

# Short introduction

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In short which is the main purpose of the EEPP and few numbers that every experimental particle physicist should have in his/her hands.

# Introduction

- The “Question to Nature” in EPP: it is the quest for the “fundamental” aspects of the Nature: not single phenomena but the common grounds of all physics phenomena.
- Historical directions of the EPP:
  - Atomic physics → Nuclear Physics → Subnuclear Physics: the only small; Nature = point-like particles interacting through forces..
  - Look at the only large: connections with cosmology, cosmic rays, etc..
  - Paradigm: unification of forces, theory of everything.
- What shall we do in this course ?
  - We concentrate on subnuclear physics and will select few experiments
  - We review some “basic statistics” and then will extend it to more “advanced” methods for data analysis EPP experiments

# The EPP experiment

- Something present through all the 20<sup>o</sup> century and continuing in 21<sup>o</sup> : the best way to understand the elementary particles and how do they interact, is to send *projectiles* on *targets*, or, more generally, “to make things collide”. And look at the *final state*:  $a+b \rightarrow X$
- “Mother-experiment” (Rutherford): 3 main elements:
  - a projectile
  - a target
  - a detector
- Main rule: the higher the momentum  $p$  of the projectile, the smaller the size  $\delta x$  I am able to resolve.

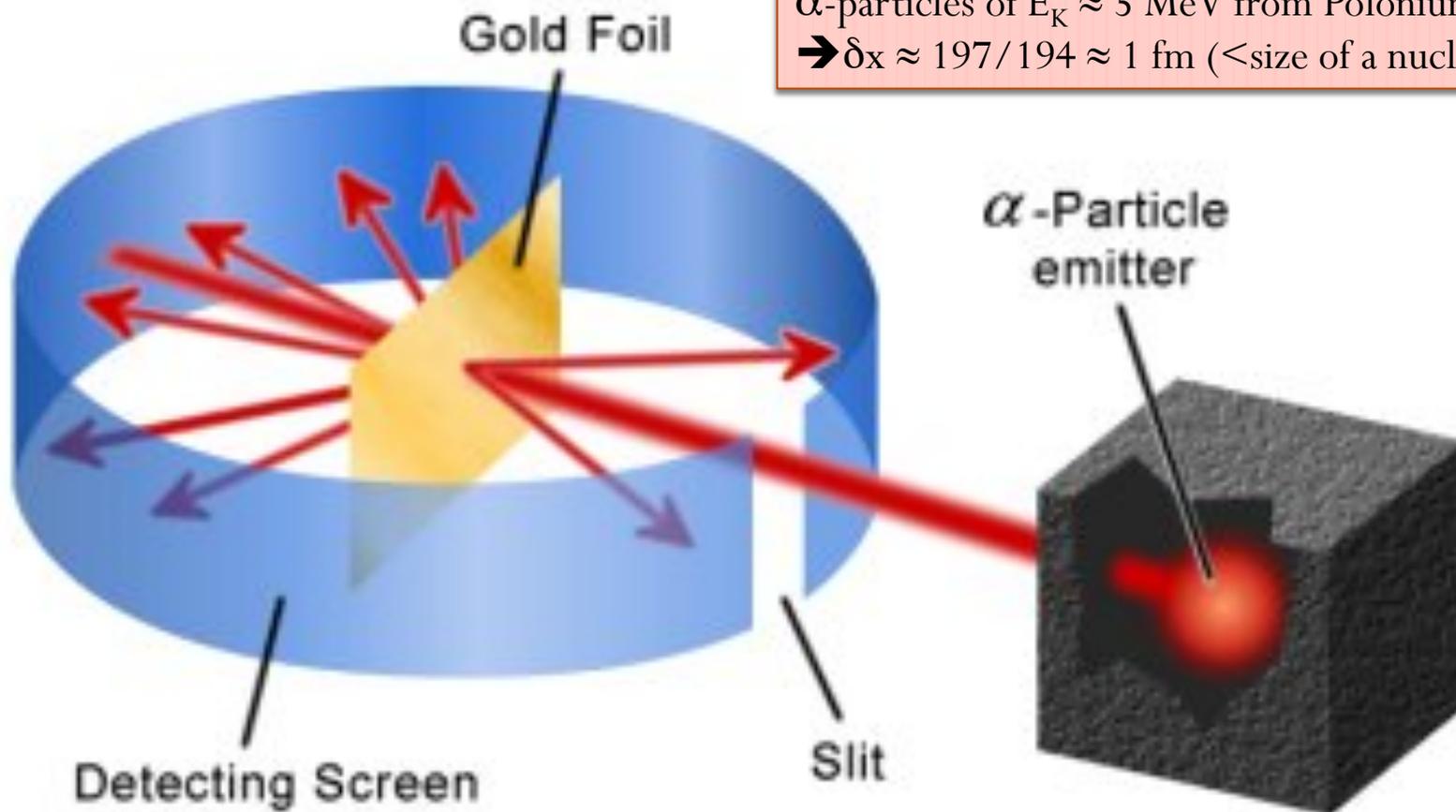
$$\delta x \approx \frac{\hbar c}{pc} \Rightarrow \delta x(fm) \approx \frac{197}{p(MeV/c)}$$

