

”On some integrable systems:

Brownian motion, normal forms and zeros of polynomials”

Among the first works of Carlo was integration of Hamiltonian

$$H_n(\vec{p}, \vec{x}) = \sum_{i=1}^n \frac{1}{2} p_i^2 + \frac{1}{2} \sum_{i=1}^n \omega^2 x_i^2 + \frac{1}{2} \sum_{i \neq j}^{1,n} g^2 \frac{1}{(x_i - x_j)^2} \quad (*)$$

for $n = 3$, following early (quantum) results of Calogero, [3, 2]:

Besides integrating the classical system also the scattering matrix of the three bodies was calculated and found equal to that of a system of impenetrable points (1970), [10].

H_n : Essentially self-adjoint on $\mathcal{D}_{\lambda, \xi} \subset L_2(\mathbb{R}^n)$

$$\prod_{i=1}^n \prod_{j=i+1}^n (x_i - x_j)^\lambda P(x_1, \dots, x_n) e^{-\xi \sum_{i=1}^n x_i^2} \quad (**)$$

with $\lambda = \frac{1}{2}(\mathbf{1} + (\mathbf{1} + \frac{4g^2}{\hbar^2})^{\frac{1}{2}})$, if $\lambda > 3$, $\xi = \frac{\omega}{2\hbar}$ and 0 unless $x_1 < x_2 < \dots < x_n$, P a symmetric polynomial in x_i .

Calogero then computed the full spectrum of the n body system with quadratic $\frac{1}{2}\omega^2 m \sum_{i<j} (x_i - x_j)^2$ determining the eigenfunctions (1971), [4].

Marchioro & G.: Eigenvectors for (*): same formula (**) with polynomials recursively generated.

Completeness: reduced to Hermite's polynomials completeness.

The H_n eigenvalues are (k_1, \dots, k_n) integers mod permutations).

$$\varepsilon_{k_1, \dots, k_n} = \omega \hbar \sum_{i=1}^n \left(k_i + \frac{1}{2}\right) + \omega \hbar \lambda \frac{n(n-1)}{2}$$

Here H_n is restricted to $\cup_P \mathcal{D}_{\lambda, \xi}^P$ where $\mathcal{D}_{\xi, \lambda}^P =$ image of $\mathcal{D}_{\lambda, \xi}$ under map $(x_1, \dots, x_n) \rightarrow (x_{P_1}, \dots, x_{P_n})$, multiplicity is $n!$.

Partition function: $Z_n = \text{Tr} e^{-\beta H_n}$: and (formally)

$$\begin{aligned} \lim_{\hbar \rightarrow 0} (2\pi \hbar)^n Z_n &\equiv \int e^{-\frac{\beta}{2} \sum_{i=1}^n p_i^2 - \frac{\beta}{2} \sum_{i<j} \frac{g^2}{(x_i - x_j)^2} - \frac{\beta}{2} \omega^2 \sum_{i=1}^n q_i^2} d\vec{p} d\vec{q} \\ &= e^{-\beta \omega g \frac{n(n-1)}{2}} \left(\frac{2\pi}{\beta \omega}\right)^n \end{aligned}$$

Z_n can be **directly computed** from spectrum \Rightarrow classical limit

$$\int e^{-\frac{\beta}{2} \sum_{i=1}^n p_i^2 - \frac{\beta}{2} \sum_{i < j} \frac{g^2}{(x_i - x_j)^2} - \frac{\beta}{2} \omega^2 \sum_{i=1}^n q_i^2} d\vec{p} d\vec{q} = e^{-\beta \omega g \frac{n(n-1)}{2}} \left(\frac{2\pi}{\beta \omega} \right)^n$$

This formula is another contribution of Marchioro[9]: it is non trivial shown by its non inclusion in Gradshteyn-Ryzhik tables.

Apparently after writing to Carlo that it would be included **it was not**.

I think that the referees had trouble (still have?) understanding the delicate analysis of the limit as $\hbar \rightarrow 0$.

Nevertheless the formula has had considerable influence in the literature. I had the chance of collaborating with Carlo in its development.

A sketch of the logic on the formula.

