

Consideriamo due particelle di massa m non interagenti in una scatola posizionate fra 0 ed L .

Una particella si trova nello stato $|n\rangle$ della scatola, l'altra nello stato $|l\rangle$ (supponiamo $l \neq n$).

Calcolare

$$\langle (x_1 - x_2)^2 \rangle$$

nel caso in cui le particelle siano

1/ distinguibili

2/ due fermioni

3/ due bosoni.

Sapendo che, per una particella nella scatola, si ha $\langle x \rangle = \langle n | x | n \rangle = \frac{L}{2}$ e $\langle x^2 \rangle = \langle n | x^2 | n \rangle = L^2 \left[\frac{1}{3} - \frac{1}{2(n\pi)^2} \right]$

Lo stato del sistema si scrive

$|l n\rangle$ part. distinguibili

$$\frac{1}{\sqrt{2}} (|l n\rangle + |n l\rangle) \text{ bosoni}$$

$$\frac{1}{\sqrt{2}} (|l n\rangle - |n l\rangle) \text{ fermioni}$$

In rappresentazione delle coordinate, le funzioni d'onda sono quindi

$$\Psi_D(x_1, x_2) = \Psi_e(x_1) \Psi_n(x_2)$$

$$\Psi_S(x_1, x_2) = \frac{1}{\sqrt{2}} \langle x_1 x_2, | \left(\frac{1}{\sqrt{2}} (|l n\rangle + |n l\rangle) \right) \rangle$$

$$= \frac{1}{2\sqrt{2}} \left[\langle x_1 x_2 | + \langle x_2 x_1 | \right] \left[|l n\rangle + |n l\rangle \right]$$

$$= \frac{1}{2\sqrt{2}} \left[\langle x_1 x_2 | l n \rangle + \langle x_1 x_2 | n l \rangle + \langle x_2 x_1 | l n \rangle + \langle x_2 x_1 | n l \rangle \right]$$

$$= \frac{1}{2\sqrt{2}} \left[\Psi_e(x_1) \Psi_n(x_2) + \Psi_n(x_1) \Psi_e(x_2) + \Psi_e(x_2) \Psi_n(x_1) + \Psi_n(x_2) \Psi_e(x_1) \right]$$

$$= \frac{1}{\sqrt{2}} \left[\Psi_e(x_1) \Psi_n(x_2) + \Psi_n(x_2) \Psi_e(x_1) \right]$$

$$e \quad \Psi_A(x_1, x_2) = \frac{1}{\sqrt{2}} \left[\Psi_e(x_1) \Psi_n(x_2) - \Psi_n(x_1) \Psi_e(x_2) \right]$$

Ora

$$\langle (x_1 - x_2)^2 \rangle = \langle x_1^2 \rangle + \langle x_2^2 \rangle - 2 \langle x_1 x_2 \rangle$$

1/ Distinguishability

$$\langle x_1^2 \rangle = \int dx_1 dx_2 x_1^2 |\Psi_D(x_1, x_2)|^2$$

$$= \int dx_1 x_1^2 |\Psi_e(x_1)|^2 \int dx_2 |\Psi_n(x_2)|^2$$

$$= \int dx_1 x_1^2 |\Psi_e(x_1)|^2 = \langle x^2 \rangle_e$$

$$\langle x_2^2 \rangle = \langle x^2 \rangle_n$$

$$\langle x_1 x_2 \rangle = \int dx_1 x_1 |\Psi_e(x_1)|^2 \int dx_2 x_2 |\Psi_n(x_2)|^2$$

$$= \langle x \rangle_e \langle x \rangle_n$$

$$\rightarrow \langle (x_1 - x_2)^2 \rangle_D = L^2 \left[\frac{1}{3} - \frac{1}{2\pi^2 \ell^2} \right] + L^2 \left[\frac{1}{3} - \frac{1}{2\pi^2 n^2} \right]$$