The scale:  $\hbar c = 197 MeV \times fm$

- From Rutherford, a major line of approach to nuclear and nucleon structure using electrons as projectiles and different nuclei as targets.

# The Rutherford experiment

$\alpha$ -particles of  $E_K \approx 5 \text{ MeV}$  from Polonium  
 $\rightarrow \delta x \approx 197/194 \approx 1 \text{ fm}$  ( $<$ size of a nucleus)

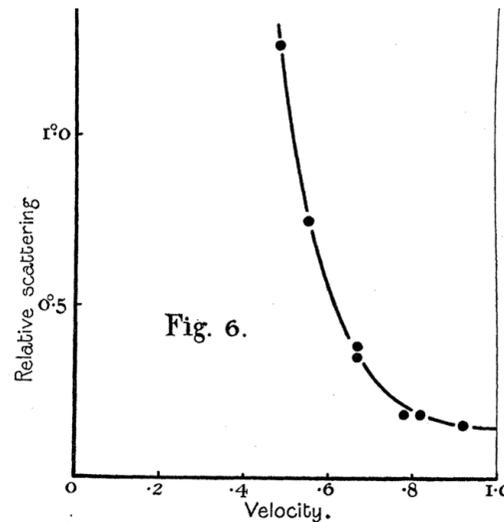
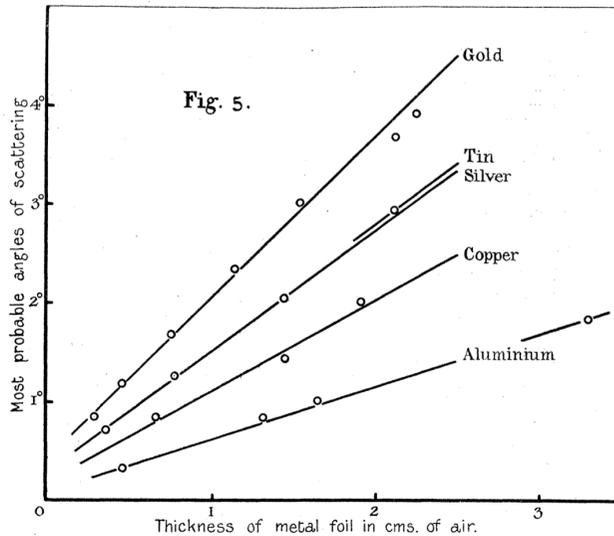
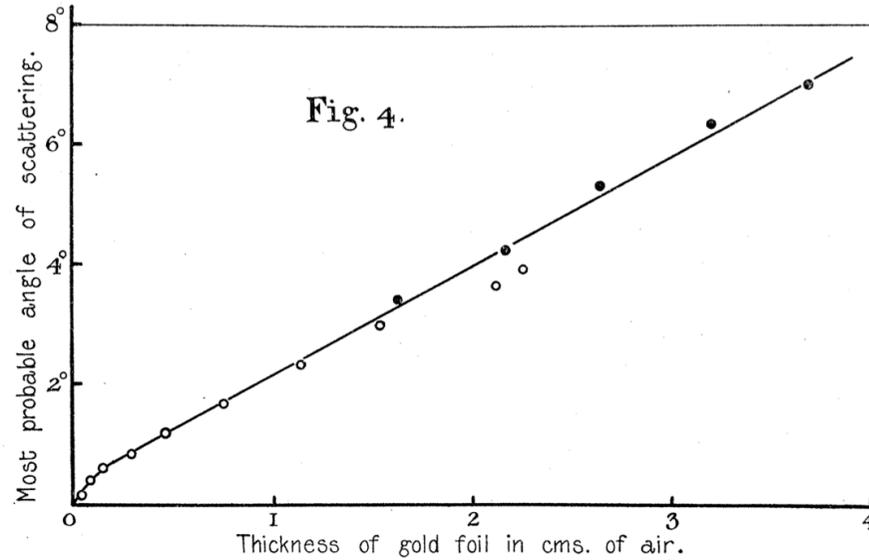
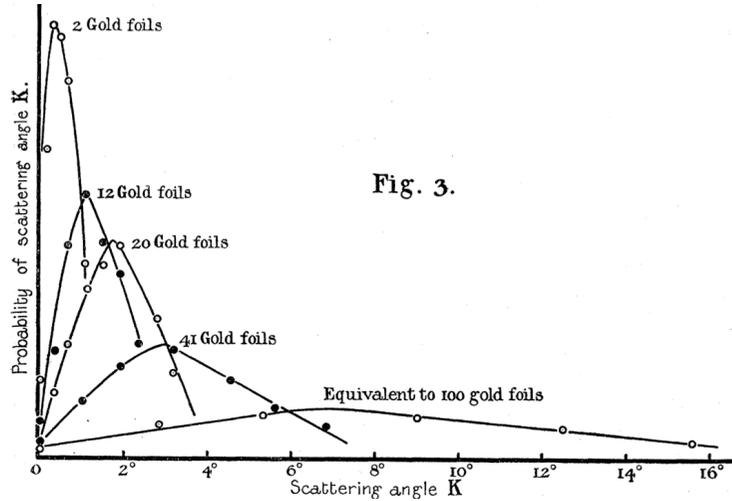