Feynman-Kac for the kernel of $e^{-\beta H_n}$:

formally:

$$K_\beta(\vec{x}, \vec{y}) = \int \left(\prod_{i=1}^n P_{x_i, y_i}(d\omega_i) \right) e^{-V_\beta(\omega)} \leq K_\beta^{g=0}$$

$$V_\beta(\omega) \stackrel{\text{def}}{=} \frac{g^2}{2} \sum_{i \neq j} \int_0^\beta \frac{d\tau}{(\omega_i(\tau) - \omega_j(\tau))^2} + \frac{\omega^2}{2} \sum_{i=1}^n \int_0^\beta \omega_i(\tau)^2 d\tau$$

K_β is the kernel of a semigroup but is it that of $e^{-\beta H_n}$?
 $n = 2$ would be sufficient.

Lemma \rightarrow **theorem**: *Strong L_2 convergence, if $\lambda > 3$:*

$$\begin{aligned} & \lim_{\beta \rightarrow 0} \int d\vec{y} \int P_{\vec{x}, \vec{y}}(d\omega) e^{-\frac{1}{\beta}(V_\beta(\omega) - 1)} \\ &= -\frac{g^2}{2} \sum_{i < j} \frac{1}{(x_i - x_j)^2} - \frac{\omega^2}{2} \sum_i x_i^2 \end{aligned}$$

and K_β is strongly continuous $\frac{1}{2}$ -group $\equiv e^{-\beta H_n}$ on L_2 .

Then the inequality (F-K formula)

$$K_\beta(\vec{x}, \vec{y}) \leq K_\beta(\vec{x}, \vec{y})_{g=0} \stackrel{def}{=} K_\beta^0(\vec{x}, \vec{y})$$

and K_β^0 is **well known**: at $\vec{x} = \vec{y}$ (needed for the trace) it is

$$K_\beta^0(\vec{x}, \vec{x}) = \left(\frac{\omega \hbar}{\pi(1 - e^{-2\beta\omega\hbar})} \right)^{\frac{n}{2}} e^{-n \frac{\beta\omega\hbar}{2}} \prod_{i=1}^n e^{-\frac{\omega}{\hbar} \frac{1 - e^{-\beta\omega\hbar}}{1 + 2e^{-\beta\omega\hbar}} x_i^2}$$

Yields *a priori* bounds leading directly to

$$\begin{aligned} & \lim_{\hbar \rightarrow 0} (2\pi\hbar)^n \text{Tr} e^{-\beta H_n} \\ &= \left(\frac{2\pi}{\beta} \right)^{\frac{n}{2}} \int d\vec{x} e^{-\frac{1}{2} \sum_{i=1}^n \omega^2 x_i^2 + \frac{1}{2} \sum_{i \neq j}^{1,n} g^2 \frac{1}{(x_i - x_j)^2}} \end{aligned}$$

An important remark was that the **same result would follow if the Hamiltonian was integrable** and admitted, in each of the $n!$ sectors of $L_2(\mathbb{R}^n)$ canonical action angle variables

$$(A_1, \dots, A_n, \varphi_1, \dots, \varphi_n), A_j \geq 0 \sim (\vec{p}, \vec{q})$$

and with normal form presented as **integrability conjecture**:

$$\tilde{H}(\vec{A}, \phi) = \sum_{k=1}^n k \omega A_k + \frac{\omega g}{2} n(n-1)$$

By symplectic invariance **THEN** the integral would be

$$\int d\vec{A} d\phi e^{-\beta \sum_{k=1}^n k \omega A_k + \frac{\beta \omega g}{2} n(n-1)}$$

I learnt about the Lax pairs at a workshop and insisted with Moser about H_n should be integrable. I tried to find the pair : failing.

But **Moser** few weeks later discovered H. equations with $\omega = 0$ could be written in terms of matrices

$$M_{ij} = \delta_{ij}(p_i - \omega^2 x_i) + \frac{g^2 \sqrt{-1}}{x_i - x_j}, \quad N_{ij} = \delta_{i \neq j} \frac{1}{(x_i - x_j)^2}$$

and become $\dot{M} = \sqrt{-1} [M, N]$: which implies n eigenvalues of the matrix M are **n consts of motion**, indep. & in involution.

Shortly later Adler realized integrability of Calogero's model [11, 1].

Among other results Moser was able to prove **Marchioro's conjecture** about the n -particles scattering.

