

# Entropy, FT, Irreversibility

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How much irreversible is a process  $\mu_0 \rightarrow \mu_\infty$  ?  
Question seems to require a notion of entropy

Entropy is naturally defined only for equilibrium

Recent controversial proposal is to define (only)  
*entropy creation* and identify it with the phase  
space contraction

$$\dot{x} = f(x)$$

**Chaotic Hypothesis:** *Motion on the attractor can be regarded as “hyperbolic” (Anosov)*

(Ruelle 73, Cohen-G 95)

(*Ambitious* → identification btwn  
Equilibrium and Nonequilibrium ⇒ ergodic)

**Th** *If evolution is hyperbolic a statistics  $\mu$  of motion exists and is unique:*

$$\lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T F(S_t x) \stackrel{def}{=} \int_M \mu(dy) F(y)$$

*“SRB distribution”. Admits a formal expression!*

Extension to nonequilibrium of Boltzmann–Gibbs

+ reversibility, i.e.  $I^2 = 1$ ,  $IS_t = S_{-t}I \rightarrow \mathbf{FT}$

$$\sigma(x) = -\text{diverg}(x) \equiv -\sum \partial_{x_i} f_i(x) \text{ and } \sigma_+ \stackrel{\text{def}}{=} \langle \sigma \rangle_{SRB}$$

$$\text{and } p \equiv p(x) \stackrel{\text{def}}{=} \frac{1}{\tau} \int_{-\frac{\tau}{2}}^{\frac{\tau}{2}} \frac{\sigma(S_t x)}{\sigma_+} dt$$

dimensionless phase space contraction". Then:

**TH: FT** *If evolution hyperbolic+reversible then*  
 $\mu(p(x) \in \Delta) = e^{\tau \max_{p \in \Delta} \zeta(p) + O(1)}$

and

$$\zeta(-p) = \zeta(p) - p\sigma_+ \quad |p| < p^*$$

$$\text{or } \frac{P(-p)}{P(p)} \simeq e^{-p\sigma_+ \tau}$$

(Cohen, Evans, Morriss (exp) Cohen, G (theory)

This is only an example of theorem. One more:

$F_1(x), \dots, F_n(x)$  even/odd:  $F_i(Ix) = \pm F_i(x)$

“patterns”:  $t \rightarrow \varphi_j(t), t \in (-\frac{\tau}{2}, \frac{\tau}{2}), j = 1, \dots, n$

$$\frac{P_\tau \left( F_j(S_t x) \sim \varphi_j(t), p \right)}{P_\tau \left( F_j(S_t x) \sim \pm \varphi_j(-t), -p \right)} = e^{p \sigma_+ \tau + O(1)}, |p| < p^*$$

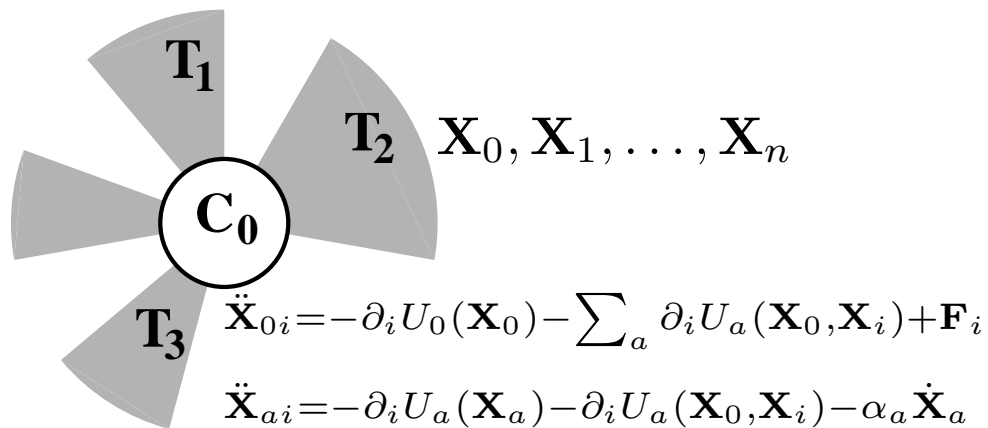
$P_\tau =$  SRB probability; for all  $f, f_j, \varphi_j, n \Rightarrow !$

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If  $\sigma_+ \rightarrow 0 ? \Rightarrow$  (Onsager reciprocity)

$$\partial_A \langle \partial_B \sigma \rangle |_{A,B=0} \equiv \partial_B \langle \partial_A \sigma \rangle |_{A,B=0}$$

Why is  $\sigma(x)$  “entropy”? How to connect with Physics ? (Experim. or Simul.). Example:



$$\alpha_a = \text{Gauss for } K_a = \frac{1}{2} \sum_i \dot{\mathbf{x}}_{ia}^2 = \frac{3}{2} k_B T_a N_a$$

$$\alpha_a = (\text{exercise}) = \frac{W_a - \dot{U}_a}{3N_a k_B T_a}$$

$$W_a = - \sum \dot{\mathbf{x}}_a \cdot \partial_{\mathbf{x}_a} U_a(\mathbf{x}_0, \mathbf{x}_a), W_a = \dot{Q}_a = \text{heat}$$

