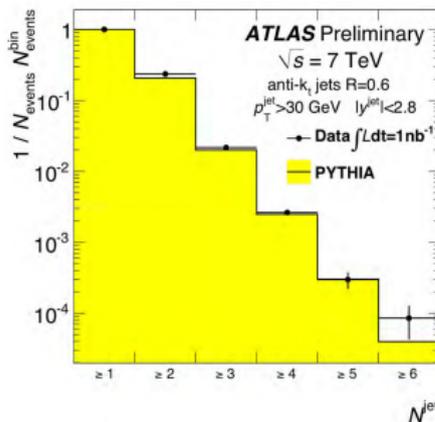
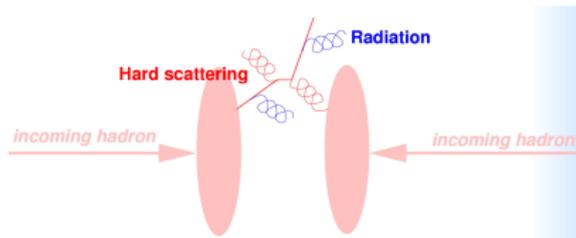
A six-jets event at $\sqrt{s} = 7$ TeV



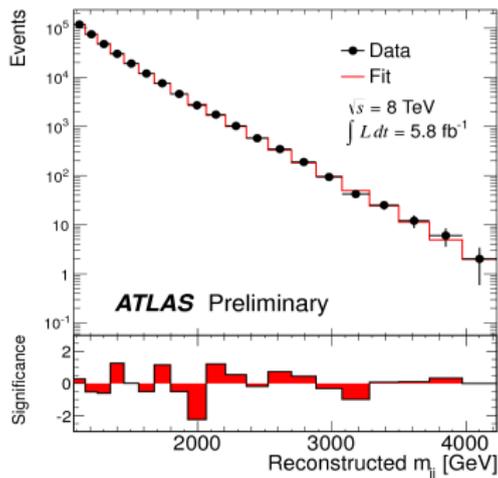
• Jet multiplicity

- Another possible test of QCD is obtained by checking the jet multiplicity.
- tests also the modelling of the radiation.

Figure : 6 Jet event in 7 TeV Collisions



Jet production



- Crucial energy measurements in the quest for new physics
- ATLAS and CMS use different approaches:
 - Calorimetric (ATLAS)
 - Particle Flow (CMS)

Particle flow use the most appropriate system to measure particle in the event.

Jet composition:

~ 60% charged particles

~ 30% photons

~ 10% neutral hadrons

~ 1% neutrinos

Particle Flows basic

- **Jet composition** : $\sim 60\%$ charged particles, $\sim 30\%$ photons, $\sim 10\%$ neutral hadrons, $\sim 1\%$ neutrinos
- **Traditional calorimetric approach**
 - Measure all components of jet energy in ECAL/HCAL
 - $\sim 70\%$ of energy measured in HCAL, limits jet energy resolution

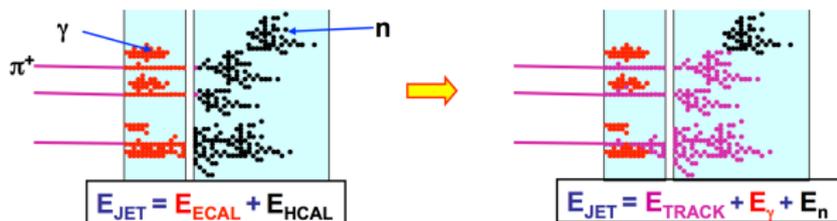


Figure : Particle Flow paradigm

- **Particle Flow Calorimetry paradigm :**
 - charged particles measured in tracker (essentially perfectly)
 - Photons in ECAL
 - Neutral hadrons (ONLY) in HCAL
 - **Only 10 % of jet energy from HCAL** \implies much improved resolution

Particle Flows basic

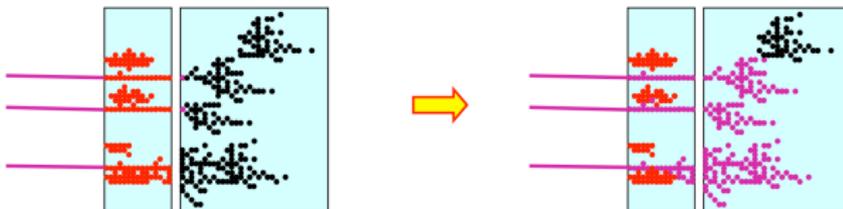
Hardware

Need to be able to **resolve energy deposits from different particles** \Rightarrow **Highly granular detectors**



Software

Need to be able to **identify energy deposits from individual particles** \Rightarrow **Sophisticated reconstruction software**



Particle Flows technique

Used in Aleph and CMS.