This proved **integrability** but not yet the **integrability conjecture**: which was **proved** (much) **later** by Françoise.[6]

Recently, with Françoise and Garrido, a series of papers was then started (actually 2!) to **analyze properties of integrable systems** starting from the most elementary, with long run aim at the Calogero-Marchioro-Moser classical Hamiltonian.

Here I continue towards discussing some remarkable properties that arose, for the **pendulum and the Poincot motions**, and that I think might be the expression of interesting structures

However we remain **still far** from getting close to the CMM-system or to other systems like the Kowalevskaja gyroscope or the Toda lattice.

Hamiltonian of a solid with inertia I_1, I_2, I_3 , in Depri't canonical coord. $(K_z, A, B, \gamma, \varphi, \beta)$, is

$$\tilde{H}(K_z, A, B, \gamma, \varphi, \beta) = \frac{1}{2} \frac{B^2}{I_3} + \frac{1}{2} \left(\frac{\cos^2 \beta}{I_1} + \frac{\sin^2 \beta}{I_2} \right) (A^2 - B^2)$$

with $B = 0, \beta = 0$ the unstable uniform rotation if $I_3 < I_1 < I_2$.

$$J = I_1, \quad \tilde{J} = I_3^{-1} - I_2^{-1}, \quad J_{31} = I_3^{-1} - I_1^{-1}, \quad r^2 \stackrel{\text{def}}{=} \frac{J_{12}^{-1}}{\tilde{J}^{-1}},$$

Theorem: Variables p, q can be constructed so that the Hamiltonian takes the form

$$\frac{A^2}{2I_2} - \frac{x A^2 r^2}{2J_{31}} \mathcal{H}(x, r^2)$$

where $x = pq$ so that evolution is

$$p \rightarrow p e^{g(x, r^2)t}, \quad q \rightarrow q e^{-g(x, r^2)t}, \quad g(x) = \partial_x \mathcal{H}(x)$$

[7, 8]

The normal form $\mathcal{H}(x)$ is **analytic in x** (small) with coefficients which are **polynomials in r^2** :

$$\begin{aligned}\mathcal{H}(x) &= 4x - 2(1+r^2)x^2 - (-1+r^2)^2x^3 - \frac{5(1+r^2)(-1+r^2)^2}{4}x^4 \\ &- \frac{3(-1+r^2)^2}{16}(11+10r^2+11r^4)x^5 - \frac{7(-1+r^2)^2(1+r^2)}{16}(9-2r^2+9r^4)x^6 \\ &- \frac{(-1+r^2)^2}{64}(527+332r^2+330r^4+332r^6+527r^8)x^7 \\ &- \frac{9(1+r^2)(-1+r^2)^2}{512}(1043-548r^2+1058r^4-548r^6+1043r^8)x^8 + \dots\end{aligned}$$

The coefficients can be expressed via combinatorial coefficients and empirically up to **order x^{15}** have all roots on the unit circle.

Attempts to show that the polynomials (at least up to order 15) are of **Lee-Yang type** *i.e.* that can be written as

$$Q(z^2) = \sum_{\sigma_1, \dots, \sigma_n = \pm 1} e^{\sum_{i,j} J_{ij} \left(\frac{\sigma_i + \sigma_j}{2}\right)^2} z^{\sum_i \sigma_i}$$

with $J_{ij} > 0$. Then conjecture implied by **Lee-Yang's theorem**

In this connection there is a theorem

Theorem: (Chen, 1995) Polynomials with symmetric, positive and monotonic coefficients have all roots on the unit circle.[5]

Is it true that such polynomials are Lee-Yang polynomials?

The latter is a conjecture that can be checked rigorously for $n \leq 5$ and we found some preliminary numerical evidence that it should also work at least for $n \leq 8$.

The proof for $n = 4$ emerged from a discussion with A. Giuliani; he also has a proof solving the case $n = 5$ (and found the reference to Chen's theorem).

The polynomials in r^2 appearing as coefficients of x^k in the normal form are sums of products of polynomials P_{n_j} of degrees which add up to $2(k-1) = \sum_j n_j$: an expression for related polynomials can be found in terms of a “tree expansion”: hence have “quite explicit combinatorial expressions” (useful for the conjecture?).

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