$$p^2 = (m_\alpha + E_K)^2 - m_\alpha^2 = 194 \text{ MeV}$$

# Key elements in the Rutherford experiment – physical quantities

- **Energy of the collision** (driven by the kinetic energy of the  $\alpha$  particles) the meaning of  $\sqrt{s}$
- **Beam Intensity** (how many  $\alpha$  particles /s)
- **Size and density of the target** (how many gold nuclei encountered by the  $\alpha$  particles);
- **Deflection angle  $\theta$**
- **Probability/frequency of a given final state** (fraction of  $\alpha$  particles scattered at an angle  $\theta$ );
- **Detector efficiency** (do I see all scattered  $\alpha$  particles ?)
- **Detector resolution** (how well do I measure  $\theta$  ?)

# The Rutherford experiment – original results



Plots from the original Geiger paper of 1910  
 → MS formula coming out from data:  $\theta \approx Z \delta X / v$   
 NB: no mention of measurement uncertainties..

# Break: the Rutherford experiment only ?

- Actually more than the Rutherford experiment
- Particle Physics without beams
  - → cosmic ray based experiments
    - In space
    - In Underground Laboratories
    - In DeepSea Detectors
  - → Search for very rare or forbidden decays of ordinary matter
    - Mostly in underground detectors
- Examples during the course
- NOW: let's concentrate on EPP with beams

# Energy: what is $\sqrt{s}$ ?

- This is a fundamental quantity to define the “effective energy scale” you are probing your system. It is how much energy is available for each collision in your experiment.
- It is relativistically invariant.
- If the collision is  $a+b \rightarrow X$

$$\begin{aligned} s &= (\tilde{p}_a + \tilde{p}_b)^2 = M_a^2 + M_b^2 + 2\tilde{p}_a \cdot \tilde{p}_b \\ &= M_a^2 + M_b^2 + 2[E_a E_b - \vec{p}_a \cdot \vec{p}_b] \end{aligned}$$

- $M_X$  cannot exceed  $\sqrt{s}$ .
- What about Rutherford experiment ?  $a=\alpha$ ,  $b=\text{Au}$ ,  $X=a+b$

$$\begin{aligned} s &= M_\alpha^2 + M_{\text{Au}}^2 + 2E_\alpha M_{\text{Au}} = \\ \sqrt{s} &= 188.5 \text{ GeV} \end{aligned}$$

Maybe Rutherford produced a Higgs ??

# Development along the years

- **WARNING:** Not only Rutherford: in the meantime EPP developed several other lines of approaches.
- More was found: It was seen that going up with the projectile momentum something unexpected happened: more particles and also new kinds of particles were “**created**”.
- **→** high energy collisions allow to create and study a sort of “**Super-World**”. The properties and the spectrum of these new particles can be compared to the theory of fundamental interactions (the Standard Model).
- Relation between projectile momentum and “creation” capability:
- **→** Colliding beams are more effective in this “creation” program.
  - ep colliders (like HERA)
  - $e^+e^-$  storage rings
  - p-pbar or pp colliders