$$- 2 \frac{L^2}{4}$$

$$= \left[\frac{2}{3} - \frac{1}{2} \right] L^2 - \frac{L^2}{2\pi^2} \left[\frac{1}{\ell^2} + \frac{1}{n^2} \right]$$

$$= \frac{4-3}{6} L^2 - \frac{L^2}{2\pi^2} \left(\frac{1}{\ell^2} + \frac{1}{n^2} \right) = L^2 \left[\frac{1}{6} - \frac{1}{2\pi^2} \left(\frac{1}{\ell^2} + \frac{1}{n^2} \right) \right]$$

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Partielle identische

$$\langle x_1^2 \rangle_{S,A} = \frac{1}{2} \int dx_1 dx_2 [\psi_e^*(x_1) \psi_n^*(x_2) + \psi_n^*(x_1) \psi_e^*(x_2)]$$

$$x_1^2 [\psi_e(x_1) \psi_n(x_2) + \psi_n(x_1) \psi_e(x_2)]$$

$$\begin{aligned} &= \frac{1}{2} \left\{ \int dx_1 x_1^2 |\psi_e(x_1)|^2 \int dx_2 |\psi_n(x_2)|^2 \right. \\ &\quad + \int dx_1 x_1^2 \psi_e^*(x_1) \psi_n(x_1) \underbrace{\int dx_2 \psi_n^*(x_2) \psi_e(x_2)}_0 \\ &\quad + \int dx_2 x_2^2 \psi_e^*(x_2) \psi_n(x_2) \underbrace{\int dx_1 \psi_e^*(x_1) \psi_n(x_1)}_0 \\ &\quad \left. + \int dx_1 x_1^2 |\psi_n(x_1)|^2 \int dx_2 |\psi_e(x_2)|^2 \right\} \\ &= \frac{1}{2} \left\{ \langle x^2 \rangle_e + \langle x^2 \rangle_n \right\} \end{aligned}$$

Analogements

$$\langle x_2^2 \rangle_{S,A} = \frac{1}{2} \left\{ \langle x^2 \rangle_e + \langle x^2 \rangle_n \right\}$$

$$\begin{aligned}
 \langle x_1 x_2 \rangle_{S,A} &= \frac{1}{2} \left\{ \int dx_1 x_1 |t_e(x_1)|^2 \int dx_2 x_2 |t_n(x_2)|^2 \right. \\
 &\quad \pm \int dx_1 x_1 t_e^*(x_1) t_n(x_1) \int dx_2 x_2 t_n^*(x_2) t_e(x_2) \\
 &\quad \pm \int dx_1 x_1 t_n^*(x_1) t_e(x_1) \int dx_2 x_2 t_e^*(x_2) t_n(x_2) \\
 &\quad + \left. \int dx_1 x_1 |t_n(x_1)|^2 \int dx_2 x_2 |t_e(x_2)|^2 \right\} \\
 &= \frac{1}{2} \left\{ \langle x \rangle_e \langle x \rangle_n \pm \underbrace{\langle x \rangle_{en} \langle x \rangle_{ne}}_{|\langle x \rangle_{en}|^2} \right. \\
 &\quad \pm \underbrace{\langle x \rangle_{ne} \langle x \rangle_{en}}_{|\langle x \rangle_{ne}|^2} + \underbrace{\langle x \rangle_n \langle x \rangle_e}_{\text{cancel}} \left. \right\} \\
 &= \langle x \rangle_e \langle x \rangle_n \pm |\langle x \rangle_{ne}|^2 \\
 \rightarrow \langle (x_1 - x_2)^2 \rangle_{S,A} &= \underbrace{\langle x^2 \rangle_e + \langle x^2 \rangle_n - 2 \langle x \rangle_e \langle x \rangle_n}_{\langle (x_1 - x_2)^2 \rangle_D} \\
 &\quad \mp 2 |\langle x \rangle_{ne}|^2
 \end{aligned}$$