Use the **best system** you have to measure particle in the event

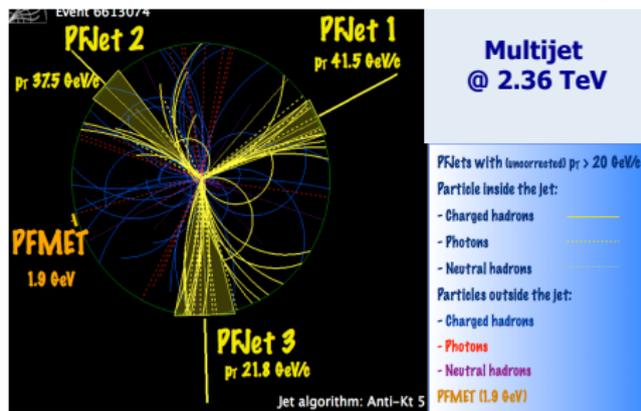


Figure : CMS-Particle Flow jets ²

CMS: high B, excellent tracker, granular ECAL



Strong improvement in JET and Missing energy (MET) resolution

²kindly from R.Paramatti presentation

CMS particle ID in Calorimeters

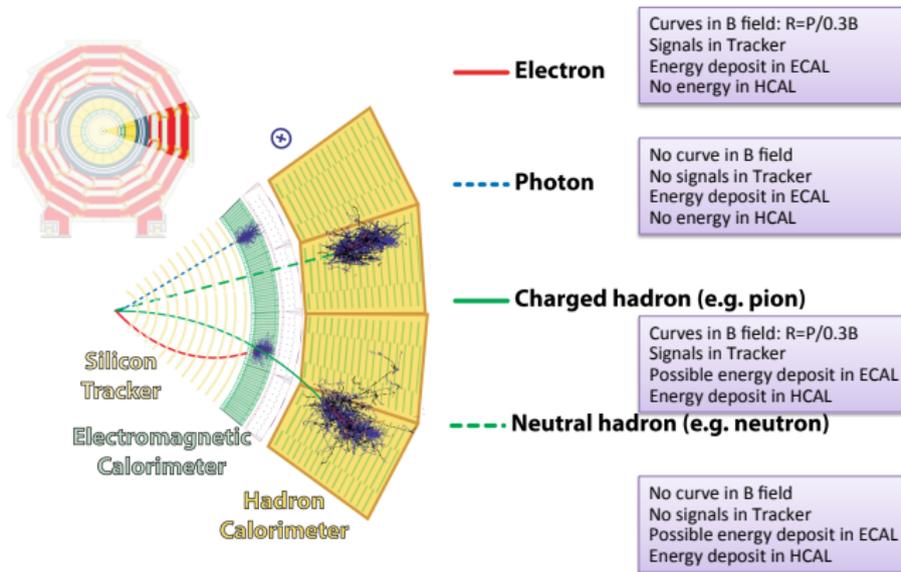


Figure : CMS-Particle ID in Calorimeters ³

³kindly from R.Paramatti presentation

Jet Energy resolution

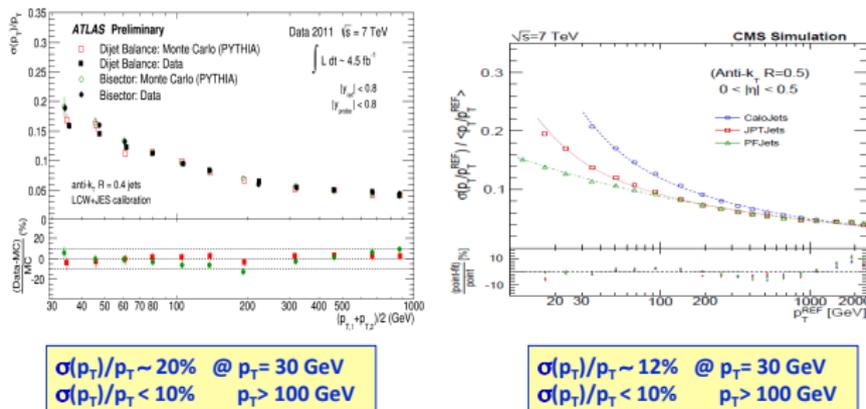


Figure : ATLAS CMS jet resolution ⁴

Similar performance in the region relevant for new physics with CMS Particle Flow. Atlas is better comparing calorimetric resolution

⁴kindly from R.Paramatti presentation

Particle Flow vs Calorimetric Jets in CMS

$$\text{Jet response} : (p_T^{\text{rec}} - p_T^{\text{gen}}) / p_T^{\text{gen}}$$

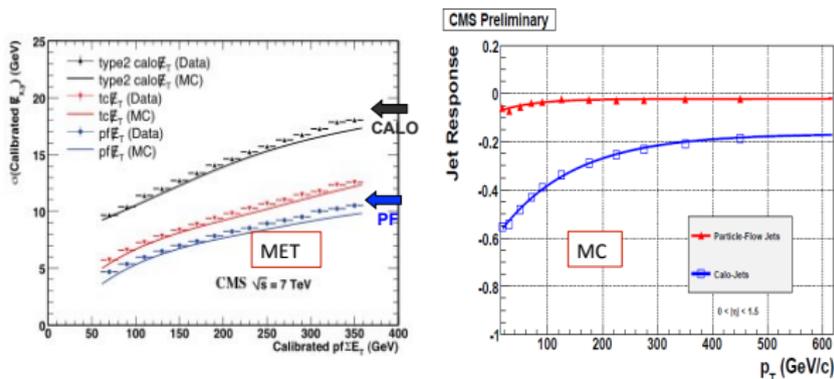


Figure : Resolution for the calibrated MET for multijet events with two jets with $p_T > 25$ GeV (left) $\approx 60\%$ of Jet Energy measured with tracks (right)

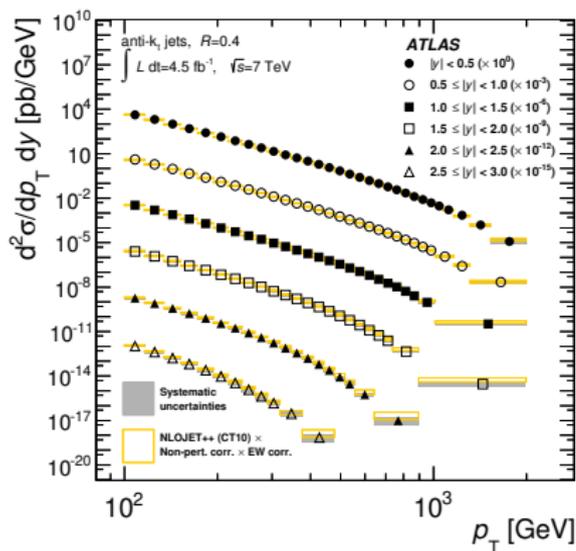


Figure : Double-differential inclusive jet cross-sections as a function of the jet pT in bins of rapidity, for anti-kt jets
 Rate spans 10 order of magnitude⁵