Divergence

$$\sigma(\mathbf{X}) = (\text{exercise}) = \sigma_0(\mathbf{X}) + \dot{U}(\mathbf{X})$$

$$\sigma_0(\mathbf{X}) = \sum_{a=1}^n \frac{\dot{Q}_a}{k_B T_a} \frac{3N_a - 1}{3N_a} = \sum_{a=1}^n \frac{\dot{Q}_a}{k_B T_a}$$

$$U(\mathbf{X}) = \sum_{a=1}^n \frac{\dot{U}_a}{k_B T_a} \frac{3N_a - 1}{3N_a}$$

If volume  $V^N \times R^{3N}$  changes

$$\sigma^\Gamma(\mathbf{X}) = \sigma_0(\mathbf{X}) + \dot{U} - N \frac{\dot{V}_t}{V_t}$$

EPR,Ja 99:  $\infty$ -thermostats, no  $K_a$ -constraint

Assume  $U_a$  bounded and thermostats *efficient*

$\lim_{t \rightarrow \infty} \frac{1}{t} \int_0^t F(S_t \mathbf{X}) = \int_{\mathcal{F}} F(\mathbf{Y}) \mu(d\mathbf{Y})$ , then

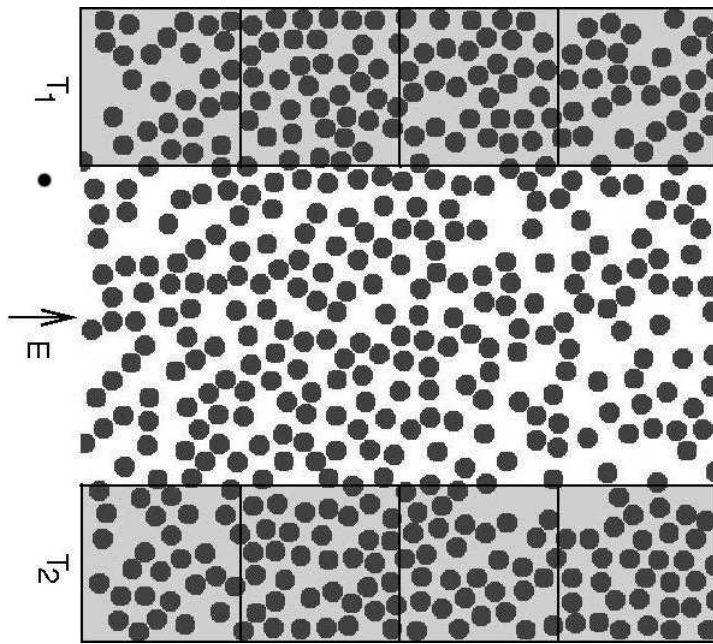
$$\sigma_+ = \langle \sigma \rangle_{\mu} \equiv \langle \sigma_0 \rangle + \lim_{t \rightarrow \infty} \frac{U(t) - U(0)}{t} \equiv \langle \sigma_0 \rangle$$

$$\text{CH} \rightarrow \text{FT} : \zeta(-p) = \zeta(p) - p\sigma_+$$

$$p = p_0 + \frac{U(t) - U(0)}{t\sigma_+} \Rightarrow \text{FT for } p_0!$$

$\frac{U(t) - U(0)}{t} \rightarrow 0$  on large time scale ! (CvZ, BGGZ)

but  $\sigma_+$  and  $p_0$  bear no more reference to thermostats and have time scale of  $O(\text{surface})$



Efficient thermostat ?



## Def Irreversibility time scale for process $\Pi$

$$\varepsilon(X) \stackrel{def}{=} \sigma_0(\mathbf{X}) - N \frac{\dot{V}_t}{V_t}$$

$$Quasi\ static\ \Pi : \mathbf{F}(t) = \mathbf{F}_0 + (1 - e^{-\gamma t})(\mathbf{F}_\infty - \mathbf{F}_0)$$

Stationary  $\mu_0$  evolves to  $\mu_t$  and to  $\mu_\infty$ .

$\mu_{srb,t}$  = SRB parameters “frozen” at value taken at time  $t$ . Let *irreversibility time scale* for  $\Pi$

$$\frac{1}{\tau(\Pi)} \stackrel{def}{=} \frac{1}{N^2} \int_0^\infty \left( \langle \varepsilon_t \rangle_{\mu_t} - \langle \varepsilon_t^0 \rangle_{\mu_{srb,t}} \right)^2 dt$$

If quasi static  $\tau(\Pi) = O(\gamma^{-1} \log \gamma^{-1}) \xrightarrow{\gamma \rightarrow 0} \infty$  (*i.e.*  
 $\rightarrow$  reversible)

If “Joule-Thomson” with vol. doubling at speed  $w$  it is  $\tau(\Pi) = O(\frac{L}{w})$

If “Joule-Thomson” with volume *suddenly* larger  
 $\tau(\Pi) = 0$

If “Joule-Thomson” with volume side from  $L$  to  
 $L(1 + \delta)$  at speed  $w$ :  $\tau(\Pi) = \frac{L}{w\delta}$