$$\sqrt{s} = \sqrt{M_1^2 + M_2^2 + 2E_1M_2} \approx \sqrt{2E_1M_2}$$

$$\sqrt{s} = 2\sqrt{E_1E_2}$$

# Units - I

- $\Delta E_k = q\Delta V$
- Joule “=”  $C \times V$  in MKS
- Suppose we have an electron  $q = e = 1.602 \times 10^{-19} \text{ C}$  and a  $\Delta V = 1 \text{ V}$ :  $\rightarrow \Delta E_k = 1.6 \times 10^{-19} \text{ J} = 1 \text{ eV}$
- Particularly useful for a linear accelerator
  - Electrons are generated through cathodes by thermoionic effect;
  - Protons and ions are generated through ionization of atoms;
  - Role of “electric field”: how many  $V/m$  can be provided ?
  - Present limit  $\approx 30 \div 50 \text{ MV/m}$  (100  $\text{MV/m}$  CLIC)
    - $\rightarrow 1 \text{ km}$  for  $30 \div 50 \text{ GeV}$  electrons !

# Units - II

- Unit system
  - By posing  $c = 1$ , **energy**, **momentum** and **mass** can all be expressed in terms of a single fundamental unit. All can be expressed using the eV.

$$E^2 = (pc)^2 + (mc^2)^2 \rightarrow E^2 = p^2 + m^2$$

- $c=1$  implies also the following dimensional equation:
  - $[L] = [T]$   
Lengths and times have the same units
- Then we also pose  $\hbar=1$ , this has implications on energy vs. l and t
  - $[E] = [L]^{-1} = [T]^{-1}$   
→ time and length are (energy)<sup>-1</sup>
- Numerically we need few conversion factors:
  - $1 \text{ MeV} == 0.00506 \text{ fm}^{-1} == 1.519 \text{ ns}^{-1}$

# Energy scales

- In the following we try to see which scales of energy correspond to different phenomenologies. We consider equivalently space and energy scales (since we know it is somehow the same..)
- This quantity is one of the driving element to design HEP experiments: you need to know first of all at which energy you have to go.

# Energy scales in the $\infty$ ly small - I

- Electromagnetic interactions have not a length scale

$$V = \frac{1}{4\pi\epsilon_0} \frac{e^2}{r}$$

- $[V \times r] = [E][L] = [\hbar c] \rightarrow$  we can define an adimensional quantity  $\alpha$ :

$$\frac{e^2}{4\pi\epsilon_0 \hbar c} = \alpha = \frac{(1.610^{-19} \text{ C})^2}{4\pi 8.8510^{-19} \text{ F/m} 1.0510^{-34} \text{ Js} 310^8 \text{ m/s}} = \frac{1}{137} = 0.0073$$

- $\alpha$  sets the scale of the *intensity* of the electromagnetic interactions. In natural units ( $\hbar = c = \epsilon_0 = \mu_0 = 1$ )  $e$  is also adimensional:

$$e = \sqrt{4\pi\alpha}$$

# Energy scales in the $\infty$ ly small - II

- Electromagnetic scales:
  - **1. Classical electron radius:** The distance  $r$  of two equal test charges  $e$  such that the electrostatic energy is equal to the rest mass  $mc^2$  of the charges

$$r_e = \frac{e^2}{4\pi\epsilon_0 m_e c^2} = \frac{\alpha \hbar}{m_e c} \rightarrow \frac{\alpha}{m} \quad \text{In natural units}$$

- **Electron Compton wavelength:** which wavelength has a photon whose energy is equal to the electron rest mass.

$$\tilde{\lambda}_e = \frac{\hbar}{m_e c} = \frac{r_e}{\alpha} \rightarrow \frac{1}{m_e}$$

- **Bohr radius:** radius of the hydrogen atom orbit

$$a_\infty = \frac{4\pi\epsilon_0 \hbar^2}{m_e e^2} = \frac{r_e}{\alpha^2} \rightarrow \frac{1}{\alpha m_e}$$

# Energy scales in the $\infty$ ly small - III

- Weak interactions: Fermi theory introduces the constant  $G_F$  with dimensions  $[E]^{-2}$  (making the theory non-renormalizable). In the electroweak theory  $G_F$  is:

$$\frac{G_F}{\sqrt{2}} = \frac{g_W^2}{8m_W^2}$$

- Where  $g_W$  is the “fundamental” adimensional coupling directly related to  $e$  through the Weinberg angle:  $e = g_W \sin \theta_W$
- The “Electroweak scale” is the scale at which the electroweak unification is at work. By convention it is given by  $v$ , the Higgs vacuum expectation value:

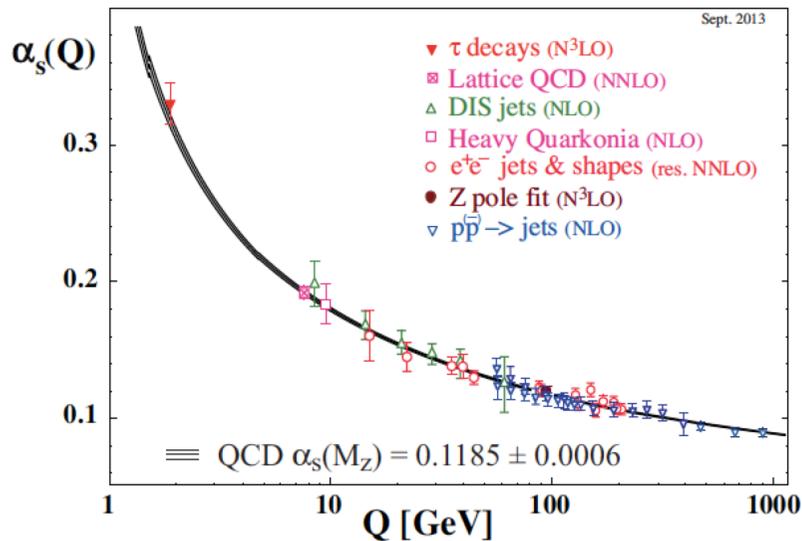
$$v = \frac{1}{\sqrt{\sqrt{2}G_F}} = 246 \text{ GeV} \quad r_{EW} \approx \sqrt{\sqrt{2}G_F} (\hbar c)$$

# Energy scales in the $\infty$ ly small - IV

Strong interaction: Yukawa potential

$$V(r) = \frac{g^2}{4\pi r} \exp\left(-\frac{r}{\lambda}\right)$$

$\lambda$  is  $1/m(\text{pion})$



- Strong Interaction scale:  $\alpha_s$  depends on  $q^2$ . There is a natural scale given by the “confinement” scale, below which QCD predictions are not reliable anymore.

$$r_{QCD} = \frac{1}{\Lambda_{QCD}} \approx \langle r_{proton} \rangle$$

# Energy scales in the ∞ly small - V

- Gravitational Interaction scale: the “problem” of the gravity is that the coupling constant is not adimensional, to make it adimensional you have to multiply by  $m^2$ . The adimensional quantity here is

$$\frac{Gm^2}{\hbar c}$$

depending on the mass. For typical particle masses it is  $\ll 1$ . The mass for which it is equal to 1 is the “Planck Mass”  $M_{Planck}$ .  $\lambda_{Planck}$  is the “Planck scale” (Compton wavelength of a mass  $M_{Planck}$ )

$$M_{Planck} = \sqrt{\frac{\hbar c}{G}} \quad \lambda_{Planck} = \sqrt{\frac{\hbar G}{c^3}}$$

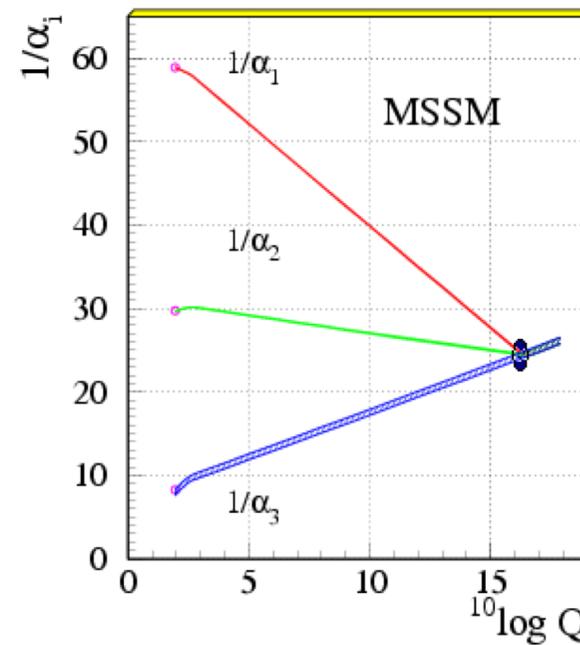
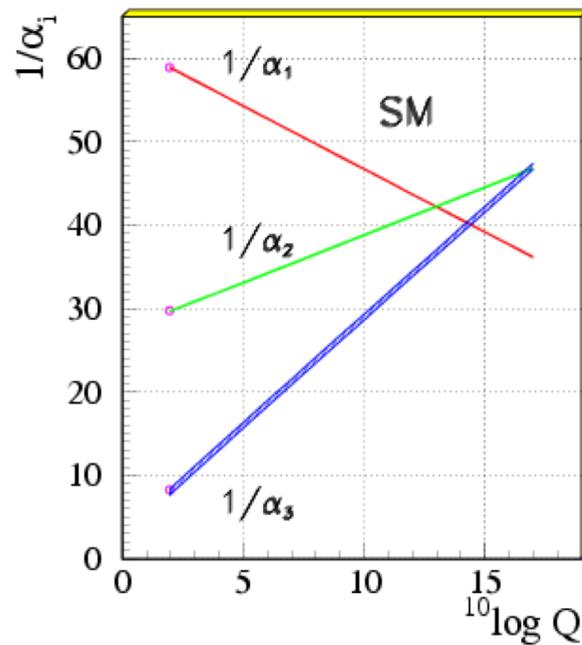
$M_{planck}$  is  $\approx 20 \mu\text{g}$ , a “macroscopic” quantity.