⁵Atlas Coll. Measurement of the inclusive jet cross-section in proton-proton collisions at $\sqrt{s}=7 \text{ TeV}$ using $4.5 \sim \text{fb}^{-1}$ of data with the ATLAS detector JHEP02(2015)153
<http://arxiv.org/abs/1410.8857>

Identification of objects: Transverse missing energy

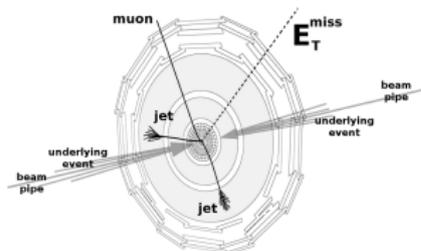
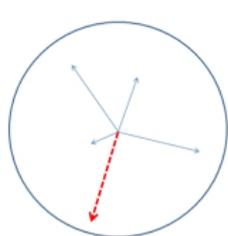
The primary motivation of E_T^{miss} is to provide as complete of a picture as possible of the event, including the presence of one or more energetic neutrinos or other weakly-interacting stable particles though apparent missing energy.

Energetic particles produced in the direction of the beam pipe make it impossible to directly measure missing energy longitudinal to the beam direction, however, the **transverse energy balance can be measured**.

- Transverse momentum, \vec{p}_T is the momentum of an object transverse to the beam.
- Transverse energy is defined as $E_T = \sqrt{m^2 + p_T^2}$.

- The **initial \vec{p}_L in a parton collision is unknown**, because the partons that make up a proton share the momentum.
- The **initial \vec{p}_T was zero**.
- Missing transverse momentum $\vec{p}_T^{miss} = -\sum_i \vec{p}_T(i)$, i extended to all visible particles. $E_T^{miss} = -\sum_i \vec{p}_T(i)$ is missing transverse energy or MET (only for massless missing particle(s))

A significant missing transverse energy results from a non-null vectorial sum of all observed particles momenta coming out of a proton-proton collision: as the colliding protons do not possess any momentum in the direction transverse to the beams, the product of their collision must retain that property; hence the vectorial sum must vanish, give or take the experimental resolution with which momenta are determined by the experimental apparatus.



Finding missing transverse momentum would indicate that new, unaccounted for particle(s) had escaped the detector.

Identification of objects: Transverse missing energy

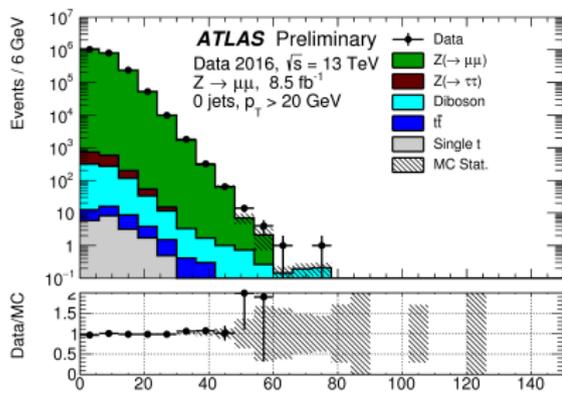


Figure : E_T^{miss} in $Z \rightarrow \mu\mu$ (ATLAS)

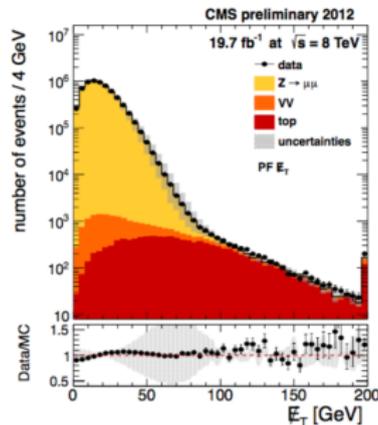
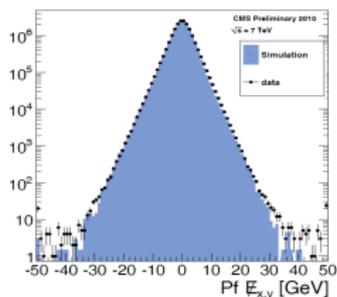
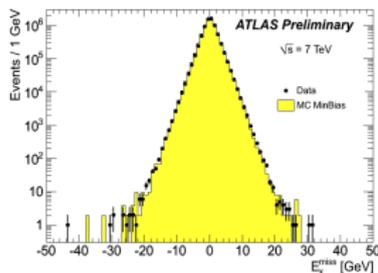


Figure : E_T^{miss} in $Z \rightarrow \mu\mu$ (CMS)

Distributions after detector cleaning, including tails, very well reproduced by MC.

Identification of objects: Transverse missing energy



With the Particle Flow reconstruction, CMS recovers ATLAS performance on jets and MET and circumvents the non brilliant performance in hadron energy measurement

Our master equation

$$\sigma_{\text{meas}} = \frac{N_{\text{obs}} - N_{\text{bkg}}}{\epsilon \mathcal{L}}$$

- N_{obs}
 - Event rate
 - Stat vs syst errors, backgrounds from data or MC?
 - Resolution, Energy scale; Signal significance
- N_{bkg} as N_{obs}
- ϵ
 - **Experimental issues:** Trigger, reconstruction, isolation cuts, low- p_T jets, (jet veto), acceptance, efficiency determination, (tag & probe)
 - **Theoretical issues:** p_T distributions at NL + resummation; differential calculations for detectable acceptance
- \mathcal{L} pp Luminosity uncertainty $< 2\%$

