

Universal large fluctuations and time reversal

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Equilibrium states as special **stationary** states.

Basic: is extension of classical thermodynamics possible?

Aim: Thermodyn. \Rightarrow general principles **independent** of microscopic properties

Major understanding of Therm. through Stat. Mech. and the assumption of **Gibbs distribution**

Isolated system (in equil.) \Rightarrow probability on $(\mathbf{p}, \mathbf{q}) \in R^{3N} \times V^N$

$$\mu_0(d\mathbf{p}d\mathbf{q}) = \frac{1}{Z} e^{-\beta H(\mathbf{p}, \mathbf{q})}$$

$$H(\mathbf{p}, \mathbf{q}) = \sum_i p_i^2 + \sum_{i,j} V(q_i - q_j) + \text{walls}$$

If $V(\mathbf{q})$ subject to physical cond. **short range and** $V(\mathbf{q}) \geq -B_N$
therm. follows

$$U = (\text{int. energy}) = H, \quad S = (\text{entropy}) = -k_B \frac{\partial}{\partial \beta^{-1}} (\beta^{-1} \log Z)$$

equivalent to

$$dU = dQ + dL, \quad \frac{dQ}{T} = dS \longleftrightarrow \frac{\partial p}{\partial T} = \frac{\partial S}{\partial V}$$

Why ? : Ergodic Hypothesis implies Gibbs Distr. **But**

(1) EH is not a theorem (**in any generality**)

(2) even if true: long way to apply (**how long wait until..? FPU!**)

Alexandrinian way out: proceed “as if”. Rationalize as

(1) stationarity is attained locally (few degrees of freedom)

(2) motion is chaotic (short time scale to stionarity)

Can this be done in NEQ??? where eq. motion **not Hamiltonian**

Nonconservative forces act: energy must be dissipated by
thermostats i.e. environment

Idea: stationary states are attained locally \Rightarrow study small systems \Rightarrow since '980's simulations \Rightarrow built to understand NESS

Problem: *no Gibbs formula*

(a) equation of motion non Hamiltonian,

(b) phase space volume not conserved.

(c) (too) **many** models of thermostats.

Simplest example (electric resistance Drude, 1890's)

$$\dot{q} = p, \quad \dot{p} = E - \alpha(p)\dot{p} - \partial V(q), \quad \alpha(p) = \frac{E \cdot p}{p^2}$$

α s.t. $p^2 = \text{const.}$ or $\frac{1}{2}p^2 + V = \text{const}$

α **not constant** but same friction effect in the average

BUT reversible

More elaborated example in fluids (NS)

$$\dot{\mathbf{u}} + \underline{u} \cdot \underline{\partial} \mathbf{u} = \alpha(\mathbf{u}) \Delta \mathbf{u} - \partial p + \mathbf{f}$$

incompressible and $\int \mathbf{u}^2 = \text{const}$ if

$$\alpha(\mathbf{u}) = \frac{\int \mathbf{f} \cdot \mathbf{u}}{\int (\underline{\partial} \mathbf{u})^2}$$

equivalent to NS if $\langle \alpha \rangle = \nu$ (??). **But reversible !!**.

Many thermostats, typically with unphysical forces: problem “**their equivalence**”. I.e. ensembles equiv.

Idea: reversibility important even in presence of dissipation: Onsager Recip., Green-Kubo formulae, Fluct.-Dissip. theorem (at infinitesimal forcing!!).

Chaotic hypothesis (Cohen-GG): “it can be assumed that motion is chaotic in the math. sense”, i.e. it is “Anosov” (by Ruelle this **implies** a natural formula for stationary states)

Key observation: in all NE thermostat models the phase space contraction rate can be interpreted as **entropy production rate** of the thermostats

$$\varepsilon = \frac{\dot{Q}}{k_B T}$$

- (a) $\langle \varepsilon \rangle \geq 0$ (“friction” \Rightarrow as entropy increase of thermostats)
(b) $\mathbf{p} \stackrel{\text{def}}{=} \frac{1}{t} \int_0^t \frac{\varepsilon(\tau)}{\langle \varepsilon \rangle} d\tau$ **fluctuates** in NESS with average $\langle p \rangle = 1$
(c) (Sinai’s theorem): $Prob(p \in \Delta) = e^{tf(p)}$ for some “large deviation $f(p)$ ”

Fluctuation theorem (Cohen, GG): **reversible & chaotic** \Rightarrow

$$f(p) - f(-p) = \langle \varepsilon \rangle p, \quad \longleftrightarrow \quad \frac{Prob(p \in dp)}{Prob(-p \in dp)} = e^{t\langle \varepsilon \rangle p}$$

no free parameters. If $\langle \varepsilon \rangle \rightarrow 0$ this reduces to **Gree-Kubo**, **Onsager Rec.**, and **FD-theorem**

What about NS? and its **constant friction?** \longleftrightarrow irreversible??

Reversible NS should be equivalent to the irreversible one:
however on a macroscopic scale the fluctuations are not
observable. Philosophy is that reversibility can be observed on
small scale

In NS where can one look for reversibility?

OK theory: friction appears on the Kolmogorov scale

$$\lambda_k = LR^{-\frac{3}{4}}$$

where L = scale over which forcing occurs.

$$\text{Average of } \alpha(\mathbf{u}) = \frac{\int \mathbf{f}\dot{\mathbf{u}}}{\int (\partial\mathbf{u})^2}$$

if \mathbf{u} is the velocity field without the viscosity components
should be equal to \mathbf{v} and its fluctuations should obey the FR

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