# The Planck scale

- When you increase a mass
  - $\rightarrow$  you are reducing its Compton wavelength (that is the scale at which quantum effects are relevant)
  - $\rightarrow$  you increase the Schwarzschild radius (that is the radius of the event horizon of the black hole with that mass)
- The mass for which Compton wavelength = Schwarzschild radius is the Planck Mass  $\rightarrow$  is the domain of the quantum gravity

# Energy scales in the $\infty$ ly small - VI

- Grand Unification Scale. From the observation that weak, em and strong coupling constants are “running” coupling constants, if we plot them vs.  $q^2$  we get:



Around  $10^{16}$  GeV meeting point ??

# Energy scales in the ∞ly small - Summary

quantity	value	Energy
Bohr radius	$0.53 \times 10^{-10} \text{ m}$ (0.5 Å)	3.7 keV
Electron Compton wavelength	$3.86 \times 10^{-13} \text{ m}$ (386 fm)	0.51 MeV
Electron classical radius	$2.82 \times 10^{-15} \text{ m}$ (2.8 fm)	70 MeV
Proton radius – QCD confinement scale	$0.82 \times 10^{-15} \text{ m}$ (0.8 fm)	240 MeV
Fermi scale	$7 \times 10^{-19} \text{ m}$	250 GeV
“New Physics” scale		1 TeV
GUT Scale		$10^{16} \text{ GeV}$
Planck scale	$1.62 \times 10^{-35} \text{ m}$	$1.2 \times 10^{19} \text{ GeV}$

The TeV scale is the maximum reachable with the present accelerator technology

# Energy scales in the ∞ly small - VII

- Why LHC is concentrate on the O(TeV) scale ?
- There is an intermediate scale around the TeV. It is motivated by the “naturalness” – “fine tuning” – “hierarchy” problem connected to the properties of the Higgs Mass.

$$m_H^2 \sim -2\mu^2 + \frac{g^2}{(4\pi)^2} M^2$$

Mass parameter in the SM lagrangian

Quantum corrections

- The Higgs mass  $m_H$  is UV sensitive (its value depends on quantum corrections)
- $M$  is the scale up to which we have the UV theory.
- If no other scale is there btw Higgs and Planck,  $M=M_{Planck}$ , so that strong cancellations are needed between  $-2\mu^2$  and  $g^2 M^2 / (4\pi)^2$  to give the observed Higgs Mass
- This is un-natural..
- If  $M \approx O(\text{TeV})$  all becomes natural, e.g. MSSM, Technicolor, ...

$$\Delta \gtrsim \left( \frac{m_{NP}}{0.5 \text{ TeV}} \right)^2$$

# More in detail

- $m_H$  is the Higgs mass;  $\mu$  is the Higgs “bare” mass (the parameter in the lagrangian).  $m_H = \mu + \text{“RC”}$  (radiative corrections due to fermion and boson loops). If “RC”  $\gg m_H$  it means that also  $|\mu| \gg m_H$  and a cancellation btw RC and  $\mu$  is needed.
- Structure of “RC”. For every particle  $p$  in the loop it is  $= g_p^2 (\Lambda^2 + m_p^2)$ .  $\Lambda$  is the “cut-off” of the integration, it is the next scale that nature gives to us.
- Supersymmetric solution. In “RC”  $N$  particle-antiparticle pairs with opposite sign couplings enter  $= N_p g_p^2 (\Lambda^2 + m_p^2) - N_{\text{antip}} g_{\text{antip}}^2 (\Lambda^2 + m_{\text{antip}}^2) = N_p g_p^2 (m_p^2 - m_{\text{antip}}^2)$ ;  $\Lambda$  is cancelled