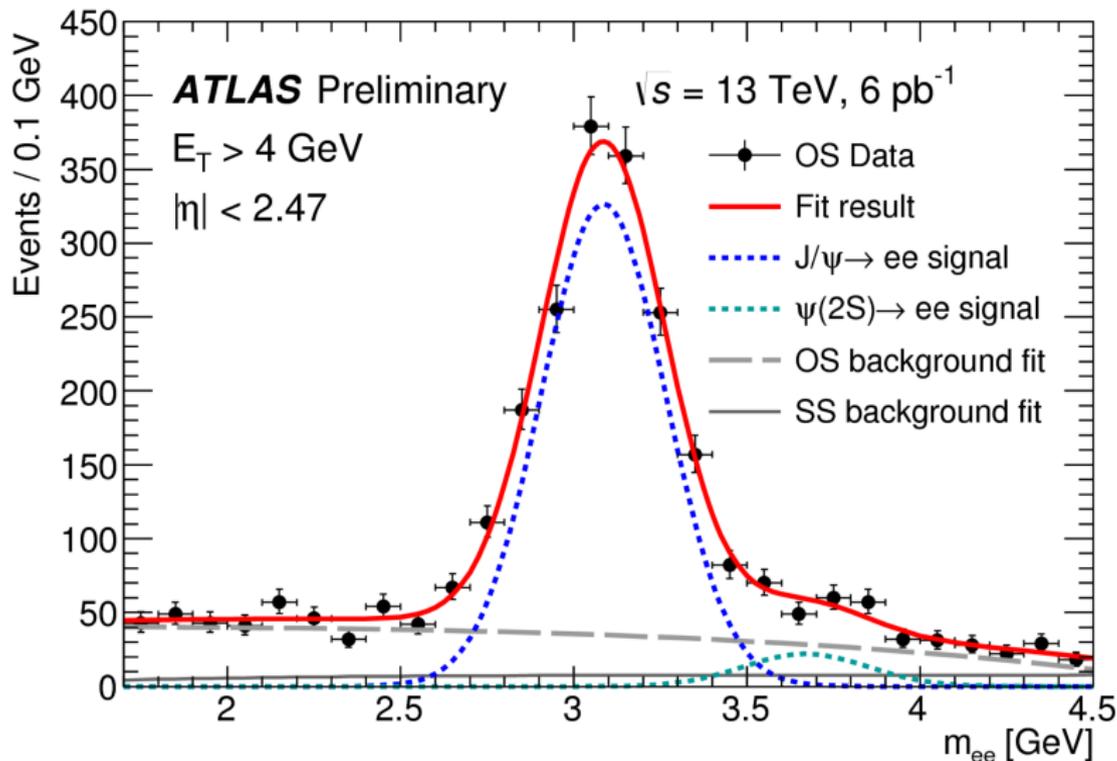
$$\sigma_{\text{theo}} = \text{PDF}(x_1, x_2, Q^2) \otimes \hat{\sigma}_{\text{hard}}$$

- $\text{PDF}(x_1, x_2, Q^2)$, constrains, define uncertainties
- $\hat{\sigma}$ NLO calculations, implemented in MC

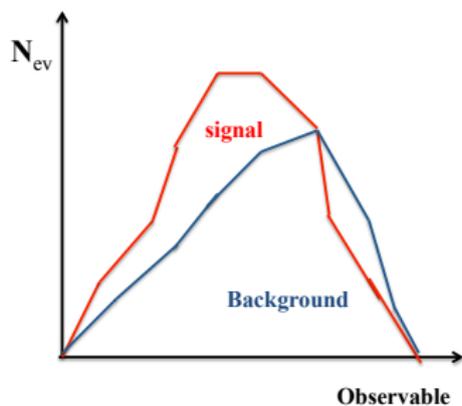
$$\sigma_{\text{theo}} = \text{PDF}(x_1, x_2, Q^2) \otimes \hat{\sigma}_{\text{hard}} \Rightarrow \sigma_{\text{meas}} = \frac{N_{\text{obs}} - N_{\text{bkg}}}{\epsilon \mathcal{L}}$$

The easy case: side bands

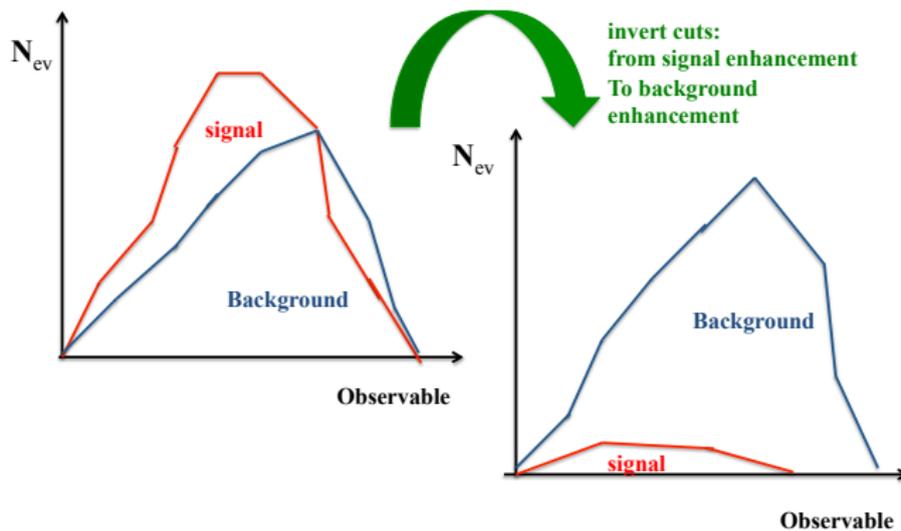
Side bands: a easy case for background evaluation