L'unico termine da calcolare esplicitamente è
l'ultimo

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$$\begin{aligned}
 \langle x \rangle_{nl} &= \frac{2}{L} \int_0^L dx \times \sin \frac{n\pi}{L} x \sin \frac{l\pi}{L} x \\
 &= \frac{2}{L} \int_0^L dx \times \frac{1}{2} \left[\cos \frac{(n-l)\pi}{L} x - \cos \frac{(n+l)\pi}{L} x \right] \\
 &= \frac{1}{L} \int_0^L dx \times \cos \frac{(n-l)\pi}{L} x - \frac{1}{L} \int_0^L dx \times \cos \frac{(n+l)\pi}{L} x
 \end{aligned}$$

Or

$$\int_0^L dx \times \cos \alpha x = \frac{1}{\alpha} \times \sin \alpha x \Big|_0^L - \frac{1}{\alpha} \int_0^L dx \sin \alpha x$$

$$= \frac{1}{\alpha} \times \sin \alpha x \Big|_0^L + \frac{1}{\alpha^2} \cos \alpha x \Big|_0^L$$

$$\begin{aligned}
 \langle x \rangle_{nl} &= \frac{1}{L} \left\{ \frac{L}{\pi(n-l)} \times \sin \frac{\pi(n-l)}{L} x \Big|_0^L + \frac{L^2}{\pi^2(n-l)^2} \cos \frac{\pi(n-l)}{L} x \Big|_0^L \right. \\
 &\quad \left. - \frac{L}{\pi(n+l)} \times \sin \frac{\pi(n+l)}{L} x \Big|_0^L - \frac{L^2}{\pi^2(n+l)^2} \cos \frac{\pi(n+l)}{L} x \Big|_0^L \right\}
 \end{aligned}$$

$$= \frac{L}{\pi^2} \left\{ \frac{1}{(n-l)^2} \left[\cos[(n-l)\pi] - 1 \right] - \frac{1}{(n+l)^2} \left[\cos[(n+l)\pi] - 1 \right] \right\}$$

$$\cos[(n \pm l)\pi] = (-1)^{l+m}$$

$$\begin{aligned}
 n+l &= 2k+1 \Rightarrow m = 2k+1-n \\
 n-l &= n - (2k+1-n) \\
 &= 2n-2k-1 \\
 &= 2(n+k)-1
 \end{aligned}$$

$$= \frac{L}{\pi^2} \left[(-1)^{l+m} - 1 \right] \left\{ \frac{1}{(n-l)^2} - \frac{1}{(n+l)^2} \right\}$$

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$$\langle (x_1 - x_2)^2 \rangle_s = \langle (x_1 - x_2)^2 \rangle_D - \left| 2 \frac{L}{\pi^2} \left[(-1)^{\ell+n} - 1 \right] \times \left\{ \frac{1}{(n-\ell)^2} + \frac{1}{(n+m)^2} \right\} \right|^2$$

$$\langle (x_1 - x_2)^2 \rangle = \langle (x_1 - x_2)^2 \rangle_D + \left| 2 \frac{L}{\pi} \left[(-1)^{\ell+n} - 1 \right] \left\{ \frac{1}{(n-\ell)^2} - \frac{1}{(n+m)^2} \right\} \right|^2$$

n se $\ell+n$ è pari non posso distinguere
un verso dall'altro

se $\ell+n$ è dispari i bosoni sono "più vicini"
delle particelle distinguibili mentre i fermioni
sono "più lontani",..

$$\begin{aligned} \frac{1}{(n-\ell)^2} - \frac{1}{(n+m)^2} &= \frac{n^2 + m^2 + 2nm - n^2 - m^2 - 2nm}{(n-\ell)^2(n+m)^2} \\ &= \frac{4nm}{(\gamma^2 \gamma^2)^2} \end{aligned}$$