# Probability/Frequency of a final state: the cross-section and the decay width

- The **cross-section** measures the “probability” of a given final state in a collision (actual definition will be in a later lecture). It is a  $[L]^2$ .
- The **decay width** and the **branching ratio** measure the “probability” of a given final state in a deca. The decay width is the inverse of the lifetime so that it is a  $[T]^{-1}$ . The branching ratio is an adimensional quantity
- If we include **cross-sections** and **decay widths**, we enter in the quantum field theories where the normalized Planck constant enters in the game.
- In the “natural system” the units are  $\hbar = c = 1$ 
  - **cross-section** is a  $(\text{length})^2$  so an  $(\text{energy})^{-2}$ .
  - **decay width** is a  $(\text{time})^{-1}$  so an  $(\text{energy})$
  - $1 \text{ GeV}^{-2} = 3.88 \times 10^{-4} \text{ barn}$

# Cross-section scales

- Relation between an experimental cross-section and the theory (same applies for branching ratios)

$$\sigma = \int \left| \text{Feynman Diagrams} \right|^2 d\phi$$

The equation shows the cross-section  $\sigma$  as an integral over phase space  $d\phi$  of the squared magnitude of the sum of two Feynman diagrams. The first diagram is a t-channel exchange of a photon ( $\gamma$ ) between two electron-positron pairs. The second diagram is an s-channel exchange of a photon ( $\gamma$ ) between two electron-positron pairs. Each external line is labeled with 'e' and has an arrow indicating the direction of particle flow.

Two ingredients in the theory calculations:

→ dynamics (amplitude from lagrangian, Feynman diagrams... mainly the coupling constants);

→ phase space  $d\phi$

NB: the integration on the phase space **DEPENDS** in general

on the experiment details (accessible kinematic region) → **Montecarlo**

# Cross-section order of magnitude estimates

- Based on dimensional arguments and few numbers (neglects phase-space and more...)
  - Electromagnetic processes:  $e^+e^- \rightarrow \mu^+\mu^-, \gamma\gamma$
  - Weak processes:  $\nu N$  scattering
  - Hadron strong interaction scattering:  $pp$  scattering

$\alpha$	1/137
$G_F$	$10^{-5} \text{ GeV}^{-2}$
$r_p$	1 fm
$1 \text{ GeV}^{-2}$	$3.88 \times 10^{-4} \text{ b}$

$$\sigma(e^+e^- \rightarrow \mu^+\mu^-, \gamma\gamma) \approx \frac{\alpha^2}{s}$$

$$\sigma(\nu e \rightarrow \nu e) \approx G_F^2 2m_e E_\nu$$

$$\sigma(pp) \approx \pi r_p^2$$

$S=(1 \text{ GeV})^2$	$S=(100 \text{ GeV})^2$
20 nb	2 pb
40 fb	4 pb
30 mb	30 mb

# LifeTime (or Width) of a particle vs. theory

- As for the cross-section the value depends on two ingredients:
  - Decay type (weak, em, strong) through decay matrix element
  - Volume of the available phase space
- The Width  $\Gamma$  is an additive quantity: you have to add the *partial widths* of the single decays to get the *total width*
- Useful formulas: two-body decay phase-space (rest system)

$$\Gamma = \frac{1}{8\pi} \frac{p}{M^2} |\mathfrak{M}|^2. \quad \text{NB Dimensions: If } \Gamma \text{ is [E]} \rightarrow |M| \text{ is also [E]}$$

$$|\vec{p}_1| = |\vec{p}_2| = \frac{[(M^2 - (m_1 + m_2)^2)(M^2 - (m_1 - m_2)^2)]^{1/2}}{2M},$$

# Width (LifeTime) order of magnitude estimates

- The amplitude square has the dimensions of  $E^2$ .
  - Weak  $\rightarrow |Ampl|^2 \approx G_F^2 \times M^6$
  - E.m.  $\rightarrow |Ampl|^2 \approx \alpha^2 \times M^2$
  - Strong  $\rightarrow |Ampl|^2 \approx \alpha_s(M)^2 \times M^2$
- Examples of estimates (wrong by factor  $\approx 10$  maximum):

Interaction	Decay	Phase space ( $\text{MeV}^{-1}$ )	$ Ampl ^2$ ( $\text{MeV}^2$ )	$\Gamma$ ( $\text{MeV}$ )	$\tau$ (s)
Weak	$\pi^\pm \rightarrow \mu^\pm \nu$	$6.0 \times 10^{-5}$	$6.0 \times 10^{-10}$	$3.6 \times 10^{-14}$	$1.8 \times 10^{-8}$ <b><math>(2.6 \times 10^{-8})</math></b>
e.m.	$\pi^0 \rightarrow \gamma\gamma$	$1.5 \times 10^{-4}$	0.97	$1.4 \times 10^{-4}$	$4.6 \times 10^{-18}$ <b><math>(8.5 \times 10^{-17})</math></b>
strong	$\rho^0 \rightarrow \pi^+\pi^-$	$2.4 \times 10^{-5}$	$6.0 \times 10^5$	13 <b>(150)</b>	$5.0 \times 10^{-23}$