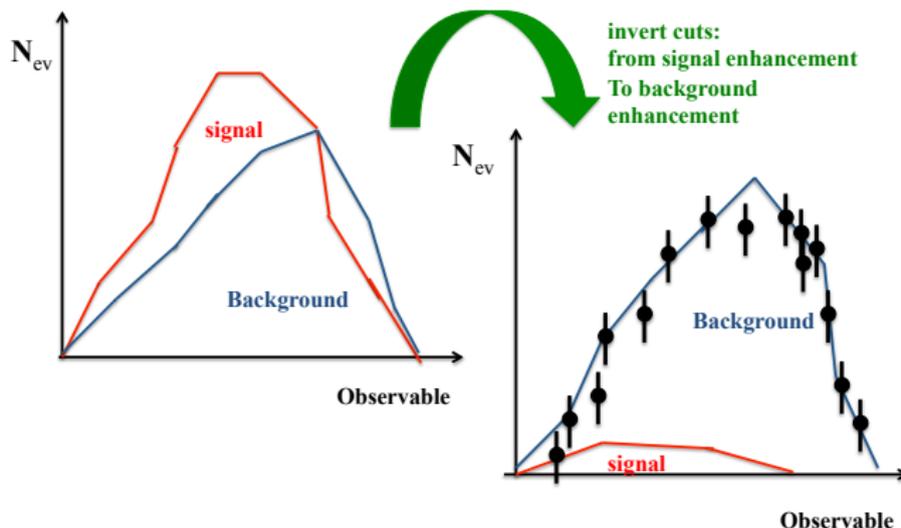
Data-driven Background estimate



Data-driven Background estimate

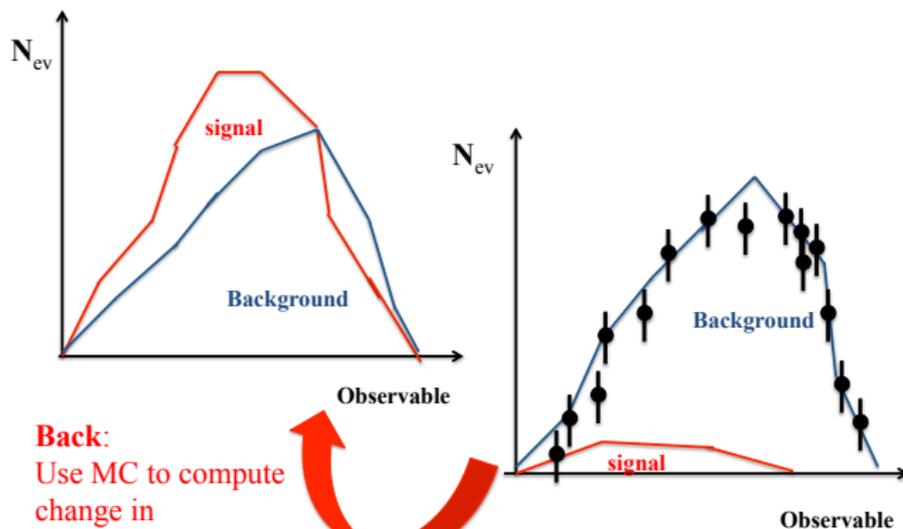


Data-driven Background estimate



Use data to normalize background

Data-driven Background estimate

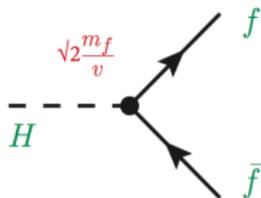


Back:
Use MC to compute
change in
background, when
inverting the cuts

Issues:

- Is left any signal after the inversion?
- Any bias in the control region?
- How well does theory/MC model cut the inversion?

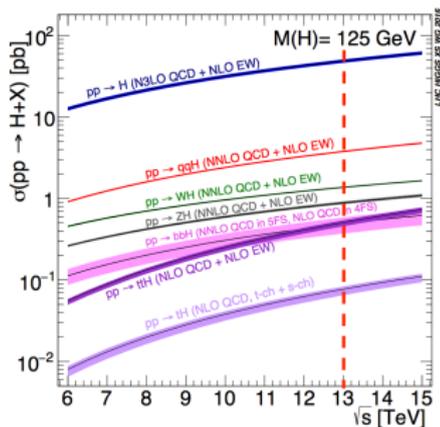
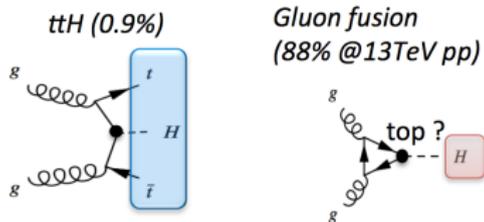
Measuring Yukawa coupling quark top-Higgs



Yukawa coupling: **proportional to fermion mass**. Top is the heavy fermion in SM \rightarrow largest Yukawa coupling $\lambda_t = 2m_t/v \approx 1$.

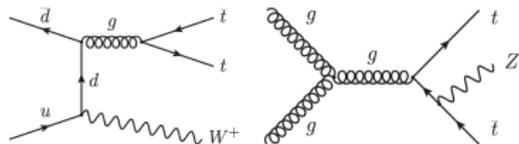
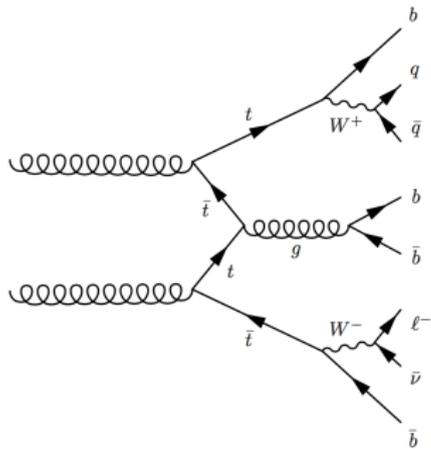
- Indirect constrain on top-Higgs Yukawa coupling derived from gluon-gluon fusion and $H \rightarrow \gamma\gamma$
- $t\bar{t}H$ production: best way to measure top-quark Higgs Yukawa coupling
- **Any deviation might be hint for new physics**

Data-driven Background: example $t\bar{t}H$



- $t\bar{t}H$ production cross-section at $\sqrt{s} = 13$ TeV:
 $0.507 \text{ pb}^{+5.8\%}_{-9.2\%} (\text{QCD scales}) \pm 3.6\% (\text{PDF}, \alpha_s)$
- Only $\sim 1\%$ of the total Higgs cross-section.
- Complex final states, with many objects: jets, b-jets, light leptons (ℓ), hadronic taus (τ_{had}), photons

Data-driven Background: example $t\bar{t}H$



Large irreducible backgrounds:

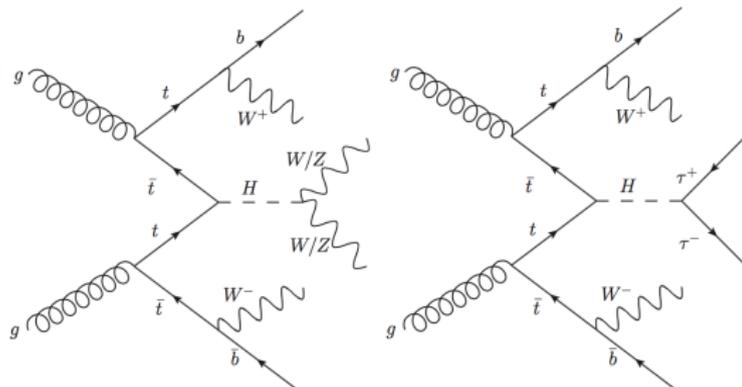
- $t\bar{t}b\bar{b}$: $\mathcal{O}(15)$ pb
- $t\bar{t}W, t\bar{t}Z$: $\mathcal{O}(1.5)$ pb

ATLAS Coll. *Evidence for the associated production of the Higgs boson and a top quark pair with the ATLAS detector* [ttH multileptons](#) submitted for pub.

$t\bar{t}H$ multilepton

Choose a channel as example multilepton final state:

- Targeting $H \rightarrow WW, \tau\tau, ZZ$ decay modes, combined with leptonic $t\bar{t}$ decays
- Main background: $t\bar{t} \rightarrow$ Rely on signature with **same-sign (SS)** or three leptons
- Light lepton channel most sensitive to $H \rightarrow WW$
- τ_{had} channel more sensitive to $H \rightarrow \tau\tau$

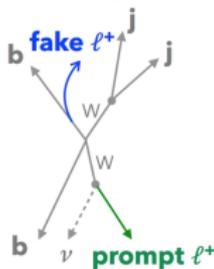


ttH multilepton: background estimation

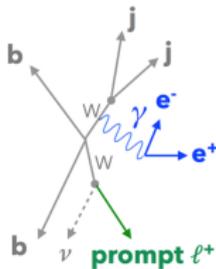
Main background sources:

- Processes with all prompt ℓ/τ_{had} \rightarrow Estimated with MC
- Electron charge mis-identification for 2lSS and 2lSS+1 τ_{had}
 - Estimated from data using charge mis-id rates measured in $Z \rightarrow e^+e^-$ and $Z \rightarrow e^\pm e^\pm$
- Events with fake/non-prompt light leptons:
 - Mainly from semileptonic b-hadron decays
 - Sizeable contribution also from photon conversions
- Events with fake tau leptons:
 - Mainly from light flavour jets and electrons mis-identified as τ_{had}

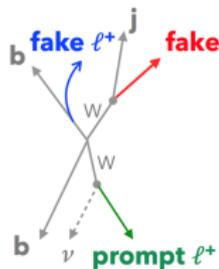
Semileptonic
b-decay



Photon
conversions

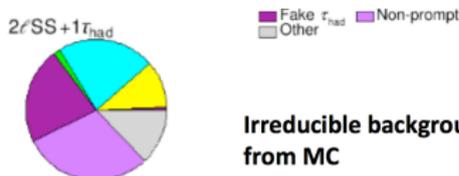
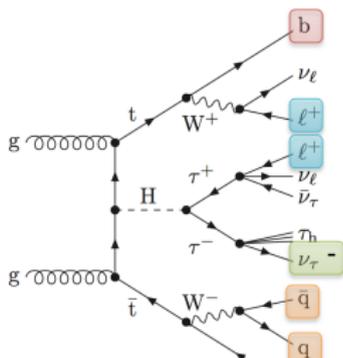


Non-prompt lepton
& fake τ



ttH multilepton: hadronic tau

Multilepton channel with a τ decaying hadronically:



Irreducible backgrounds from MC

- **ttW, ttZ, diboson**

Reducible backgrounds: → Data-driven estimates

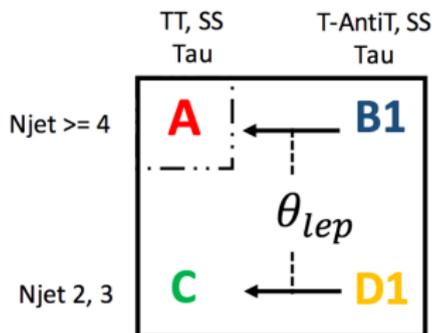
- **Non-prompt leptons**
- **Fake hadronic τ** from light/b-jets
- **Mis-identified charge Lepton**

- **Irreducible backgrounds: $t\bar{t}W$, $t\bar{t}Z$, other SM rare process**

The 2 same sign ℓ and 1 τ_{had} has both.

- Fake lepton majority from b-jets
- Fake taus majority from light additional jet
- ABCD method to determine data-driven background from non-prompt lepton
- data/MC fake tau rate derived from opposite sign τ_{had} control region
- the main problem is statistics in control region

ttH multilepton: hadronic tau, background



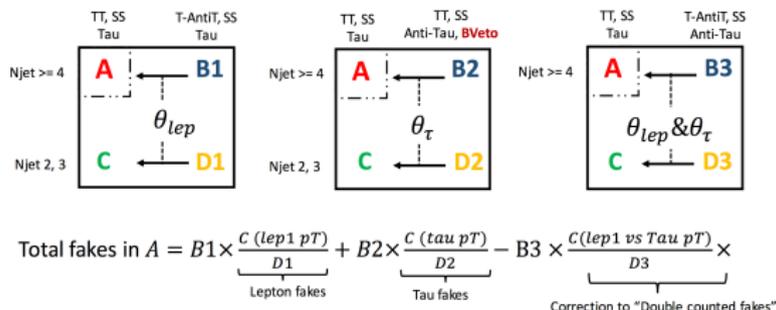
T= lepton selection *tight*,
 SS= same sign, Njet= numer
 of jets, θ_{lep} correction factor
A = signal region, **B1** = control
 region, **C** = control re-
 gion, **D1** = control region

$$\begin{aligned} \text{Total lepton fakes in A} \\ = B1 \times \frac{C(\text{lep1pT})}{D1} \end{aligned}$$

- A control region (CR)(D1) with an anti-tight selection is constructed by reversing the identification or isolation variables of the lepton, enriched in the fake leptons and it is orthogonal to the signal region (SR) A *Tight lepton* selection.

ttH multilepton: hadronic tau, background

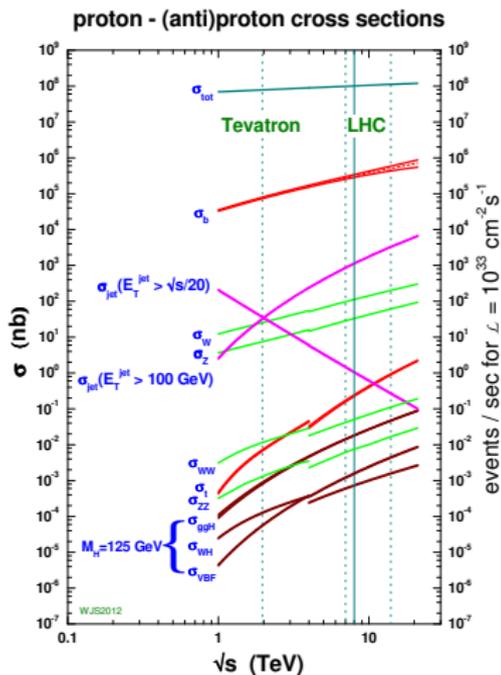
- The *Fake Factor CR* is constructed using low number of jets and the fake factor $\frac{C(\text{lep1}pT)}{D1}$ is calculated as a function of pT of the fake lepton.
- The fake factor is then applied to a CR B1 which is identical to the SR but with at least one anti-tight lepton.
- why we have not used as well for τ fakes ?



1 Identification of jets and leptons

2 Expected physics

Expectect physics: act 1



Example of calculation of expect events rate:

$$\text{Event rate: } \frac{dN}{dt} = \mathcal{L} \cdot \sigma$$

$$\sqrt{s} = 7 \text{ TeV}, \mathcal{L} = 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$$

$$\text{nb} = 10^{-33} \text{ cm}^{-2}$$

process	\sqrt{s} 8 TeV rate[1/s]	\sqrt{s} 14 TeV rate[1/s]
inel pp coll.	10^8	10^8
$b\bar{b}$	$3 \cdot 10^5$	$5 \cdot 10^5$
$t\bar{t}$	0.2	1
$W \rightarrow e\nu$	10	20
$Z \rightarrow ee$	1	2
Higgs	~ 0.02	~ 0.05

$$W \rightarrow e\nu \sim 10.71\%, Z \rightarrow ee \sim 3.36\%$$

Figure : from W.J. Stirling,
private communication

[http://www.hep.ph.ic.ac.uk/
~wstirlin/plots/plots.html](http://www.hep.ph.ic.ac.uk/~wstirlin/plots/plots.html)

Expectect physics: act 1

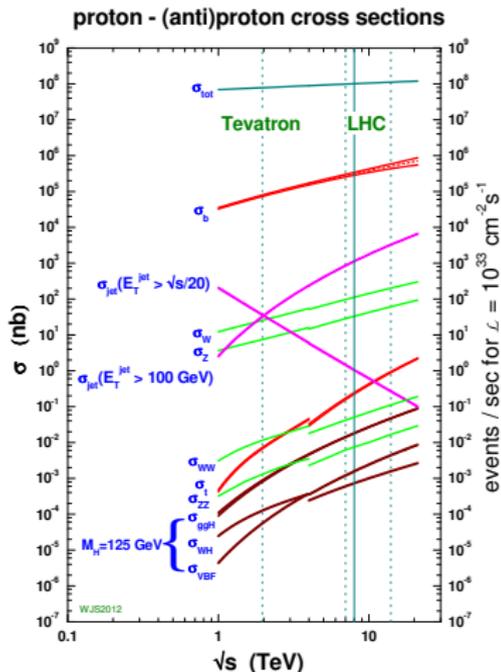


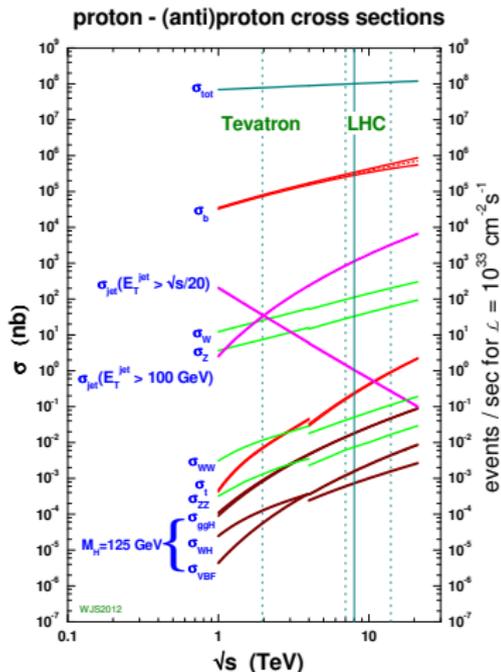
Figure : from W.J. Stirling,
private communication
[http://www.hep.ph.ic.ac.uk/
~wstirlin/plots/plots.html](http://www.hep.ph.ic.ac.uk/~wstirlin/plots/plots.html)

⇒ **Inelastic low- p_T collisions**

- Most processes are due to **soft** and semi-soft **interactions** between incoming protons.
- particle in final state have **large longitudinal**, but **small transverse momentum** \approx **several hundreds MeV**.

Low- p_T inelastic pp-collisions are Minimum Bias events Some parameters (multiplicity etc.) have to be tuned by MC simulation.