	Lifetime $\tau$	Width $\Gamma$
<b>Weak decays</b>		
$K_s, K_L$	$0.89564 \times 10^{-10} \text{ s}, 5.116 \times 10^{-8} \text{ s}$	
$K^\pm$	$1.2380 \times 10^{-8} \text{ s}$	
$\Lambda$	$2.632 \times 10^{-10} \text{ s}$	
B-hadrons	$\approx 10^{-12} \text{ s}$	
Muon	$2.2 \times 10^{-6} \text{ s}$	
Tau-lepton	$2.9 \times 10^{-13} \text{ s}$	
Top-quark	$\approx 5 \times 10^{-25} \text{ s}$	2 GeV
<b>e.m. decays</b>		
$\pi^0$	$8.52 \times 10^{-17} \text{ s}$	8 eV
$\eta$	$\approx 10^{-19} \text{ s}$	1.30 keV
<b>Strong decays</b>		
$J/\psi$		92.9 keV
$\Upsilon$		54.02 keV
$\rho$		149.1 MeV
$\omega$		8.49 MeV
$\phi$		4.26 MeV
$\Delta$		114 ÷ 120 MeV

# Recap - fundamental interactions

- Electromagnetic interaction:
  - Can be studied at all energies with “moderate” cross-sections;
  - Above  $O(100 \text{ GeV})$  becomes electro-weak
- Weak interactions:
  - At low energies it can be studied using decays of “stable” particles – large lifetimes and small cross-sections;
  - Above  $O(100 \text{ GeV})$  becomes electro-weak
- Strong interactions:
  - At low energy (below  $1 \text{ GeV}$ ) “hadronic physics” based on confinement: no fundamental theory available by now
  - At high energies (above  $1 \text{ GeV}$ ) QCD is a good theory: however since partons are not directly accessible, only “inclusive” quantities can be measured and compared to theory. Importance of simulations to relate partonic quantities to observables.

# Comparison between beam possibilities

- Electrons:
  - Clean, point-like, fixed (almost) energy, but large irradiation due to the low mass. “Exclusive” studies are possible (all final state particles are reconstructed and a complete kinematic analysis can be done)
  - →  $e^+e^-$  colliders less for energy frontier, mostly for precision measurements
- Protons:
  - Bunch of partons with momentum spectrum, but low irradiation. “Inclusive” studies are possible. A complete kinematic analysis is in general not possible (only in the transverse plane it is to first approximation possible)
  - → highest energies are “easily” reachable, high luminosity are reachable but problems in the interpretation of the results; very “demanding” detectors and trigger systems.
- Anti-protons:
  - Difficult to obtain high intensities and high luminosity but no problems with energies, same problems of protons (bunch of partons)
- → p-antip limited by luminosity,  $e^+e^-$  limited by energy BUT perfect for precision studies, pp good choice for energy frontier

# Implications for experiments:

- You need high energy for
  - Probe electro-weak scales, get closer to higher scales
  - Enlarge the achievable mass spectrum (particle discoveries)
- You need high beam intensity and large/dense targets or high efficiency detectors
  - To access low probability phenomena
- You need high resolution detectors
  - To improve particle discrimination especially for rare events.

# Where do we stand now.

- The EW + QCD Standard Model allows to describe reasonably well most of the “high energy”  $> O(10 \text{ GeV})$  phenomena
- However:
  - The model is unsatisfactory under several points of view
    - Hierarchy / naturalness problem
    - Large number of unpredictable parameters
  - Left behind “ununderstood areas”
    - Strong interaction phenomena below  $O(1 \text{ GeV})$
    - Hadron spectroscopy
    - No description / no space left for dark matter
    - Still not clear picture of neutrino dynamics
    - Of course gravitation is out...

# End of the Introduction

- Present prospects of Elementary Particle experiments:
  - ENERGY frontier → LHC, HL-LHC, ILC, TLEP,....
  - INTENSITY frontier → flavour-factories, fixed target,...
  - SENSITIVITY frontier → detectors for dark matter, neutrinos,..
- The general idea is to measure quantities for which you have a clear prediction from the Standard Model, and a hint that a sizeable correction would be present in case of “New Physics”.