Expectect physics: act 2

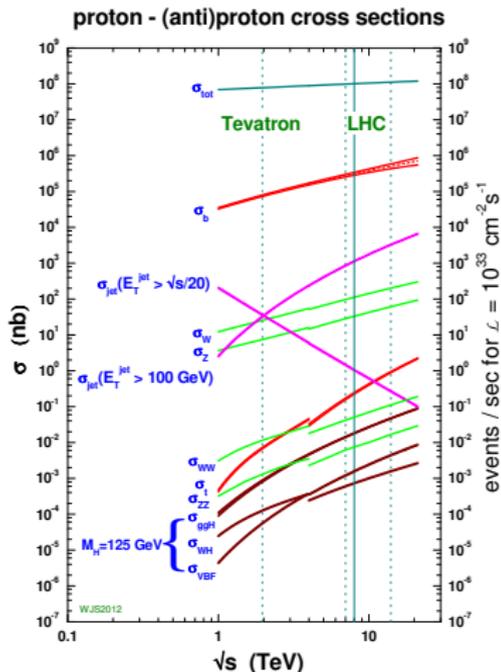


⇒ Measure Jet cross sections

- requires a good understanding of jets (jet energy scale, number of jets)
- a good calorimeter is necessary

Figure : from W.J. Stirling,
private communication
[http://www.hep.ph.ic.ac.uk/
~wstirlin/plots/plots.html](http://www.hep.ph.ic.ac.uk/~wstirlin/plots/plots.html)

Expectect physics: act 3



⇒ Electroweak sector

- test SM
- most SM cross sections are significantly higher than at the TEVATRON
 - *e.g* 100 x larger top-pair production cross sections
 - the LHC is a top,b, W, Z..... Higgs... factory.

Figure : from W.J. Stirling,
private communication
<http://www.hep.ph.ic.ac.uk/~wstirlin/plots/plots.html>

Important:

Concentrate on **final states with high- p_T** and **isolated leptons and photons** (+jets), otherwise overwhelmed by QCD jet background.

First top in Europe

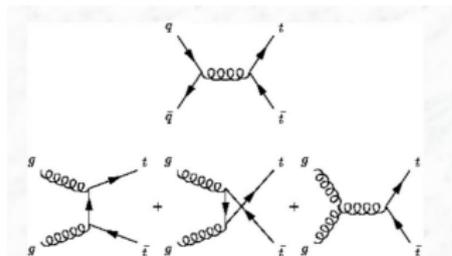
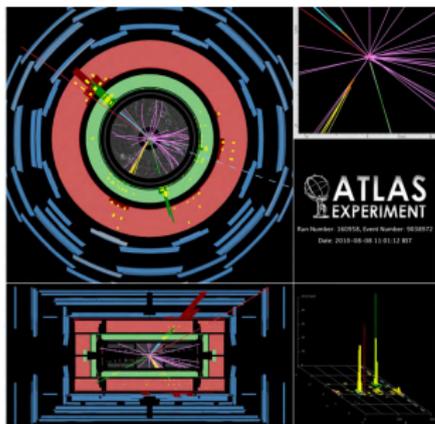


Figure : Diagrams

$tt \rightarrow Wb \ Wb \rightarrow e\nu b \ \mu\nu b$
The fragmentation products of b-quarks(B-hadrons) have a life times of 1.5 ps
= decay distance of ~ 2.5 mm.

$WW \rightarrow e\nu\nu$

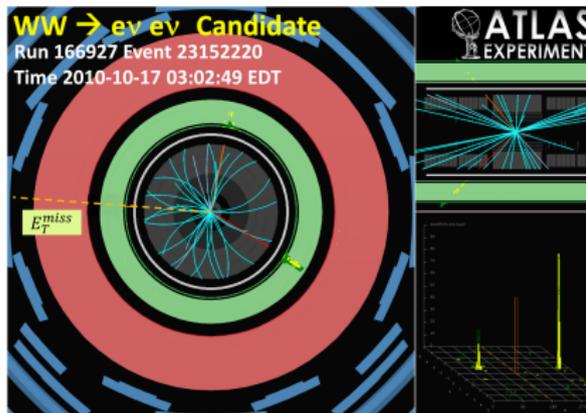


Figure : $WW \rightarrow ee\nu\nu$

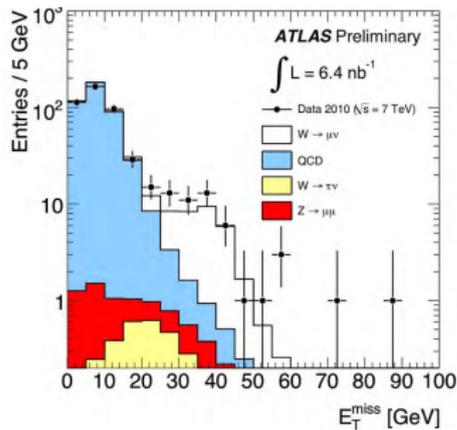


Figure : E_T^{miss}

- Benchmark processes defining main machine/detector parameters
- Basic processes relevant for study of SM
- All cross sections(\times Br) $\sim 1 - 100 fb^{-1}$ determines the statistics needed
- W mass
- top mass
- $B_s \rightarrow \mu\mu$ $B_s \rightarrow J/\psi$
- $pp \rightarrow W^+W^- \rightarrow \mu^+\nu_\mu\mu^-\bar{\nu}_\mu$
- $pp \rightarrow H \rightarrow ZZ \rightarrow \mu^+\mu^-\mu^+\mu^-$
- $pp \rightarrow H \rightarrow \gamma\gamma$
- $pp \rightarrow Z' \rightarrow \mu^+\mu^-$
- SUSY with R-parity, SUSY Higgs

The things appeared with time, accumulating statistics:

- 1 MinBias/ low- p_T Physics
- 2 jets
- 3 W/Z
- 4 Top
- 5 Search Higgs & new Physics

The Standard Model at LHC

Standard Model Total Production Cross Section Measurements

Status: July 2017

